Stabilization and asymptotic behavior of the system related to water waves problem

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Outline

- 1. Stability of a class of infinite dimensional systems
- 2. Boundary control problem in rectangle
 - 2.1 Governing equations for gravity waves;
 - 2.2. Well-posedness and stabilization (linearized equation);
 - 2.3. Asymptotic analysis;
- 3. Perspectives

Beginning from an example

Transport equation with boundary control:

$$(*) \begin{cases} \frac{\partial z}{\partial t}(t,x) = -\frac{\partial z}{\partial x}(t,x) & \forall t,x \ge 0, \\ z(t,0) = u(t). \end{cases}$$

Take $U = \mathbb{C}$, $X = L^2[0,\infty)$ and $Z = \mathcal{H}^1[0,\infty)$ Define

$$Az = -\frac{\mathrm{d}z}{\mathrm{d}x} \quad \text{with} \quad \mathcal{D}(A) = \mathcal{H}_0^1[0,\infty);$$
$$B = \delta_0 \in \mathcal{L}(\mathbb{C}, \mathcal{H}^{-1}[0,\infty))$$
$$(*) \iff \dot{z}(t) = Az(t) + Bu(t)$$

A generates the unilateral right shift seimigroup \mathbb{T}_t on X,

$$(\mathbb{T}_t z)(x) = z(x-t) \qquad \forall x \in [0,\infty).$$

B is an admissible control operator for $\mathbb{T}.$ \Longrightarrow (*) is a well-posed boundary control system !

Consider the linear control system:

$$\begin{cases} \dot{z}(t) = Az(t) + Bu(t), \\ z(0) = z_0. \end{cases}$$
(1)

with the state space X (Hilbert) and the input space U (Hilbert). $A: \mathcal{D}(A) \to X$ is skew-adjoint with compact resolvents; $u \in L^2_{\text{loc}}([0,\infty);U)$ and $B \in \mathcal{L}(U,X)$.

• Energy estimate:

$$||z_0||^2 - ||z(t)||^2 = -2\int_0^t \langle Az, z \rangle d\tau - 2\int_0^t \langle u, B^*z \rangle_U d\tau,$$

$$\implies ||z_0||^2 - ||z(t)||^2 = 2\int_0^t ||B^*z||_U^2 d\tau \ge 0,$$

if take $u = -B^*z$, i.e. colocated feedback, s.t. \implies energy non-increasing !

(1) + colocated feedback:

$$\begin{cases} \dot{z}(t) = \mathcal{A}z = (A - BB^*)z(t), \\ z(0) = z_0. \end{cases}$$
(2)

• Types of stability:

1. Exponential stability: $\exists M, \alpha > 0$, s.t.

 $||z(t)|| \leqslant M e^{-\alpha t} ||z_0|| \quad \forall z_0 \in X;$

2. Weak stability: $\lim_{t\to\infty} \langle z(t), x \rangle = 0, \quad \forall x, z_0 \in X;$

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- 3. Strong stability: $\lim_{t\to\infty} z(t) = 0$;
- 4. "Polynomial" stability: $\exists f(t)$ with $\lim_{t\to\infty} f(t) = 0$, s.t. $||z(t)|| \leq |f(t)||z_0||_{\mathcal{D}(A)}$, $z_0 \in \mathcal{D}(A)$

$$||z(t)|| \leq f(t)||z_0||_{\mathcal{D}(A)} \quad z_0 \in \mathcal{D}(A)$$

The undamped system:

$$\begin{cases} \dot{w}(t) = Aw(t), \\ w(0) = z_0. \end{cases}$$
(3)

Proposition 1 (Ammari & Tucsnak '01) If there exists T > 0, such that $\forall z_0 \in \mathcal{D}(\mathcal{A})$,

 $\int_0^T \|B^*w\|_U^2 \mathrm{d}\tau \ge C \|z_0\|_W^2 \quad \text{with} \quad \mathcal{D}(\mathcal{A}) \subset X \subset W,$

then for every $z_0 \in \mathcal{D}(\mathcal{A})$,

$$\|z(t)\|^2 \leqslant t^{-\frac{\theta}{1-\theta}} \|z_0\|_{\mathcal{D}(\mathcal{A})}^2, \quad \text{where} \quad [\mathcal{D}(\mathcal{A}), W]_{\theta} = X.$$

 $\frac{\text{Proposition 2}}{C_0\text{-semigroup }\mathbb{T}} \text{ (Borichev \& Tomilov '10) } \mathcal{A} \text{ generates a bounded}$

$$\|(\mathbf{i}\omega - \mathcal{A})^{-1}\| = O(|\omega|^s) \iff \|\mathbb{T}_t z_0\| = O(t^{-1/s}) \quad s > 0.$$

 Chill, Paunonen, Seifert, Stahn, Tomilov, Analysis & PDE (accepted), 2019.
 Su, Control of water waves and floating body system, PhD thesis, Université de Bordeaux, Dec. 2021.

Statement of the main result:

Let the eigenpair of A be $(i\mu_k, \phi_k)_{k \in J}$. We assume that for every $k, l \in J$ and $k \neq l$, there exist $\alpha, \beta > 0$, s.t.

$$\begin{split} & [H_1] \ \mu_k \neq \mu_l \ (\forall k \neq l), \quad \mu_k = k^{\alpha} + O(k^q) \ (q < \alpha) \ \text{as } k \to \infty; \\ & [H_2] \ \|B^* \phi_k\|_U \geqslant \frac{1}{k^{\beta}}; \end{split}$$

Theorem 1 (Su & Tucsnak 2019)

Assume that (A, B) satisfy $[H_1]$ and $[H_2]$ with 1. $0 < \alpha < 1$ and $\beta \ge 0$, then the solution of the damped system (2) satisfy

$$||z(t)||^2 \leqslant t^{-\frac{\alpha}{\beta-\alpha+1}} ||z_0||^2_{\mathcal{D}(A)}, \qquad \forall z_0 \in \mathcal{D}(A);$$

2. $\alpha \geqslant 1$ and $\beta > 0,$ then we have

$$||z(t)||^2 \leqslant t^{-\frac{\alpha}{\beta}} ||z_0||^2_{\mathcal{D}(A)}, \qquad \forall z_0 \in \mathcal{D}(A).$$

Remark: $\alpha \ge 1$ and $\beta = 0 \implies$ exponential stability !

Main idea of the proof:

- 1. $0 < \alpha < 1$, $\beta \ge 0$: (Resolvent estimate)
 - $[H_1] + [H_2] \Longrightarrow i\mathbb{R} \subset \rho(\mathcal{A})$
 - The structure of the eigenvalues: There exist M, $\gamma > 0$, s.t. for every $\omega \in \mathbb{R}$ and $|\omega| \ge M$, the interval

$$\left[\omega - \gamma \omega^{\frac{lpha - 1}{lpha}}, \omega + \gamma \omega^{\frac{lpha - 1}{lpha}}\right]$$

contains at most one element of $(\mu_k)_{k \in J}$.

• Resolvent estimate:

$$\|(\mathrm{i}\omega - \mathcal{A})^{-1}\| = O(|\omega|^{\frac{2(\beta - \alpha + 1)}{\alpha}}).$$

2. $\alpha \ge 1$, $\beta > 0$: (Observability inequality)

 $\int_{0}^{t} \|B^{*}w\|_{U}^{2} d\tau \ge C \|z_{0}\|_{X_{-\beta/\alpha}}^{2};$ $Y = \mathcal{D}(A) = X_{1}, W = X_{-\beta/\alpha} \Longrightarrow \theta = \frac{\alpha}{\alpha+\beta}.$ Remark: Resolvent estimate also works in case 2!!

2. Boundary control problem of a water waves system

Problem setting:

 \bullet A rectangular domain Ω filled with water:



- Question: How to describe the change of the free surface when imposing a wave maker at the lateral boundary?
- Assumptions:

Small-amplitude gravity water waves; Rigid wave maker; Incompressible, irrotational, inviscid...

2.1 Governing equations

- Notation: ζ(t, x): The elevation of the free surface;
 φ(t, x, y): The velocity potential of the fluid;
 v(t): The velocity produced by the wave maker;
 h(y): Shape function (the profile of the velocity field imposed by the wave maker).
- The governing equations of the water waves system, for every $t \ge 0$, $x \in [0, \pi]$ and $y \in [-1, \zeta(t, x)]$, are

$$\begin{cases} \Delta \phi(t, x, y) = 0, \\ (\partial_t \phi) \big|_{y = \zeta(t, x)} + \frac{1}{2} |\partial_x \phi|^2 \big|_{y = \zeta(t, x)} + \zeta(t, x) = 0, \\ (\partial_{\vec{n}} \phi) \big|_{y = \zeta(t, x)} = \partial_t \zeta(t, x), \\ (\partial_{\vec{n}} \phi) \big|_{x = 0} = -h(y)v(t), \\ (\partial_{\vec{n}} \phi) \big|_{y = -1} = 0 = (\partial_{\vec{n}} \phi) \big|_{x = \pi}. \end{cases}$$
(4)

2.2 Reformulation of the equations

• Dirichlet-Neumann operator $A_0[\zeta]$:

$$\begin{cases} \Delta \Psi(x,y) = 0, \\ \Psi\big|_{y=\zeta(t,x)} = \psi(x), \\ (\partial_{\vec{n}}\Psi)_{x=0} = 0 = (\partial_{\vec{n}}\Psi)\big|_{y=-1} = (\partial_{\vec{n}}\Psi)\big|_{x=\pi}. \\ A_0[\zeta] = \gamma_1 D : \psi \longmapsto (\partial_{\vec{n}}\Psi)\big|_{y=\zeta(t,x)}. \end{cases}$$
(5)

• Neumann-Neumann operator $B_0[\zeta]$:

$$\begin{cases} \Delta \Phi(x,y) = 0, \\ \Phi \big|_{y=\zeta(t,x)} = 0, \\ (\partial_{\vec{n}} \Phi)_{x=0} = h(y)v(t) = \mathcal{V}(t,y), \\ (\partial_{\vec{n}} \Phi) \big|_{y=-1} = 0 = (\partial_{\vec{n}} \Phi) \big|_{x=\pi}. \end{cases}$$

$$B_0[\zeta] = \gamma_1 N : \mathcal{V} \longmapsto (\partial_{\vec{n}} \Phi) \big|_{y=\zeta(t,x)}.$$
(6)

2.2 Reformulation of the equations

Remark: Let $\psi(t, x) := \phi(t, x, \zeta(t, x))$, the water waves system (4) is decomposed into (5) and (6), by verifying

$$\phi(t,x,y)=\Psi(t,x,y)+\Phi(t,x,y).$$

The Zakharov/Craig-Sulem formulation (ZCS): (ζ , ψ)-formulation:

$$\begin{cases} \partial_t \zeta - A_0[\zeta]\psi = B_0[\zeta]\mathcal{V},\\ \partial_t \psi + \zeta + \frac{1}{2}|\partial_x \psi|^2 - \frac{(A_0[\zeta]\psi + B_0[\zeta]\mathcal{V} + \partial_x \zeta \partial_x \psi)^2}{2(1+|\partial_x \zeta|^2)} = 0. \end{cases}$$
(7)

\triangleright Related results for ZCS:

1. In an infinite strip domain without control, i.e. $x \in \mathbb{R}$, $\mathcal{V} = 0$: Local well-posedness: D. Lannes '05;

2. Control of ZCS from the free surface:

T. Alazard et al. '15 (2π -peridoc domain, pressure as control on $\omega \subset \mathbb{R}$.) Boundary control system of ZCS (7) ? Completely open !!! Previous work.

- Reid & Russell '85: Null controllability for infinite time;
- Reid '86: Null controllability in an irregular domain;
- Reid '95: Null controllability in finite time for gravity-capillary;

• Mottelet '00:

Flexible generator: Approximate controllability holds (finite); Rigid generator: No approximate controllability in finite time; strong stability in non-uniform way.

Zero mean spaces:

$$H = \left\{ f \in L^2[0,\pi] | \int_0^{\pi} f(x) \, \mathrm{d}x = 0 \right\},$$

$$\varphi_k(x) = \sqrt{\frac{2}{\pi}} \cos(kx) \qquad k \in \mathbb{N}, \ x \in [0,\pi],$$

$$H_\alpha = \left\{ f \in H | \sum_{k \ge 1} k^{2\alpha} | \langle f, \varphi_k \rangle |^2 < \infty \right\} \qquad \forall \ \alpha \in \mathbb{R},$$

$$\mathcal{H}^1_{\mathrm{top}}(\Omega) = \left\{ f \in \mathcal{H}^1(\Omega) \ | f(x,0) = 0, \ x \in (0,\pi) \right\}.$$

2.3 Linearized water waves equations

• Linearized Dirichlet-Neumann operator $A_0 := A_0[0]$:

$$A_0 = \gamma_1 D : \psi \longmapsto \partial_y \Psi(x, 0)$$

Proposition 3: Dirichlet operator $D \in \mathcal{L}(H_{\frac{1}{2}}, \mathcal{H}^1(\Omega))$ and $A_0: H_1 \to H$ is strictly positive.

$$\begin{aligned} A_0\varphi_k &= \lambda_k\varphi_k \quad \lambda_k = k \tanh(k) \qquad \forall \ k \in \mathbb{N} \\ A_0\psi &= \sum_{k \in \mathbb{N}} \lambda_k \langle \psi, \varphi_k \rangle \varphi_k \qquad \forall \ \psi \in H_1. \end{aligned}$$

• Linearized Neumann to Neumann operator B₁:

$$B_1 = \gamma_1 N : \mathcal{V} \longmapsto \partial_y \Phi(x, 0)$$

Proposition 4: Neumann operator $N \in \mathcal{L}(L^2[-1,0], \mathcal{H}^1_{top}(\Omega))$, $B_1 \in \mathcal{L}(L^2[-1,0], L^2[0,\pi])$. Let $h \in L^2[-1,0]$ with $\int_{-1}^0 h(y) \, \mathrm{d}y = 0$, then, for $v \in \mathbb{C}$,

$$B_0v = vB_1h \Longrightarrow B_0 \in \mathcal{L}(\mathbb{C}, H).$$

2.3.1 Well-posedness

Let
$$A = \begin{bmatrix} 0 & I \\ -A_0 & 0 \end{bmatrix}, \quad B = \begin{bmatrix} 0 \\ B_0 \end{bmatrix},$$

the linearized water waves equations can be recast into

$$\begin{cases} \dot{z}(t) = Az(t) + Bu(t) \\ z(0) = z_0 \end{cases}$$
(8)
$$u(t) = \dot{v}(t), \ z(t) = \begin{bmatrix} \zeta & \dot{\zeta} \end{bmatrix}^{\mathsf{T}}, \ X = H_{1/2} \times H \text{ and } B \in \mathcal{L}(\mathbb{C}, X), \\ X_1 = \mathcal{D}(A) = H_1 \times H_{1/2}. \\ \text{Definition 1: Solution of the linearized water waves } (\zeta, \phi): \\ \text{Given } u \in L^2_{\text{loc}}[0, \infty) \text{ and } h \in L^2[-1, 0], \ \int_{-1}^0 h(y) \, \mathrm{d}y = 0, \\ \phi \in \mathcal{H}^1_{\text{loc}}([0, \infty); \mathcal{H}^1(\Omega)), \quad \zeta \in C([0, \infty); H_{\frac{1}{2}}) \cap C^1([0, \infty); H), \\ \partial_t \phi(t, \cdot, 0) + \zeta(t, \cdot) = 0, \text{ in } L^2_{\text{loc}}([0, \infty); L^2[0, \pi]), \\ \forall \ \Psi \in \mathcal{H}^1(\Omega): \\ \int_0^{\pi} \dot{\zeta}(t, x) \overline{\Psi(x, 0)} \mathrm{d}x - \int_0^{\pi} \dot{\zeta}(0, x) \overline{\Psi(x, 0)} \mathrm{d}x = \\ \int_0^t \int_\Omega \nabla(\partial_t \phi)(\sigma, x, y) \cdot \overline{\nabla \Psi(x, y)} \mathrm{d}x \mathrm{d}y \mathrm{d}\sigma - \int_0^t u(\sigma) \int_{-1}^0 h(y) \overline{\Psi(0, y)} \mathrm{d}y \mathrm{d}\sigma. \end{cases}$$

Theorem 2 (Su, Tucsnak and Weiss 2020)

Let $h \in L^2[-1,0]$ be such that $\int_{-1}^0 h(y) \, dy = 0$. Then $\forall \ u \in L^2_{\text{loc}}[0,\infty)$, $\zeta_0 \in H_{\frac{1}{2}}$ and $w_0 \in H$, the linearized equations admits a unique solution with $\zeta(0) = \zeta_0$ and $\dot{\zeta}(0) = w_0$. Moreover, there exists a well-posed linear control system (\mathbb{T}, Φ) with $X = H_{\frac{1}{2}} \times H$ and $U = \mathbb{C}$ s.t., setting $z_0 = \begin{bmatrix} \zeta_0 \\ w_0 \end{bmatrix}$, we have $z(\tau) = \mathbb{T}_{\tau} z_0 + \Phi_{\tau} u \qquad \forall \ \tau \ge 0$.

Finally, the generator A of \mathbb{T} is skew-adjoint, with domain $\mathcal{D}(A) = H_1 \times H_{\frac{1}{2}}$, and there exists $B \in \mathcal{L}(\mathbb{C}, X)$ s.t. $\forall \tau \ge 0$,

$$\Phi_{\tau} u = \int_0^{\tau} \mathbb{T}_{\tau - \sigma} B u(\sigma) \, \mathrm{d}\sigma \qquad \forall \ u \in L^2_{\mathrm{loc}}[0, \infty).$$

linear water waves \iff well-posed linear control system (A, B)[1] Su, Tucsnak and Weiss, *Systems & Control Letters*, 2020

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Theorem 3 (Stabilizability properties)

For $\boldsymbol{\Sigma}=(A,B)$ introduced in Theorem 2, we have

- Σ is not exponentially stabilizable;
- Σ is strongly stabilizable iff h is a strategic profile, i.e.

$$\int_{-1}^{0} h(y) \cosh \left[k(y+1)\right] \mathrm{d}y \neq 0 \qquad \forall k \in \mathbb{N};$$

• If
$$\inf_{k \in \mathbb{N}} \frac{k}{\cosh k} \left| \int_{-1}^{0} h(y) \cosh \left[k(y+1) \right] \mathrm{d}y \right| > 0,$$

then the system Σ is USSD. Precisely, $F=-B^*$ leads to the closed-loop semigroup \mathbb{T}^{cl} satisfying

$$\|\mathbb{T}_t^{cl} z_0\|_X^2 \leqslant C t^{-1/3} \|z_0\|_{\mathcal{D}(A)}^2 \qquad \forall \ z_0 \in \mathcal{D}(A), \ t \ge 0.$$

Remark: gravity-capillary waves

$$\|\mathbb{T}_t^{cl} z_0\|_{H_{3/2} \times H}^2 \leqslant C t^{-3/2} \|z_0\|_{H_3^b \times H_{3/2}}^2$$

L: The horizontal scale of the domain Ω ;

 h_0 : The typical water depth; a: the order of the surface variation; The shallowness parameter: $\mu = \frac{h_0^2}{L^2} \ll 1$ in shallow water regime;



Dimensionless quantities:

$$\tilde{x} = \frac{x}{L}, \quad \tilde{y} = \frac{y}{h_0}, \quad \tilde{t} = \frac{t}{L/\sqrt{gh_0}}, \quad \tilde{\zeta} = \frac{\zeta}{a}, \quad \tilde{\phi} = \frac{\phi}{aL\sqrt{g/h_0}},$$

"twisted" Laplace operator: $\Delta_{\mu} = \mu \frac{\partial^2}{\partial \tilde{x}^2} + \frac{\partial^2}{\partial \tilde{y}^2}$

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Omitting the tildes, for all $x \in (0, \pi)$ and $y \in (-1, 0)$,

dimensionless version :
$$\begin{cases} \Delta_{\mu}\phi_{\mu}(t,x,y) = 0, \\ \partial_{t}\zeta_{\mu}(t,x) - \frac{1}{\mu}\partial_{y}\phi_{\mu}(t,x,0) = 0, \\ \partial_{t}\phi_{\mu}(t,x,0) + \zeta_{\mu}(t,x) = 0, \\ \partial_{x}\phi_{\mu}(t,0,y) = -h(y)v(t), \\ \partial_{y}\phi_{\mu}(t,x,-1) = 0 = \partial_{x}\phi_{\mu}(t,\pi,y), \end{cases}$$
(9)
$$\iff \begin{cases} \partial_{t}^{2}\zeta_{\mu}(t,x) + \frac{1}{\mu}A_{\mu}\zeta_{\mu}(t,x) = \frac{1}{\mu}B_{\mu}u(t), \\ \zeta_{\mu}(0,x) = \zeta_{0}(x), \quad \partial_{t}\zeta_{\mu}(0,x) = \zeta_{1}(x), \end{cases}$$

Remark: No zero-mean conditions now ! Dimensionless operators:

$$D-N: A_{\mu}; \quad N-N: B_{\mu} \in \mathcal{L}(\mathbb{C}, L^2[0, \pi])$$

Proposition 5: $A_{\mu} : \mathcal{H}^{1}[0,\pi] \to L^{2}[0,\pi]$ and $A_{\mu}\varphi_{k} = \lambda_{\mu,k}\varphi_{k}$ $\lambda_{\mu,k} = \sqrt{\mu}k \tanh(\sqrt{\mu}k) \quad \forall \ k \in \mathbb{N} \cup \{0\},$ $A_{\mu}\eta = \sum_{k \in \mathbb{N}} \lambda_{\mu,k} \langle \eta, \varphi_{k} \rangle \varphi_{k} \quad \forall \ \eta \in \mathcal{H}^{1}[0,\pi].$

Remark: as $\mu \to 0$, $\frac{1}{\mu}A_{\mu} : \frac{1}{\mu}\lambda_{\mu,k} \to k^2 \dashrightarrow -\partial_x^2$?

• Wave equation with Neumann boundary control:

$$\begin{cases} \partial_t^2 \zeta(t,x) - \partial_x^2 \zeta(t,x) = 0, \\ \partial_x \zeta(t,0) = u(t), \quad \partial_x \zeta(t,\pi) = 0, \\ \zeta(0,x) = \zeta_0(x), \quad \partial_t \zeta(0,x) = \zeta_1(x). \end{cases}$$
(11)
$$(11) \Longleftrightarrow \begin{cases} \partial_t^2 \zeta(t,x) + A_w \zeta(t,x) = B_w u(t), \\ \zeta(0,x) = \zeta_0(x), \quad \partial_t \zeta(0,x) = \zeta_1(x). \end{cases}$$
(12)
with $B_w = -\delta_0$ and

 $A_w = -\partial_x^2, \qquad \mathcal{D}(A_w) = \left\{ f \in \mathcal{H}^2[0,\pi] \mid f'(0) = f'(\pi) = 0 \right\},$

Proposition 6:
$$\lim_{\mu\to 0} \frac{1}{\mu} B_{\mu} u = B_w u$$
 in $(\mathcal{H}^1[0,\pi])'$.

Theorem 4 (Su 2020)

 $u \in L^2_{loc}[0,\infty)$ and $\zeta_0 \in \mathcal{H}^1[0,\pi]$, $\zeta_1 \in L^2[0,\pi]$, let ζ_μ be the solution of (10), satisfying $\zeta_{\mu} \in C([0,\infty); \mathcal{H}^{1/2}[0,\pi]) \cap C^1([0,\infty); L^2[0,\pi]).$ Let ζ be the weak solution (12) satisfying $\zeta \in C([0,\infty); \mathcal{H}^1[0,\pi]) \cap C^1([0,\infty); L^2[0,\pi]).$ $\lim_{\mu \to 0} \sup_{t \in [0,\tau]} \|\zeta_{\mu} - \zeta\|_{\mathcal{H}^{1/2}[0,\pi]} = 0,$ $\implies \lim_{\mu \to 0} \sup_{t \in [0,\tau]} \|\partial_t \zeta_\mu - \partial_t \zeta\|_{L^2[0,\pi]} = 0.$ $H_{\alpha} = \{ f \in L^{2}[0,\pi] | \sum_{k \ge 0} (1+|\lambda_{k}|^{2\alpha}) | \langle z, \varphi_{k} \rangle |^{2} < \infty \} \Longrightarrow \begin{cases} H_{\mu,\alpha} \ \lambda_{\mu,k} = \frac{k \tanh(\sqrt{\mu}k)}{\sqrt{\mu}} \\ \mathbb{H}_{-} \ \lambda_{-} = \iota^{2} \end{cases}$

Sketch of the proof: Step 1. System transformation. For $\zeta_0 \in \mathbb{H}_{\frac{1}{2}}$ and $\zeta_1 \in H$, (9) and (11) admit ζ_{μ} and ζ satisfying $\zeta_{\mu} \in C([0,\infty); H_{\mu,\frac{1}{\alpha}}) \cap C^1([0,\infty); H),$ $\zeta \in C([0,\infty); \mathbb{H}_{\frac{1}{2}}) \cap C^1([0,\infty); \mathbb{H}).$ $\alpha_{\mu} := \partial_t \zeta_{\mu}, \qquad \beta_{\mu} := \left(\frac{1}{\mu} A_{\mu}\right)^{1/2} \zeta_{\mu},$ Define $\alpha := \partial_t \zeta, \qquad \beta := A^{1/2} \zeta.$ $\implies \alpha_{\mu}, \beta_{\mu}, \alpha, \beta \in C([0, \infty); H).$ Let $w_{\mu}(t) = \begin{vmatrix} \alpha_{\mu}(t, \cdot) \\ \beta_{\mu}(t, \cdot) \end{vmatrix}$ and $w(t) = \begin{vmatrix} \alpha(t, \cdot) \\ \beta(t, \cdot) \end{vmatrix}$,

$$\begin{cases} \dot{w_{\mu}}(t) = \mathcal{A}_{\mu}w_{\mu}(t) + \mathcal{B}_{\mu}u(t), \\ w_{\mu}(0) = w_{\mu,0}, \end{cases} & \begin{cases} \dot{w}(t) = \mathcal{A}_{0}w(t) + \mathcal{B}_{0}u(t), \\ w(0) = w_{0}, \end{cases}$$

with $X_{\mu}=X=H\times H$, \mathcal{A}_{μ} and \mathcal{A}_{0} are skew-adjoint. Using Fourier series we have

$$\lim_{\mu \to 0} (I - \mathcal{A}_{\mu})^{-1} \begin{bmatrix} f \\ g \end{bmatrix} = (I - \mathcal{A}_{0})^{-1} \begin{bmatrix} f \\ g \end{bmatrix} \quad \text{in } X.$$
$$\lim_{\mu \to 0} (I - \mathcal{A}_{\mu})^{-1} \mathcal{B}_{\mu} u = (I - \mathcal{A}_{0})^{-1} \mathcal{B}_{0} u \quad \text{in } X.$$

Step 2. Convergence of a scattering semigroup.

$$\begin{split} w_{\mu} &= \mathbb{T}_{\mu,t} \, w_{\mu,0} + \Phi_{\mu,t} \, u \quad \text{and} \quad w = \mathbb{T}_t \, w_0 + \Phi_t \, u, \\ \mathscr{A} &= \begin{bmatrix} \mathcal{A} & \mathcal{B}\delta_0 \\ 0 & \frac{d}{d\xi} \end{bmatrix} \longrightarrow \mathfrak{T}_t = \begin{bmatrix} \mathbb{T}_t & \Phi_t \\ 0 & S_t \end{bmatrix} \quad \text{on} \quad X \times L^2_{\omega}[0,\infty). \\ \mathcal{D}(\mathscr{A}) &= \left\{ \begin{bmatrix} x \\ u \end{bmatrix} \in X \times \mathcal{H}^1_{\omega}[0,\infty) \, | \, \mathcal{A}x + \mathcal{B}u(0) \in X \right\}, \end{split}$$

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By Trotter-Kato theorem,

$$\lim_{\mu \to 0} (\mathbb{T}_{\mu,t} x + \Phi_{\mu,t} u) = \mathbb{T}_t x + \Phi_t u \quad \text{in} \quad X \times \mathcal{U}.$$

 $\begin{aligned} \text{Step 3.} & \lim_{\mu \to 0} \zeta_{\mu} = \zeta \text{ in } C^{1}([0,\tau];H) \cap C([0,\tau];H_{\frac{1}{2}}). \\ \|\mathbb{T}_{\mu,t}w_{\mu,0} - \mathbb{T}_{t}w_{0}\|_{X} \leqslant \|\mathbb{T}_{\mu,t}w_{\mu,0} - \mathbb{T}_{\mu,t}w_{0}\|_{X} + \|\mathbb{T}_{\mu,t}w_{0} - \mathbb{T}_{t}w_{0}\|_{X}. \\ & \implies \lim_{\mu \to 0} w_{\mu} = w \quad \text{in} \quad X, \\ & i.e. \quad \lim_{\mu \to 0} \alpha_{\mu} = \alpha \quad \lim_{\mu \to 0} \beta_{\mu} = \beta \quad \text{in} \quad H. \\ & \implies \begin{cases} \lim_{\mu \to 0} \partial_{t}\zeta_{\mu} = \partial_{t}\zeta \quad \text{in} \quad C([0,\tau];H); \\ \lim_{\mu \to 0} \zeta_{\mu} = \zeta \quad \text{in} \quad C([0,\tau);H_{\frac{1}{2}}). \end{cases} \end{aligned}$

 Staffans & Weiss, Lax-Phillips semigroup *Trans.Amer.Math.Society*, 2002.
 Su, Asymptotic behaviour of a linearized water waves system in a rectangle, *Asymptotic Analysis*, in press, 2022.

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3. Perspectives

- ★ What is the next?
- ▷ The system in a bounded domain
 - Control of nonlinear system;
 - Water waves system in a general convex domain;
- ▷ Shallow water convergence
 - The controllability & the dispersive effect;
 - Higher regularity and regular convergence;

Thanks for your attention !

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