How Measurement Realizes Quantum Vacuum Ambiguities

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In collaboration with Álvaro Álvarez-Domínguez, José A. R. Cembranos, Luis J. Garay, Mercedes Martín-Benito, Jose M. Sánchez Velázquez



Analogue Gravity in 2023, Benasque – May 2023



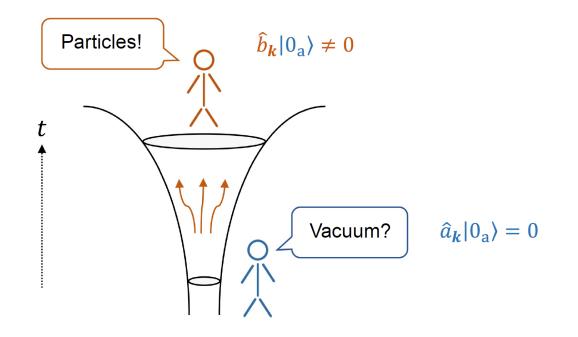


Introduction

- QFT in the presence of an external, time-dependent agent
 - Particle production (even from vacuum)
 - Vacuum/particle notion is ambiguous

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- QFT in the presence of an external, time-dependent agent
 - Particle production (even from vacuum)
 - Vacuum/particle notion is ambiguous
- Some scenarios
 - Expansion of spacetime
 - Schwinger effect



In these processes, one often wonders...

How many 'particles' have been produced?

What a problematic question...

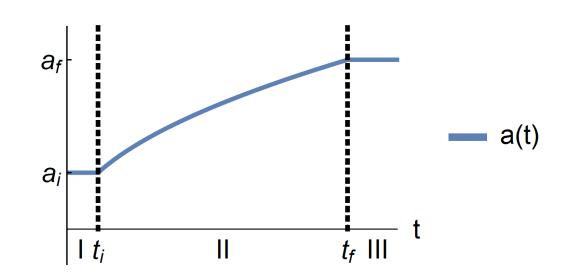
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Consider the following realization

- Preferred vacuum in I, $\hat{a}_{k}|0_{a}\rangle = 0$
- There is no preferred vacuum in II
- Preferred vacuum in III, $\hat{b}_{k}|0_{b}\rangle = 0$



In these processes, one often wonders...

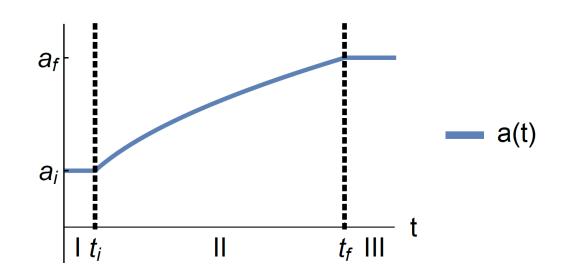
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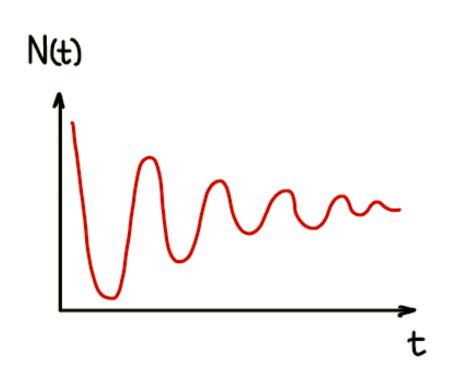
Particle production at t_f is $\langle 0_a | \hat{b}_k^+ \hat{b}_k | 0_a \rangle \neq 0$

But... How many particles are at $t < t_f$?



Depends on the choice of vacuum

• Vacua minimizing energy



Depends on the choice of vacuum

- Vacua minimizing energy
- Adiabatic vacua

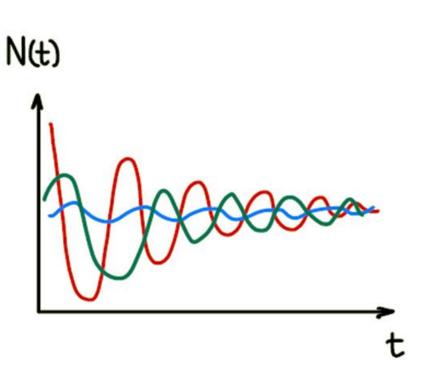
N(t) τ

Depends on the choice of vacuum

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• ...

• Vacua minimizing time oscillations of particle number

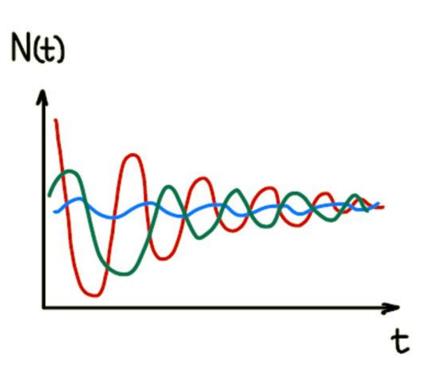


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Can we use the fact that we want to measure to give an operational meaning to some of these notions of 'particle'?

Schwinger effect in 1+1

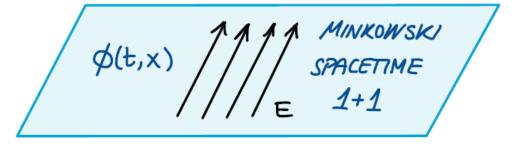
• Scalar field $\phi(t, x)$ in the presence of a homogeneous electric field E(t)

$$\phi(t,x) = \frac{1}{\sqrt{2\pi}} \int dk \ e^{ikx} \phi_k(t)$$

• EOM in *k*-space yields the mode equation

$$\ddot{\phi}_k(t) + \frac{\omega_k(t)^2 \phi_k(t)}{\downarrow} = 0$$

$$\downarrow$$
Time-dependent



Schwinger effect in 1+1

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MINKOWSKI SPACETIME - 1+1

 $\phi(t,x)$

• We can expand the field with each particular solution of the mode equation,

$$\phi(t,x) = \frac{1}{\sqrt{2\pi}} \int dk \left[a_k e^{ikx} \varphi_k(t) + b_k^* e^{-ikx} \varphi_k^*(t) \right]$$
$$= \frac{1}{\sqrt{2\pi}} \int dk \left[c_k e^{ikx} \zeta_k(t) + d_k^* e^{-ikx} \zeta_k^*(t) \right]$$

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Creation and annihilation ops.

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- In Minkowski, Poincarè invariance restricts the choices to plane waves,

$$\varphi_k(t) \sim e^{-i\omega_k t} \rightarrow |0_M\rangle; \ \hat{a}_k^M, \hat{b}_k^M$$

But in the presence of the electric field...

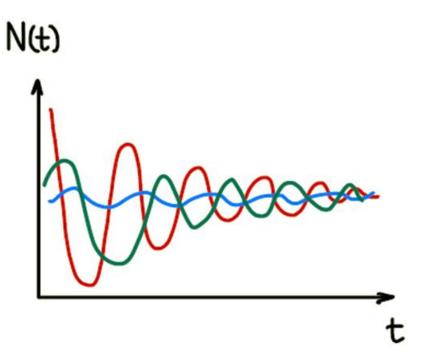
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Depends on the choice of vacuum

- Vacua minimizing energy
- Adiabatic vacua

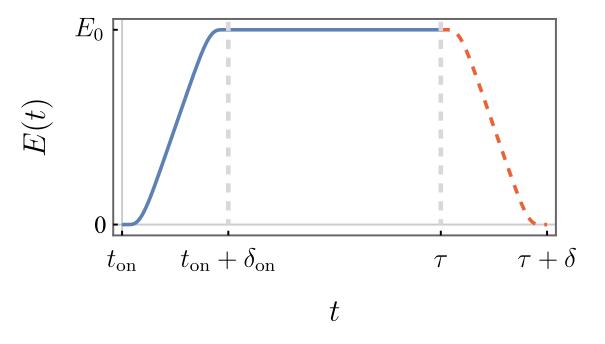
• ...

• Vacua minimizing time oscillations of particle number



IN THE 'LAB'

• Smooth switch-on and switch-off

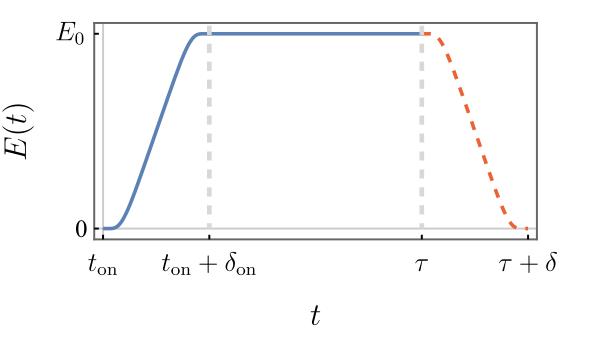


IN THE 'LAB'

- Smooth switch-on and switch-off
- Preferred notions of vacua in the asymptotic past and future

$$\varphi_k^{\text{in}}(t) \sim e^{-i\omega_k^{\text{in}}t} \Rightarrow |0_{\text{in}}\rangle$$

 $\varphi_k^{\rm out}(t) \sim e^{-i\omega_k^{\rm out}t} \Rightarrow |0_{\rm out}\rangle$

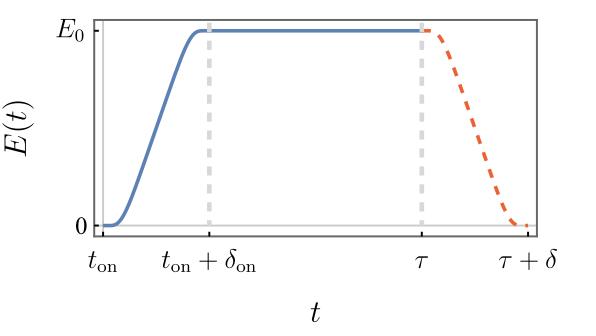


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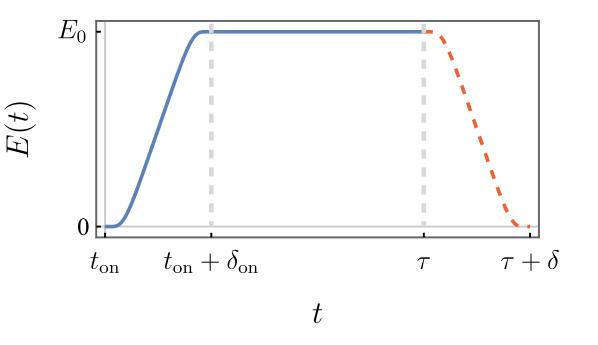
If our system is in $|0_{in}\rangle$, how many particles have been produced at $\tau + \delta$?

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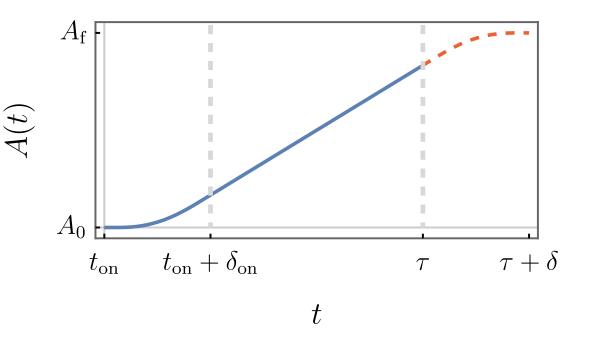
'Measured' particle number (density) $N_{\tau}^{\exp} = \langle 0_{\text{in}} | \hat{a}_{k}^{+,\text{out}} \hat{a}_{k}^{\text{out}} | 0_{\text{in}} \rangle + \langle 0_{in} | \hat{b}_{k}^{+,\text{out}} \hat{b}_{k}^{\text{out}} | 0_{\text{in}} \rangle$

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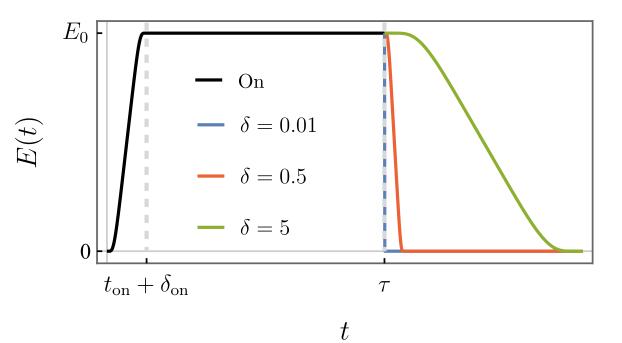
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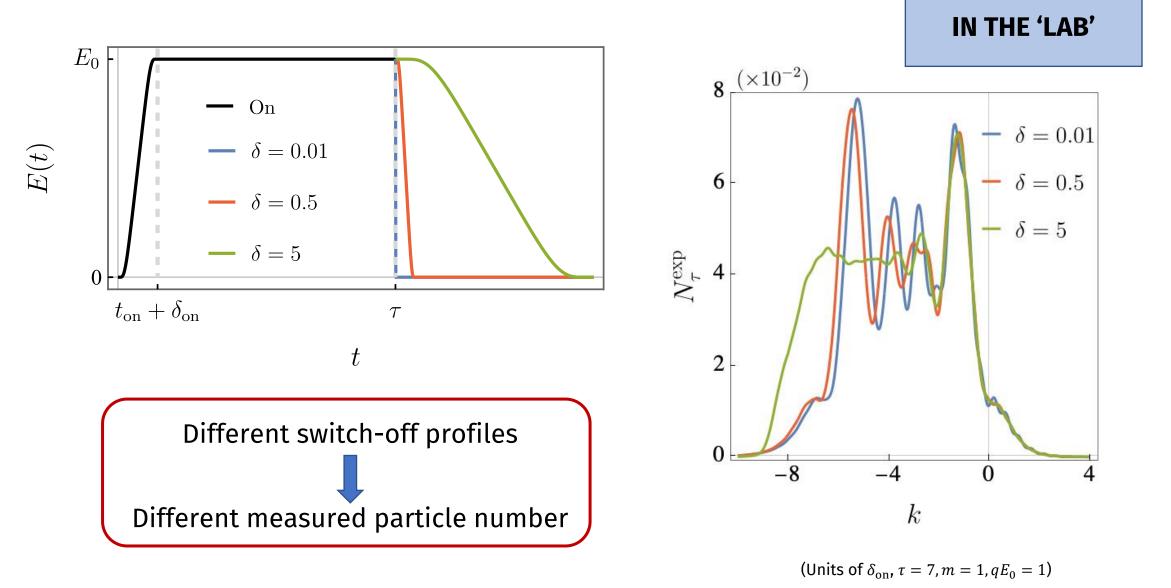


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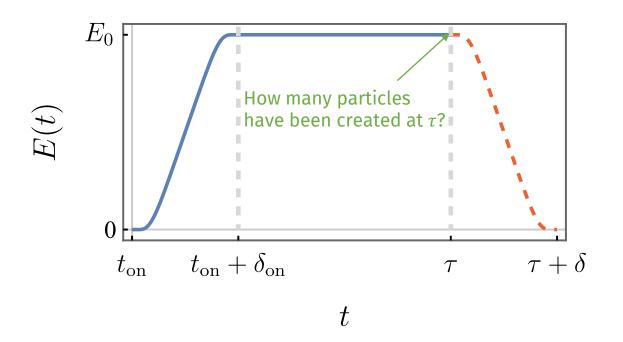
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IN THE 'LAB'





Theoretical particle number



Theoretical particle number

IN THE OFFICE

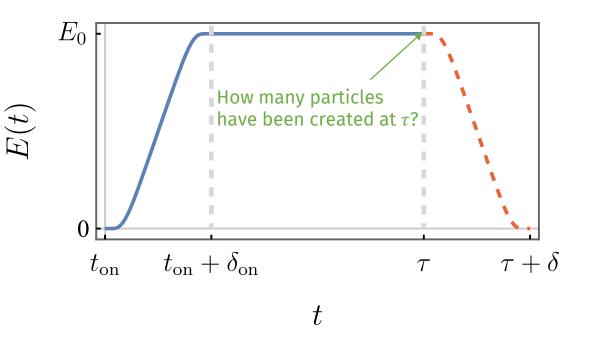
 Preferred notion of vacua in the asymptotic past

 $\varphi_k^{\rm in}(t) \sim e^{-i\omega_k^{\rm in}t} \Rightarrow |0_{\rm in}\rangle$

• We have to make a choice of φ_k^{τ}

 $\varphi_k^\tau(t) \, \Rightarrow |0_\tau\rangle$ (Initial conditions at τ but defined globally)

If out system is in $|0_{in}\rangle$, how many particles have been produced at τ ?



Theoretical particle number (density)

$$N(\tau) = \langle 0_{\rm in} | \hat{a}_k^{+,\tau} \hat{a}_k^{\tau} | 0_{\rm in} \rangle + \langle 0_{in} | \hat{b}_k^{+,\tau} \hat{b}_k^{\tau} | 0_{\rm in} \rangle$$

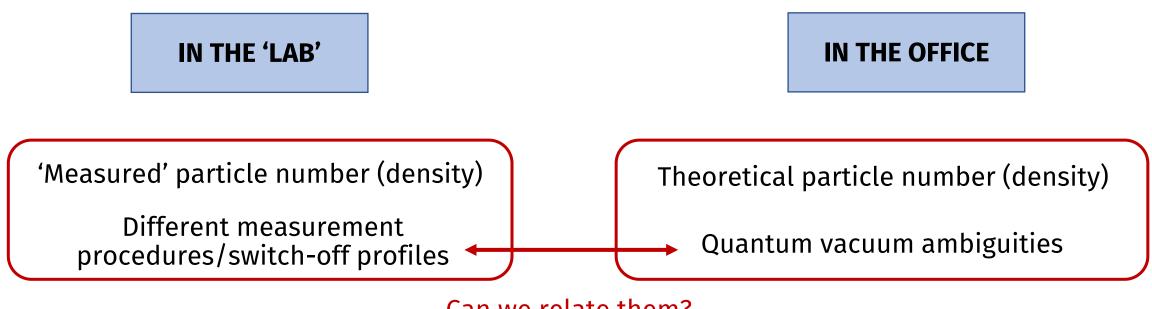
IN THE 'LAB'

IN THE OFFICE

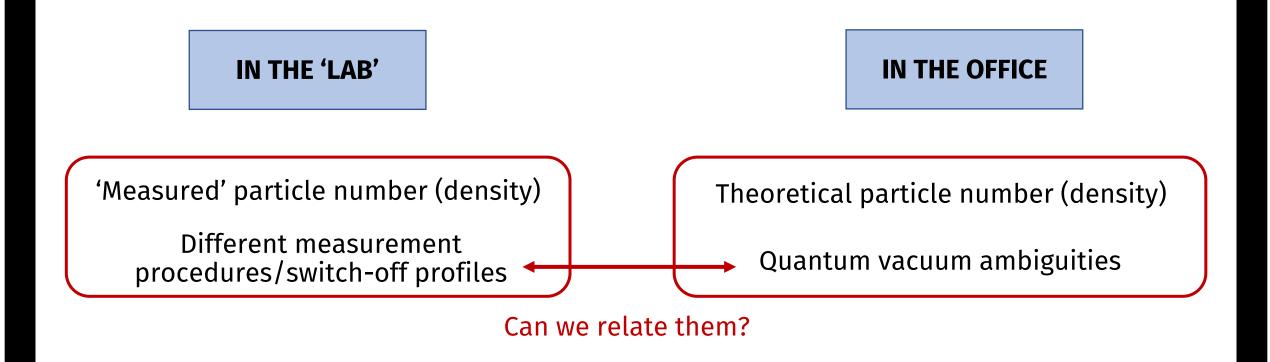
'Measured' particle number (density)

Different measurement procedures/switch-off profiles Theoretical particle number (density)

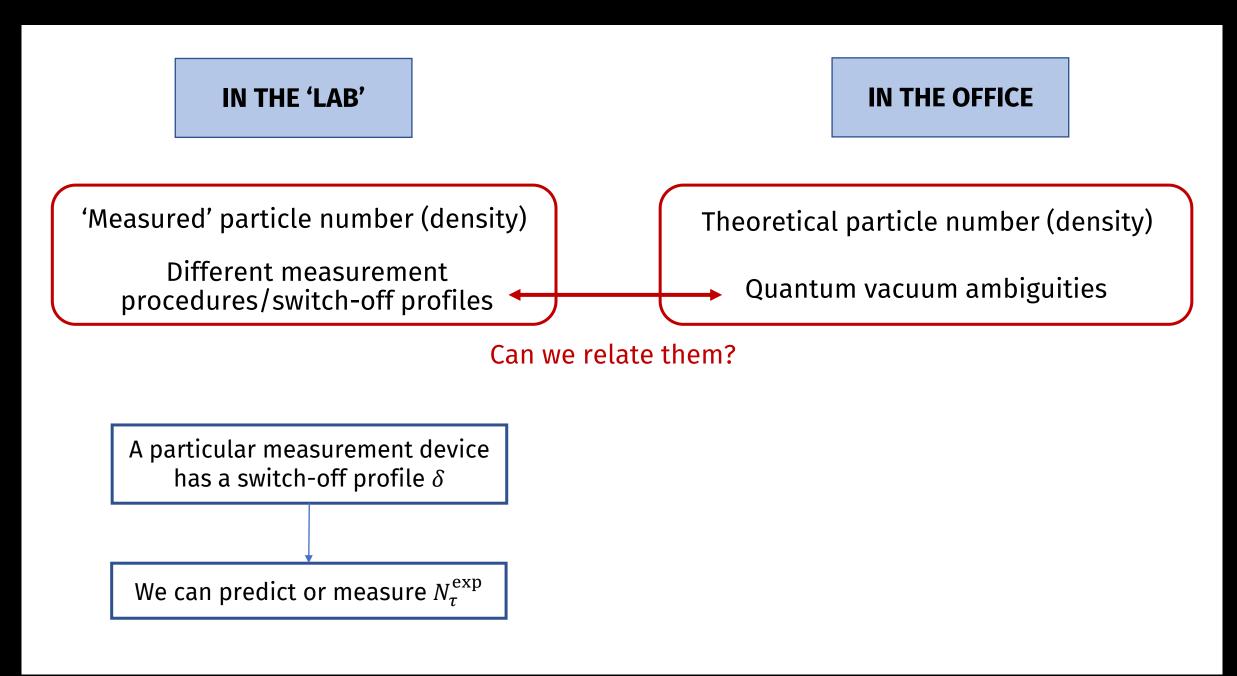
Quantum vacuum ambiguities

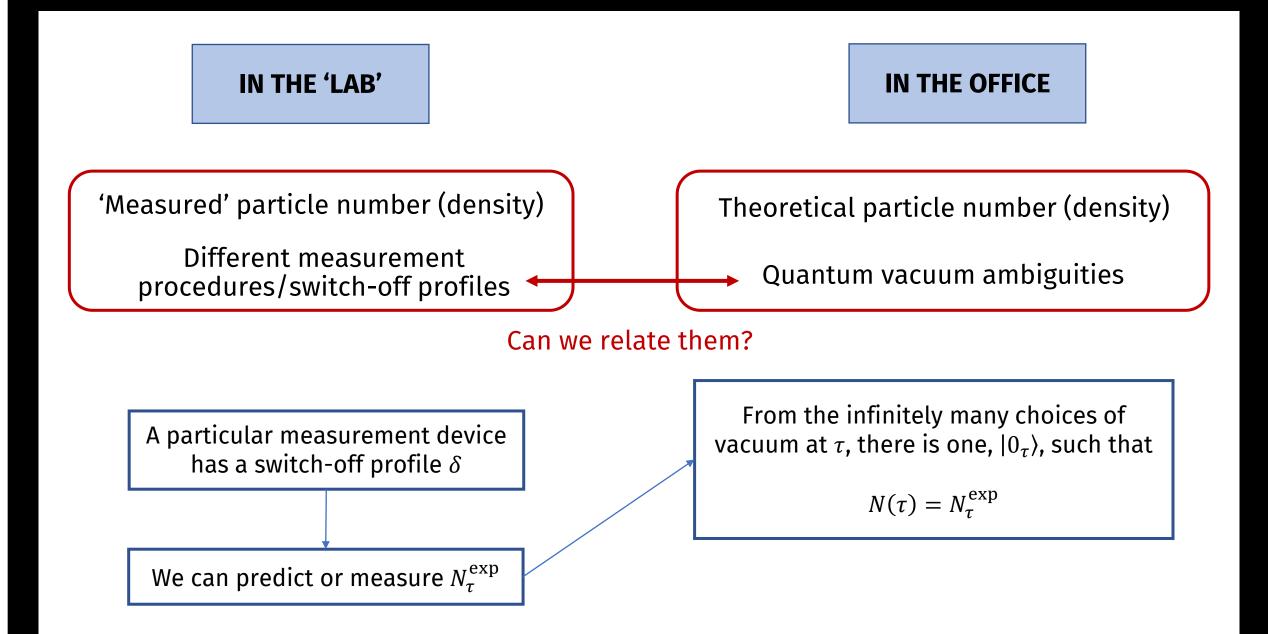


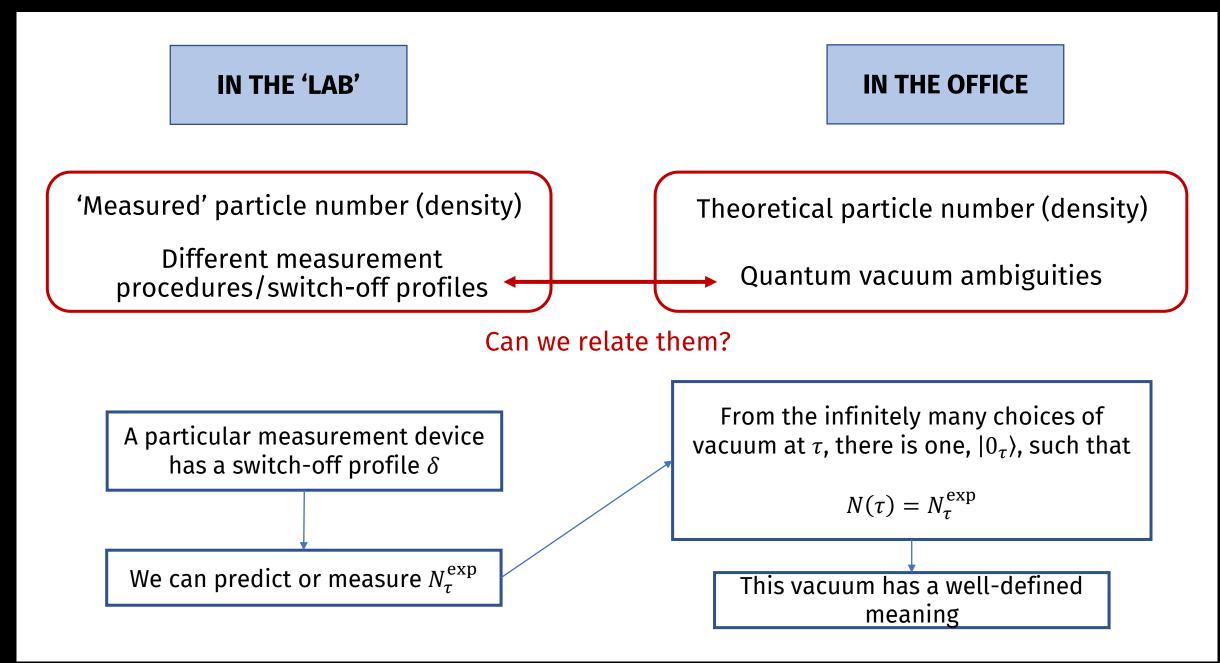
Can we relate them?



A particular measurement device has a switch-off profile δ

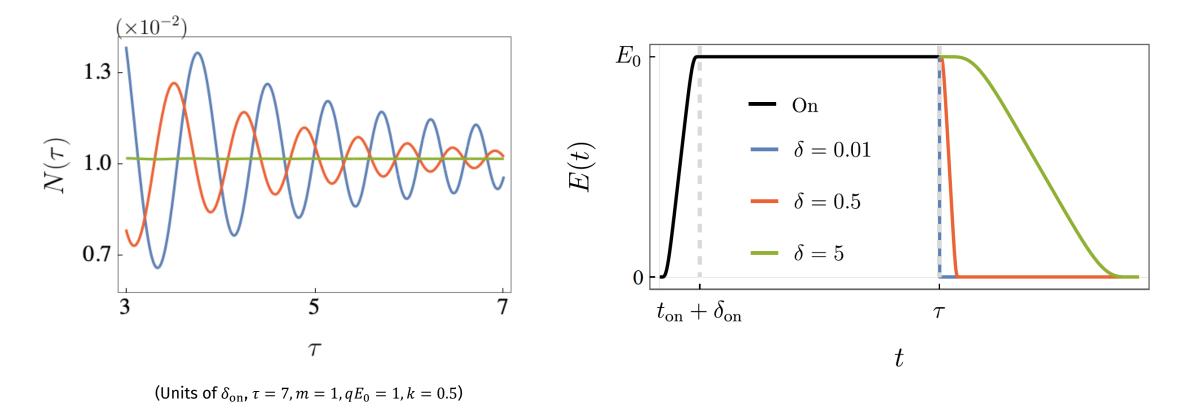






Consider we repeat this for any other $t = \tau$. Given a switch-off profile δ , I have an rule that selects a particular vacuum $|0_{t=\tau}\rangle$ such that $N(\tau)$ has a well-defined physical meaning:

 $N(\tau)$ is the number of particles that would have been measured if we had started the switch-off process at τ



Summary

- Quantum ambiguities arise in QFT in the presence of an external, time-dependent agent
- Particle number notion depends on the choice of vacuum
- Each measurement procedure/switching-off profile δ selects a vacuum $|0_{\tau}\rangle$ for which $N(\tau)$ has a well-defined physical meaning
- From the infinitely many quantizations, there is a family which accomodates information about real outcomes, i.e., they can be understood in terms of a measurement procedure
- Canonical quantum ambiguities are inherently physical: They are intimately related to the different ways of measuring

If you are still interested...

<u>arXiv:2303.07436</u>

