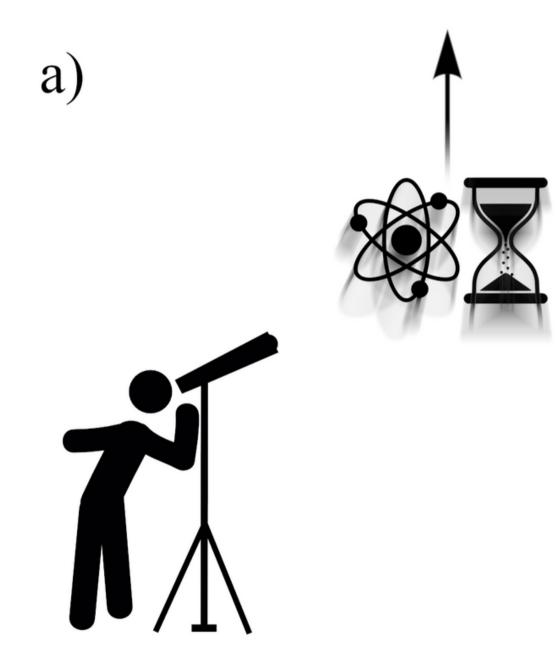
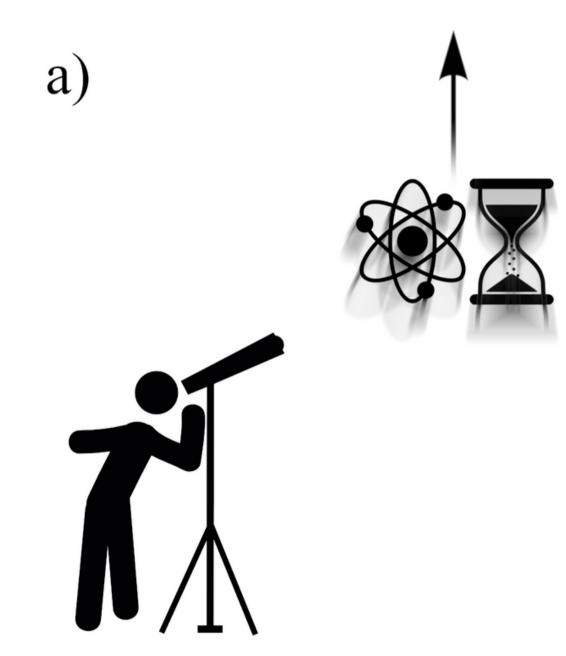
UNIVERSALITY OF QUANTUM TIME DILATION

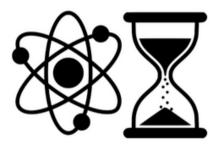
Kacper Dębski

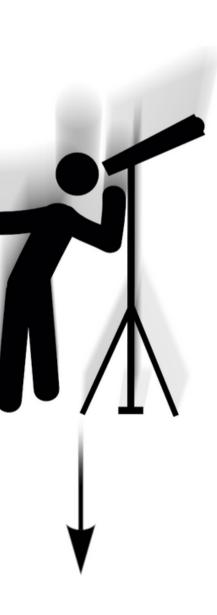


Universal = independent of the mechanism of the clock



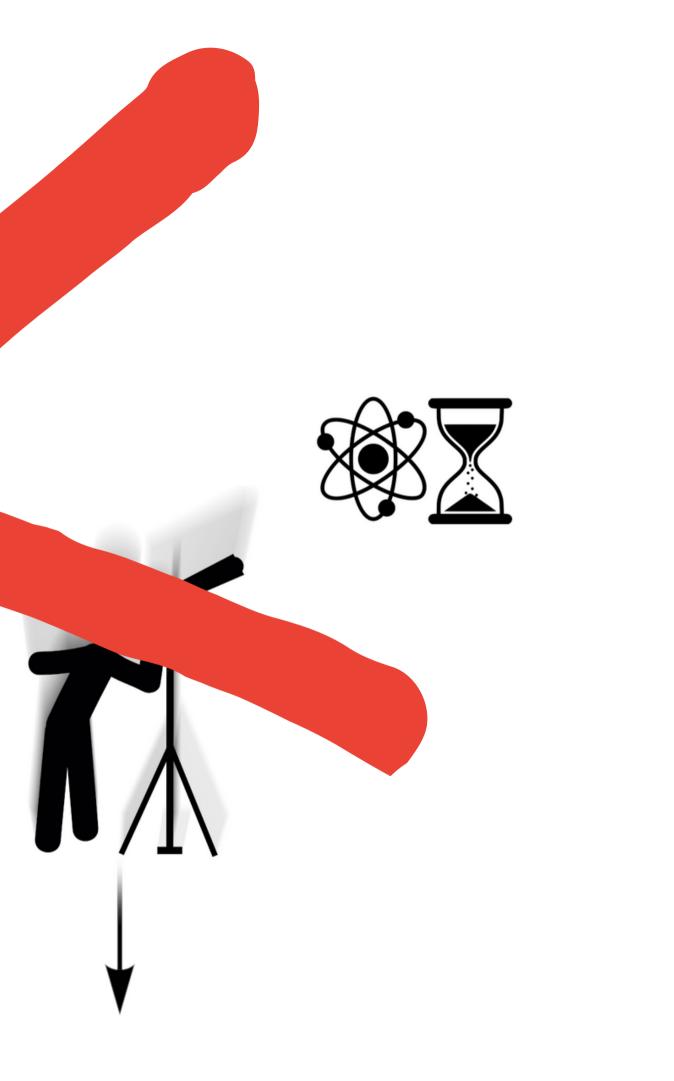


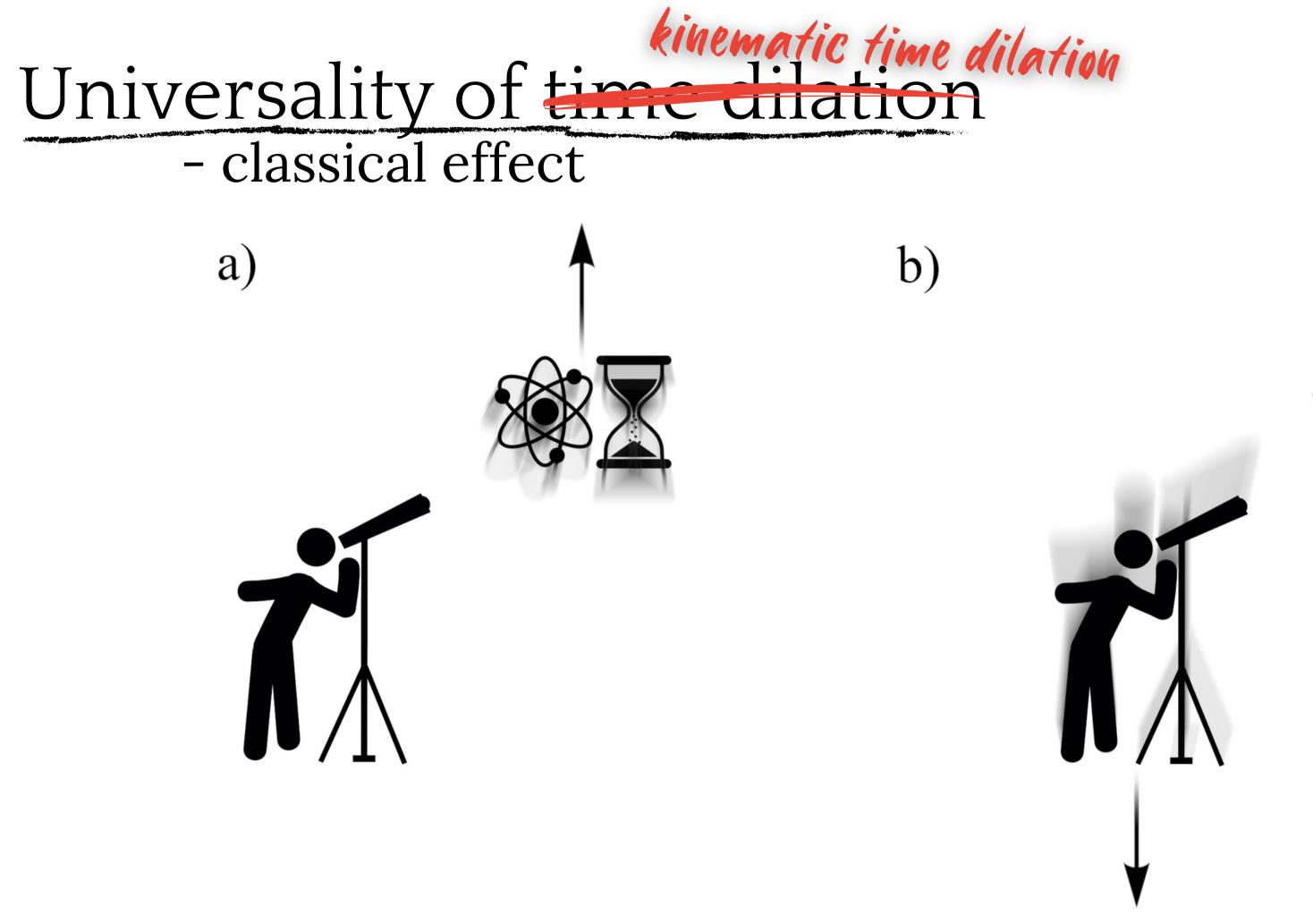


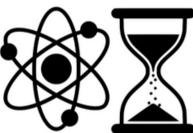


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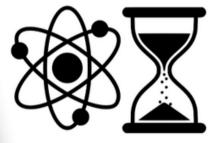
Universality of time dilation - classical effect Class. Quantum Grav. 32 (2015) 175003 (8pp) IOP Publishin Ideal clocks—a convenient fiction Krzysztof Lorek^{1,3}, Jorma Louko² and Andrzej Dragan¹ a ¹ Institute of Theoretical Physics, Faculty of Physics, University of Warsaw, Pasteura 5, 02-093 Warsaw, Poland 02-093 Warsaw, Poland ² School of Mathematical Sciences, University of Nottingham, Nottingham NG7 2RD, UK E-mail: krzysztof.lorek@fuw.edu.pl Received 24 March 2015, revised 1 July 2015 Accepted for publication 3 July 2015 Published 4 August 2015 Abstract We show that no device built according to the rules of quantum field theory one measure preserving along its path. Highly accelerated quantum clocks We show that no device built according to the rules of quantum field theory can measure proper time along its path. Highly accelerated quantum clocks experience the Unrule effect, which inevitably influences their time rate. This can measure proper time along its path. Highly accelerated quantum clocks experience the Unruh effect, which inevitably influences their time rate. This contradicts the concent of an ideal clock, whose rate should only depend on experience the Unruh effect, which inevitably influences their time rate. This contradicts the concept of an ideal clock, whose rate should only depend on the instantaneous velocity. Keywords: Unruh effect, ideal clocks, relativistic quantum information the instantaneous velocity. One immediate prediction of special relativity is that a rate of any clock moving inertially with a value in a larger $1 - y^2$ (we work in natural units such that One immediate prediction of special relativity is that a rate of any clock moving inertially with a velocity v is dilated by a Lorentz factor $\sqrt{1-v^2}$ (we work in natural units such that be a produced by a local velocity reachastic mechanics). The same law applies to a pendulum clock with a velocity v is dilated by a Lorentz factor $\sqrt{1 - v^2}$ (we work in natural units such that $\hbar = c = 1$) independent of the clock's mechanism. The same law applies to a pendulum clock as according to the principle of relativity motion of any of these clocks is $\hbar = c = 1$) independent of the clock's mechanism. The same law applies to a pendulum clock and an atomic clock, as according to the principle of relativity motion of any of these clocks is equivalent to the motion of the observer in the opposite direction. Special relativity cannot and an atomic clock, as according to the principle of relativity motion of any of these clocks is equivalent to the motion of the observer in the opposite direction. Special relativity cannot predict however how an arbitrary clock is affected by non-inertial motion, because different equivalent to the motion of the observer in the opposite direction. Special relativity cannot predict however how an arbitrary clock is affected by non-inertial motion, because different clock for example the rendulum clock or the atomic clock will be affected by acceleration predict however how an arbitrary clock is affected by non-inertial motion, because different clocks, for example the pendulum clock or the atomic clock, will be affected by acceleration in a different way. One often introduces a clock postulate defining a hypothetical ideal clock clocks, for example the pendulum clock or the atomic clock, will be affected by acceleration in a different way. One often introduces a *clock postulate* defining a hypothetical *ideal clock* as a device that measures proper time τ along its arbitrary path according to the formula [1]: in a different way. One often introduces a *clock postulate* defining a hypothetical *ideal clock* as a device that measures proper time τ along its arbitrary path according to the formula [1]: to $v(t) = A\omega \cos \omega t$, $a(t) = -A\omega^2 \sin \omega t$. Let us consider a limit of small amplitudes and high frequencies such that $A \to 0$, $A\omega \to 0$ and $A\omega^2 \to \infty$. In this non-relativistic limit the ³ Author to whom any correspondence should be addressed. 0264-9381/15/175003+08\$33.00 © 2015 IOP Publishing Ltd Printed in the UK





which does not depend on the clock's acceleration at all, only on its instantaneous velocity u(t) The assumption (1) in its idealized form leads to interesting consequences. Consider an which does not depend on the clock's acceleration at all, only on its instantaneous velocity v(t). The assumption (1) in its idealized form leads to interesting consequences. Consider an ideal clock oscillating along a cinercidal path; $x(t) = A \sin \omega t$, where A and ω are the given v(t). The assumption (1) in its idealized form leads to interesting consequences. Consider an ideal clock oscillating along a sinusoidal path: $x(t) = A \sin \omega t$, where A and ω are the given amplitude and frequency, respectively. The clock's velocity and acceleration vary according ideal clock oscillating along a sinusoidal path: $x(t) = A \sin \omega t$, where A and ω are the given amplitude and frequency, respectively. The clock's velocity and acceleration vary according to $y(t) = A\omega \cos \omega t$, $a(t) = -A\omega^2 \sin \omega t$. Let us consider a limit of small amplitude and amplitude and frequency, respectively. The clock's velocity and acceleration vary according to $v(t) = A\omega \cos \omega t$, $a(t) = -A\omega^2 \sin \omega t$. Let us consider a limit of small amplitudes and high frequencies such that $A \rightarrow 0$. $A\omega \rightarrow 0$ and $A\omega^2 \rightarrow \infty$. In this non-relativistic limit the





is

is

the difference in

is

the difference in

time measured by a clock prepared as:

is

the difference in

time measured by a clock prepared as:

quantum superposition

is

the difference in

time measured by a clock prepared as:

quantum superposition VS.

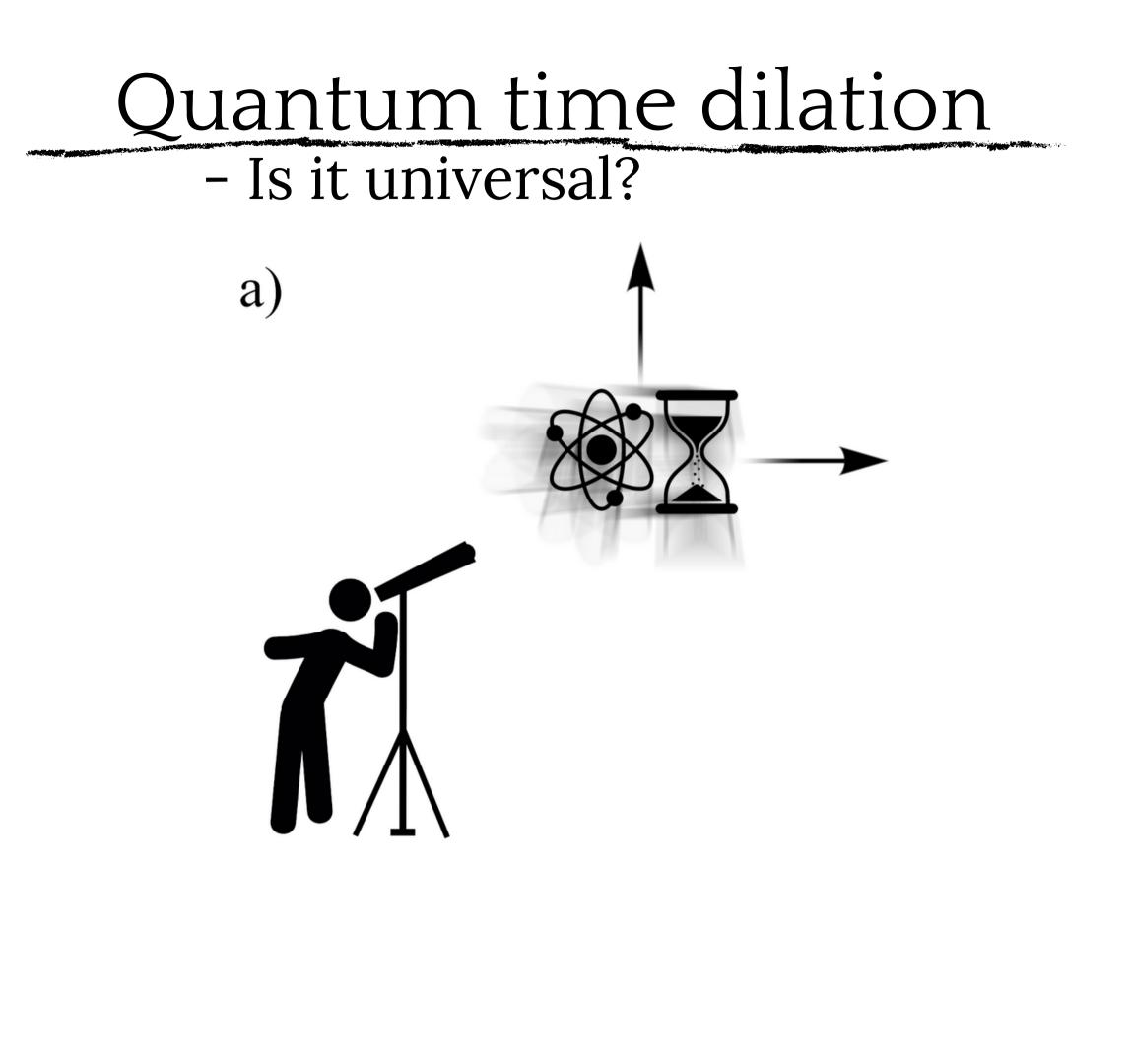
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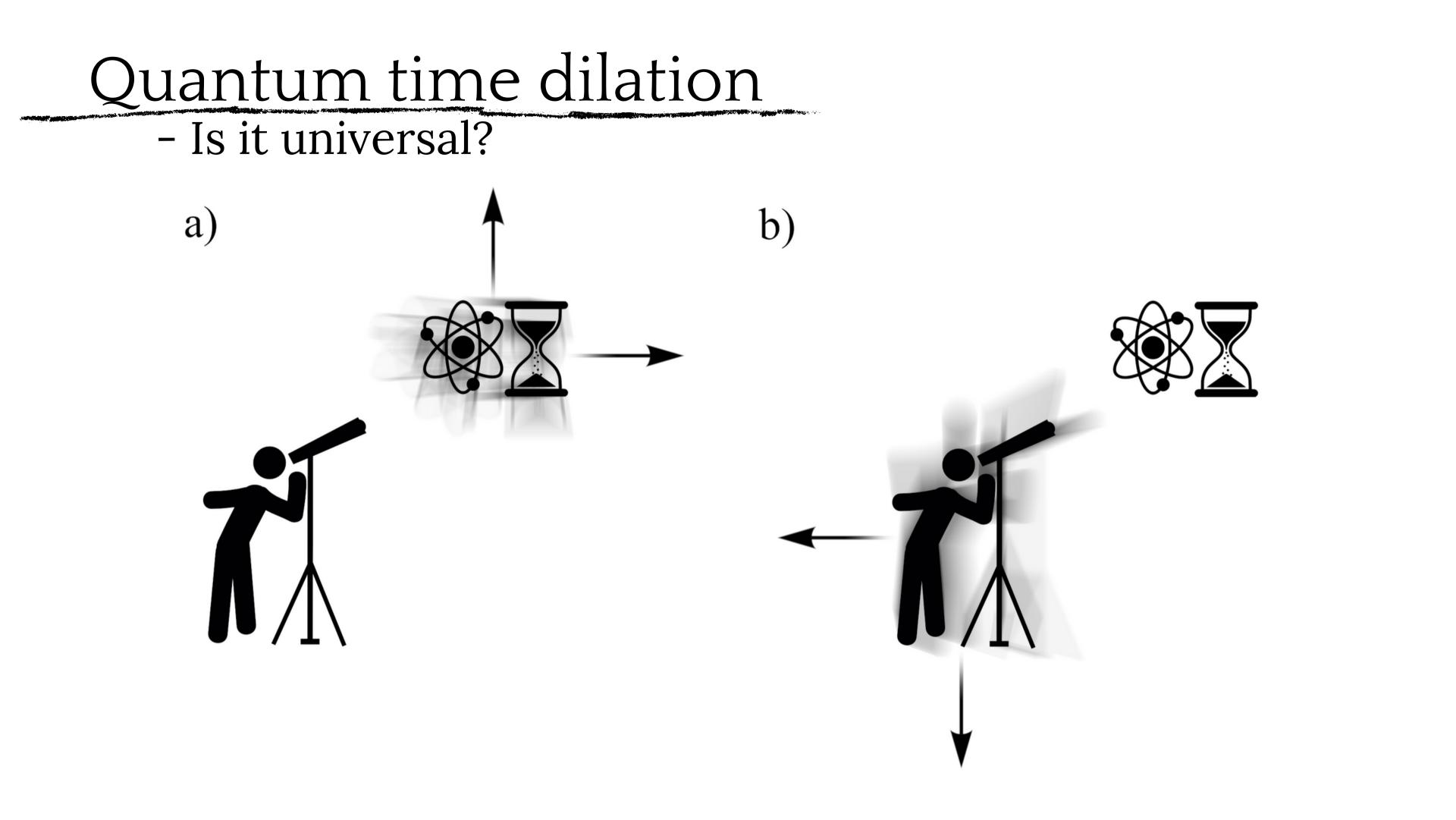
the difference in

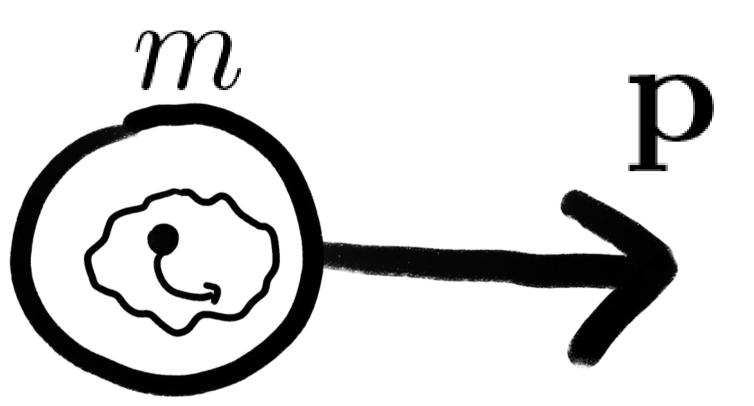
time measured by a clock prepared as:

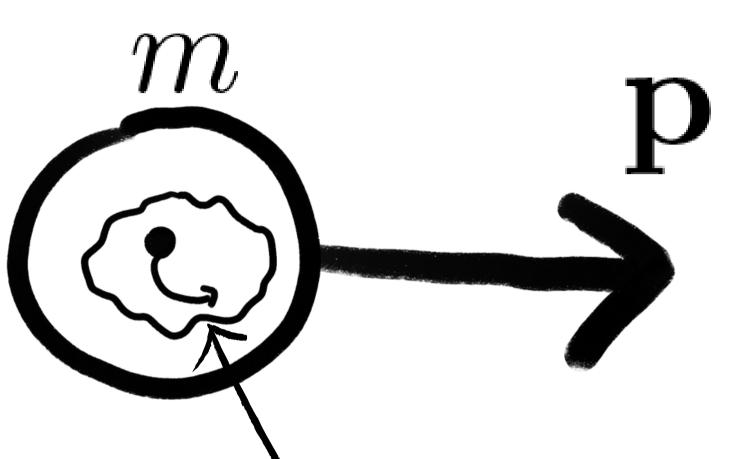
classical mixture quantum superposition VS.

Quantum time dilation - Is it universal?



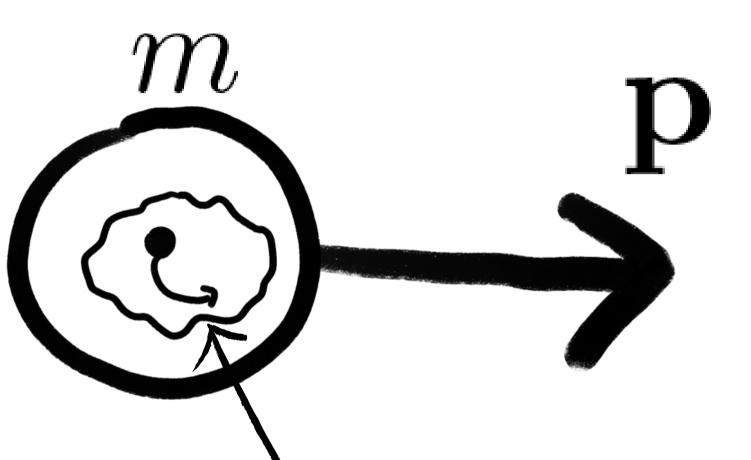




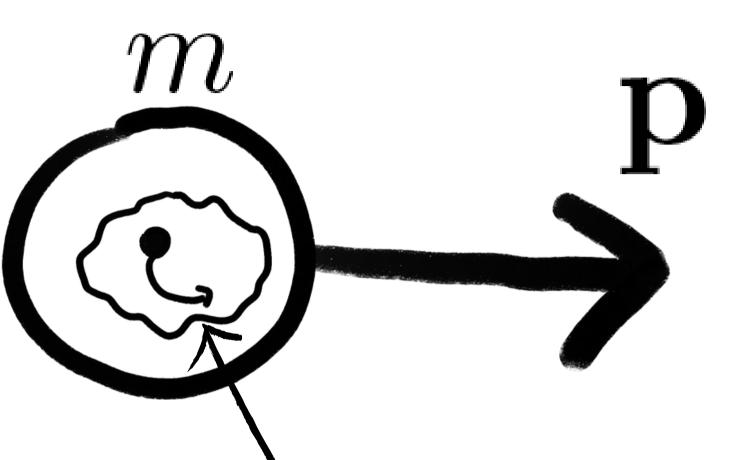


 $E = \sqrt{g_{00}(\boldsymbol{r})} \sqrt{(mc^2)^2 - g^{ij}(\boldsymbol{r})} p_i p_j c^2$

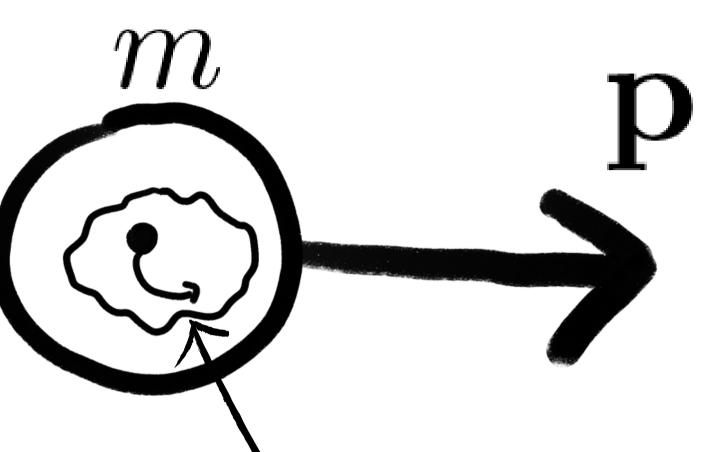
The Schwarzschild metric



 $\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{\text{clock}}\right)^2 - g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2}$



 $\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{\text{clock}}\right)^2 - g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2}$ mass-energy equivalence



 $\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{\text{clock}}\right)^2 - g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2}$

mass-energy equivalence

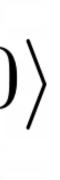
internal degrees of freedom (eg. two-level atom etc.)

ТN

- kinematic

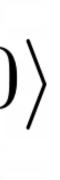
 $\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{\text{clock}}\right)^2 - g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2}$ $\hat{r}(t) = \hat{r}_0$

$|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle\otimes|0\rangle$



$|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle \otimes |0\rangle$

initial state of the system



Quantum time dilation - kinematic $|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle \otimes |0\rangle$

initial state of the system

eigenstate of momentum



Quantum time dilation - kinematic $|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle \otimes |0\rangle$

initial state of the system

 $\rangle \otimes |0\rangle$ initial state of the clock

eigenstate of momentum

Quantum time dilation - kinematic $|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle \otimes |0\rangle$

initial state of the system

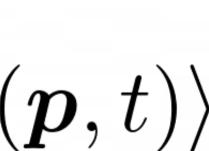
 $\rangle \otimes |0\rangle$ initial state of the clock

some distribution eg. the sum of two Gaussians eigenstate of momentum

Quantum time dilation eigenstate of momentum - kinematic $|\Psi(0)\rangle = \int \mathrm{d}^{3} \boldsymbol{p} \,\psi(\boldsymbol{p}) |\boldsymbol{p}\rangle \otimes |0\rangle$ initial state of the clock initial state of the system $arphi_{ar{\mathbf{p}}_1}$ $|\psi\rangle \sim \cos\theta \,|\varphi_{\bar{\mathbf{p}}_1}\rangle + e^{i\phi}\sin\theta \,|\varphi_{\bar{\mathbf{p}}_2}\rangle$ $\langle \boldsymbol{p} | \varphi_{\bar{\boldsymbol{p}}_i} \rangle = e^{-(\mathbf{p} - \bar{\mathbf{p}}_i)^2 / 2\Delta^2} / \pi^{1/4} \sqrt{\Delta}$ \mathbf{p}_1

$|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle\otimes|0\rangle$

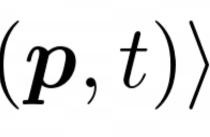
$|\Psi(t)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle\otimes|\phi(\boldsymbol{p},t)\rangle$



$|\Psi(0)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle\otimes|0\rangle$

$|\Psi(t)\rangle = \int \mathrm{d}^{3}\boldsymbol{p}\,\psi(\boldsymbol{p})|\boldsymbol{p}\rangle \otimes |\phi(\boldsymbol{p},t)\rangle$

the evolution only affects a clock's state



 $\hat{\rho}_{\rm clock}(t) = \int d^3 \boldsymbol{p} |\psi(\boldsymbol{p})|^2 |\phi(\boldsymbol{p},t)\rangle \langle \phi(\boldsymbol{p},t)|$

 $\hat{\rho}_{\text{clock}}(t) = \int d^3 \boldsymbol{p} |\psi(\boldsymbol{p})|^2 |\phi(\boldsymbol{p},t)\rangle \langle \phi(\boldsymbol{p},t)|$ we traced out the center-of-mass

degrees of freedom

$$\hat{\rho}_{\text{clock}}(t) = \int d^{3}p |\psi(p)|^{2} |\phi(t)|^{2} |$$

$(\boldsymbol{p},t)\rangle\langle\phi(\boldsymbol{p},t)|$

 $\hat{p}_{clock}(t)$

Quantum time dilation - kinematic $\hat{\rho}_{\text{clock}}(t) = \int d^3 \boldsymbol{p} |\psi(\boldsymbol{p})|^2 |\phi(\boldsymbol{p},t)\rangle \langle \phi(\boldsymbol{p},t)|$ the probability to obtain the time measurement result measurement $\overset{\checkmark}{\mathcal{P}}(\tau) = \operatorname{Tr}\left(\hat{E}(\tau)\hat{\rho}_{\mathrm{clock}}(t)\right)$

$\mathcal{P}(\tau) = \int \mathrm{d}^{3} \boldsymbol{p} \, |\psi(\boldsymbol{p})|^{2} \langle \phi(\boldsymbol{p},t) | \hat{E}(\tau) | \phi(\boldsymbol{p},t) \rangle$

$\mathcal{P}(\tau) = \int \mathrm{d}^{3} \boldsymbol{p} \, |\psi(\boldsymbol{p})|^{2} \langle \phi(\boldsymbol{p}, t) | \hat{E}(\tau) | \phi(\boldsymbol{p}, t) \rangle$ \uparrow

a weighted average

 $\mathcal{P}(\tau) = \int d^3 \boldsymbol{p} |\psi(\boldsymbol{p})|^2 \langle \phi(\boldsymbol{p}, t) | \hat{E}(\tau) | \phi(\boldsymbol{p}, t) \rangle$ a weighted average contribution from the

eigenstate of momentum

 $\mathcal{P}(\tau) = \int \mathrm{d}^{3} \boldsymbol{p} \, |\psi(\boldsymbol{p})|^{2} \langle \phi(\boldsymbol{p},t) | \hat{E}(\tau) | \phi(\boldsymbol{p},t) \rangle$

universal Contribution from the eigenstate of momentum

a weighted average

 $\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{clock}\right)^2 - g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2}$

 $\hat{p}_i = 0$

 $\widetilde{\mathcal{P}}(\tau) = \int \mathrm{d}^{3} \boldsymbol{r} \, |\widetilde{\psi}(\boldsymbol{r})|^{2} \langle \widetilde{\phi}(\boldsymbol{r},t) | \hat{E}(\tau) | \widetilde{\phi}(\boldsymbol{r},t) \rangle$

contribution from the eigenstate of position



 $\widetilde{\mathcal{P}}(\tau) = \int \mathrm{d}^{3} \boldsymbol{r} \, |\widetilde{\psi}(\boldsymbol{r})|^{2} \langle \widetilde{\phi}(\boldsymbol{r},t) | \hat{E}(\tau) | \widetilde{\phi}(\boldsymbol{r},t) \rangle$





eigenstate of position

Kinematic/gravitational quantum time dilation depends only on the probability distribution.

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It can be **simulated** by **classical states**.

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It can be **simulated** by **classical states**.

Is it possible to **find** an effect that depends on

combined distributions

of position and

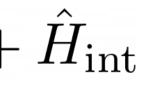
momentum?

$$\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{\text{clock}}\right)^2} \cdot$$

 $-g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2$

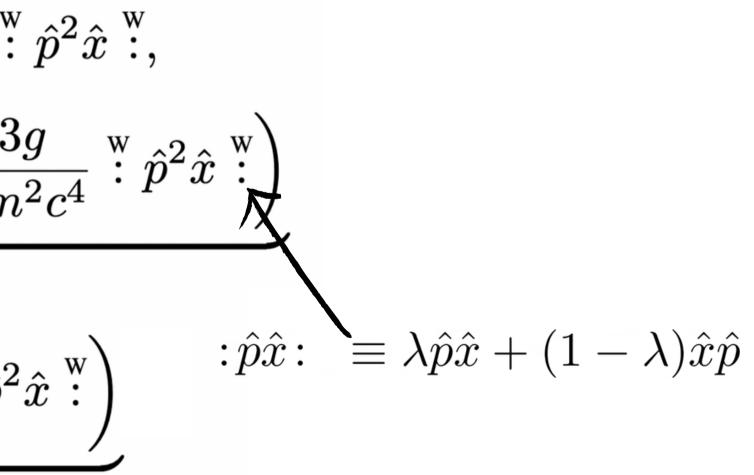
$$\hat{H} = \sqrt{g_{00}(\hat{\boldsymbol{r}})} \sqrt{\left(mc^2 + \hat{H}_{clock}\right)^2}$$
$$\int_{\hat{H}} \hat{H} \approx \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{p}}) + \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{p}}) + \hat{H}_{clock} + \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{p}}) + \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{p}}) + \hat{H}_{clock} + \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{p}}) + \hat{H}_{clock} + \hat{H}_{c$$

 $-g^{ij}(\hat{\boldsymbol{r}})\hat{p}_i\hat{p}_jc^2$



$$\hat{H} \approx \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{p}}) + \hat{H}_{int}$$

$$\begin{split} \hat{H}_{\rm cm} &\equiv mc^2 + \frac{\hat{p}^2}{2m} + mg\hat{x} + \frac{3g}{2mc^2} \stackrel{\rm w}{:} \\ \hat{H}_{\rm int} &\equiv \hat{H}_{\rm clock} \underbrace{\left(-\frac{\hat{p}^2}{2m^2c^2} + \frac{g\hat{x}}{c^2} - \frac{3g}{2m^2c^2} \right)}_{\hat{V}_1} \\ &+ \hat{H}_{\rm clock}^2 \underbrace{\left(\frac{\hat{p}^2}{2m^3c^4} + \frac{3g}{2m^3c^6} \stackrel{\rm w}{:} \hat{p}^2 \right)}_{\hat{V}_2} \end{split}$$



quantum time dilation

$$\hat{H} \approx \hat{H}_{clock} + \hat{H}_{cm}(\hat{\boldsymbol{r}}, \hat{\boldsymbol{j}})$$

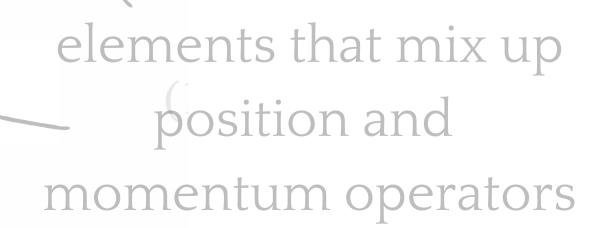
 $\hat{\boldsymbol{p}}) + \hat{H}_{\text{int}}$

elements that mix up \hat{x} position and momentum operators

quantum time dilation $H \approx H_{e} + \hat{H}_{e} + \hat{H}_{e}$ $\hat{H}_{cm} \equiv mc^{2} + \hat{p}^{2} + \tilde{p}^{2} +$ + \hat{H}^2_{clock} (FUNCTION elements that mix up







quantum time dilation

$$\hat{\rho}_{\text{clock}}(t) \propto \text{Tr} \left[\stackrel{\text{w}}{:} \hat{p}^2 \hat{x} \stackrel{\text{w}}{:} \hat{\rho}_{\text{cm}}(0) \right] = \frac{1}{3} \text{Tr} \left[\left(\frac{1}{2} \right) \right] = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left(\frac{1}{2} \right) = \frac{1}{3} \text{Tr} \left[\frac{1}{2} \right] \left($$

$\hat{p}^2 \hat{x} + \hat{p} \hat{x} \hat{p} + \hat{x} \hat{p}^2) \hat{\rho}_{\rm cm}(0)$

 $\hat{x}^{p(x-x')}\langle x|\hat{\rho}_{\rm cm}(0)|x'\rangle$

2-point correlation function

Conclusions:

- QTD depends on a weighted average of probabilities (1-point correlation function)
- Kinematic QTD is UNIVERSAL
- Gravitational QTD is NOT
- There is a higher-order effect that can't be simulated by classical states and that depends on a 2-point correlation function