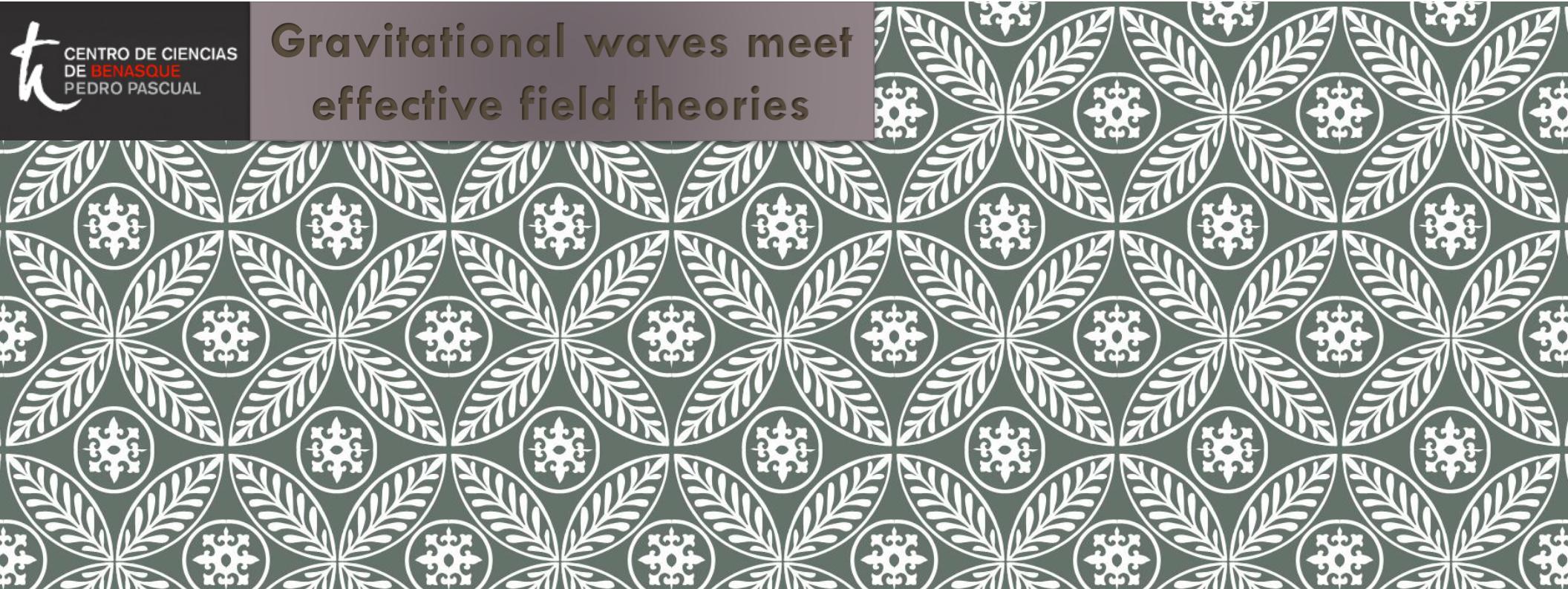
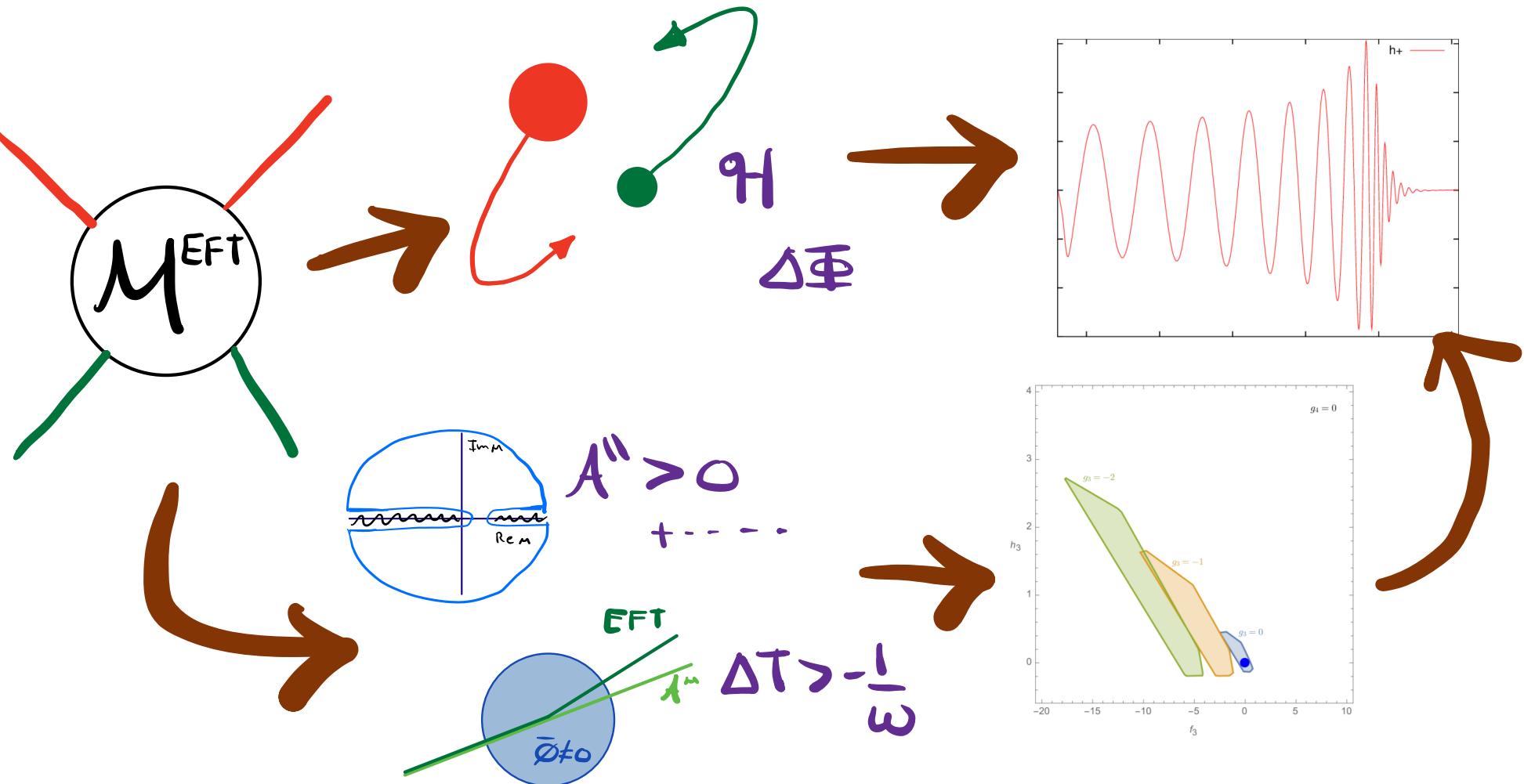


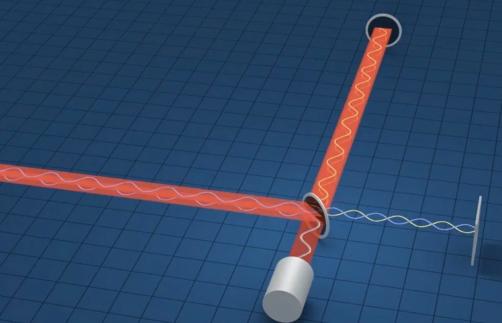
# Gravitational waves meet effective field theories



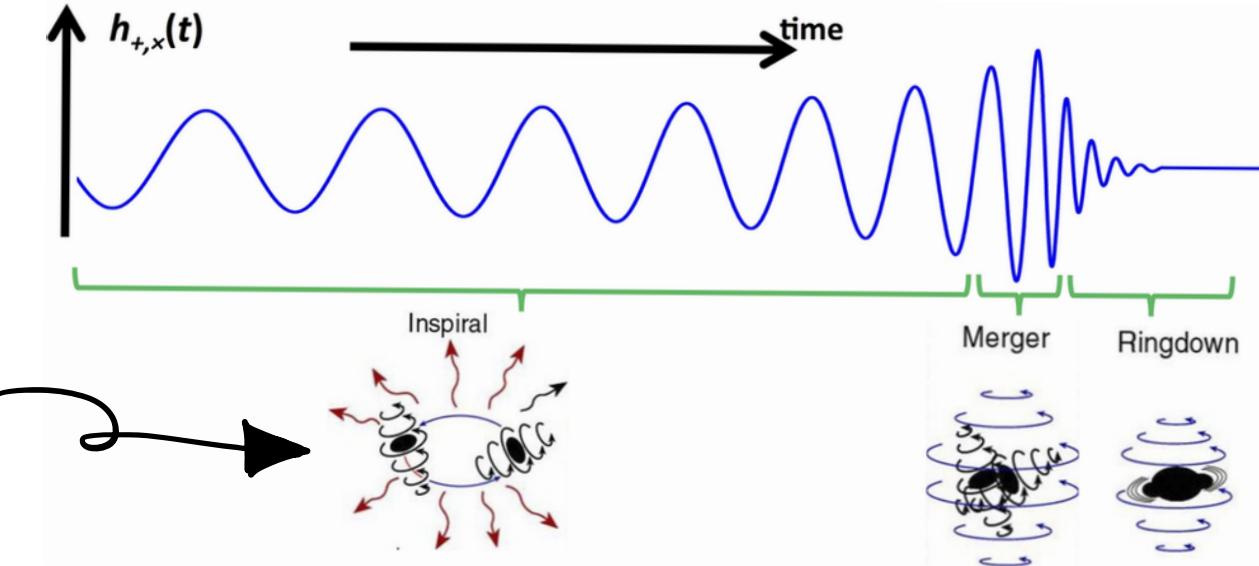
TESTING EFTS: AMPLITUDES,  
GRAVITATIONAL WAVES, AND CAUSALITY

Mariana Carrillo  
González





# MOTIVATION



Perturbative  
gravity

PN

$$\frac{Gm}{r} \sim v^2 \ll 1$$

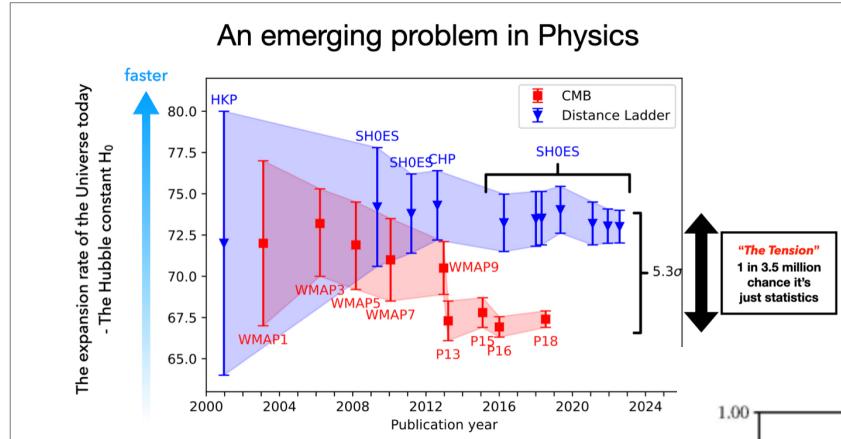
EFT methods

$$\frac{Gm}{r} \ll v^2 \sim 1 + \text{perturbative QFT regime}$$

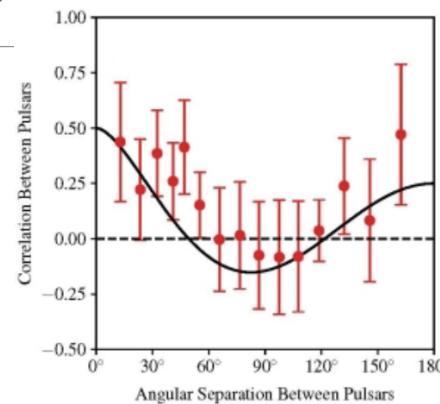
Can  
test  
GR  
extensions

# WHY GO BEYOND GR?

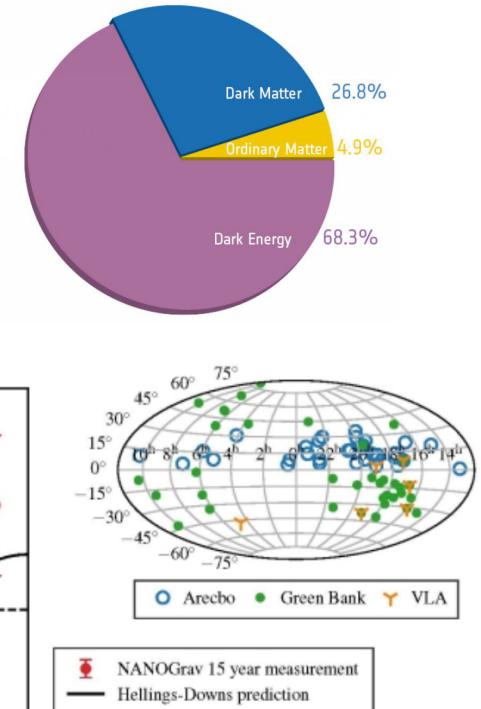
Cosmological  
tensions



Observed stochastic  
gravitational wave background  
of unknown origin



Missing explanation  
for dark matter and  
dark energy



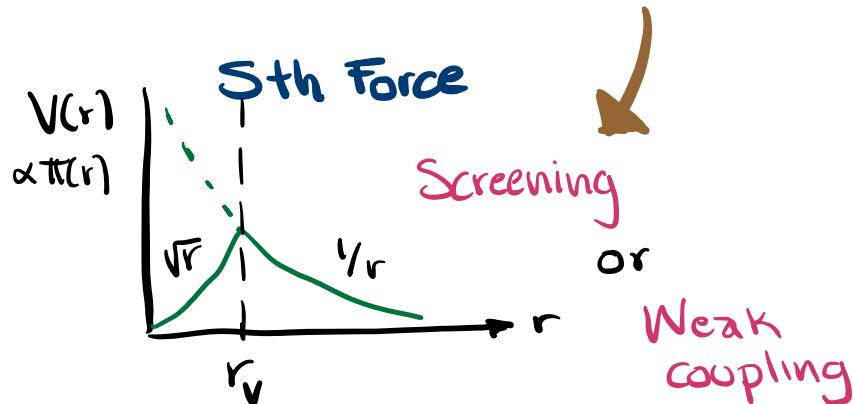
# Beyond GR

$$h_0 \approx 2 \quad \pm 1 \quad 0 \star$$

Generically arise in theories with:

- higher dim.
- higher spin states
- massive gravitons

Strong constraints  
from Solar System tests



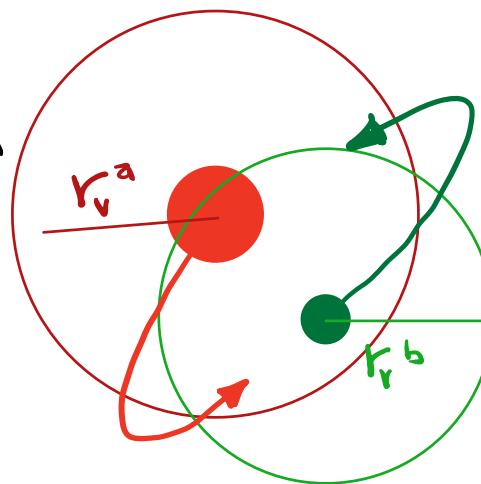
e.g. Cubic Galileon

$$\mathcal{L} = -\frac{1}{2} (\partial \pi)^2 - \frac{C}{\Lambda^3} \square \pi (\partial \pi)^2$$

shift symmetry / soft limit

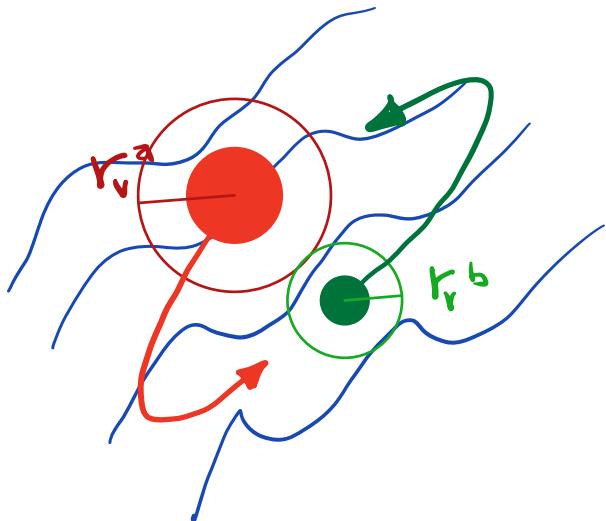
$$\delta \Pi = c + b_\mu x^\mu \Rightarrow A \sim p^2$$

Expectation  
in vacuum

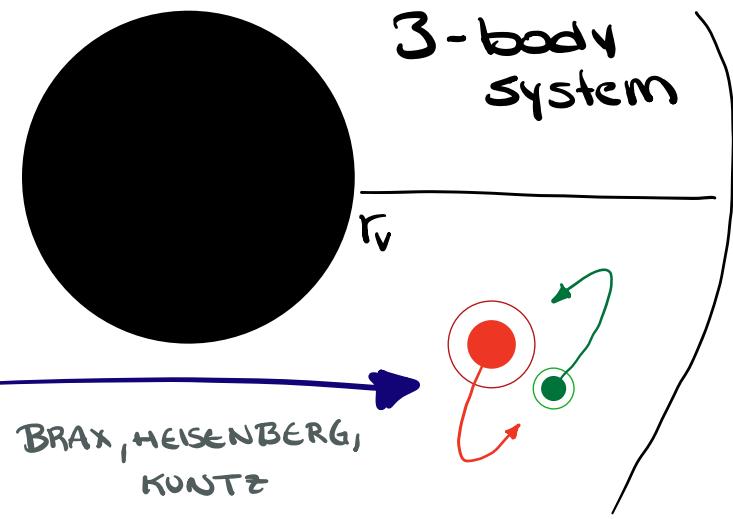


Redressed  $r_v(g)$  in  
a background:

$$\pi = \pi_0 + \delta\pi$$



Perturbative  
regime

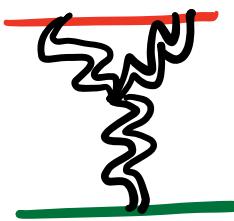


# Classical GR from Amplitudes

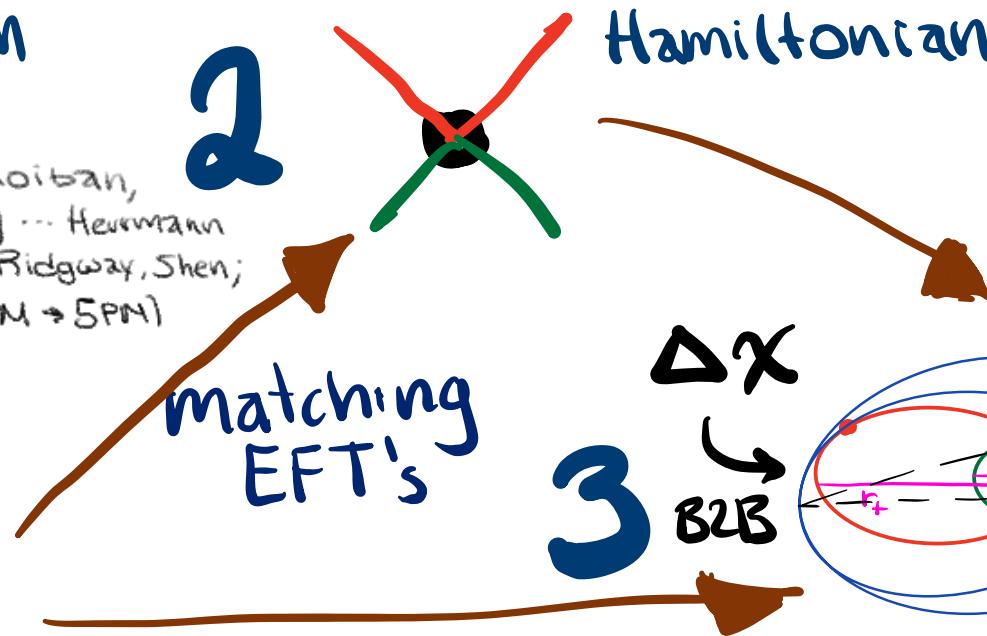
## Post-Minkowskian expansion

Bern, Parra-Martinez, Roiban,  
Ruf, Shen, Solon, Zeng ... Henrann  
Smirnov, Smirnov; Manohar, Ridgway, Shen;  
(4PM  $\rightarrow$  5PM)

1



classical loops



+ spin + radiation ... See SNOWMASS 2204.05194

Other related approaches:

- Worldline (QFT, EFT)

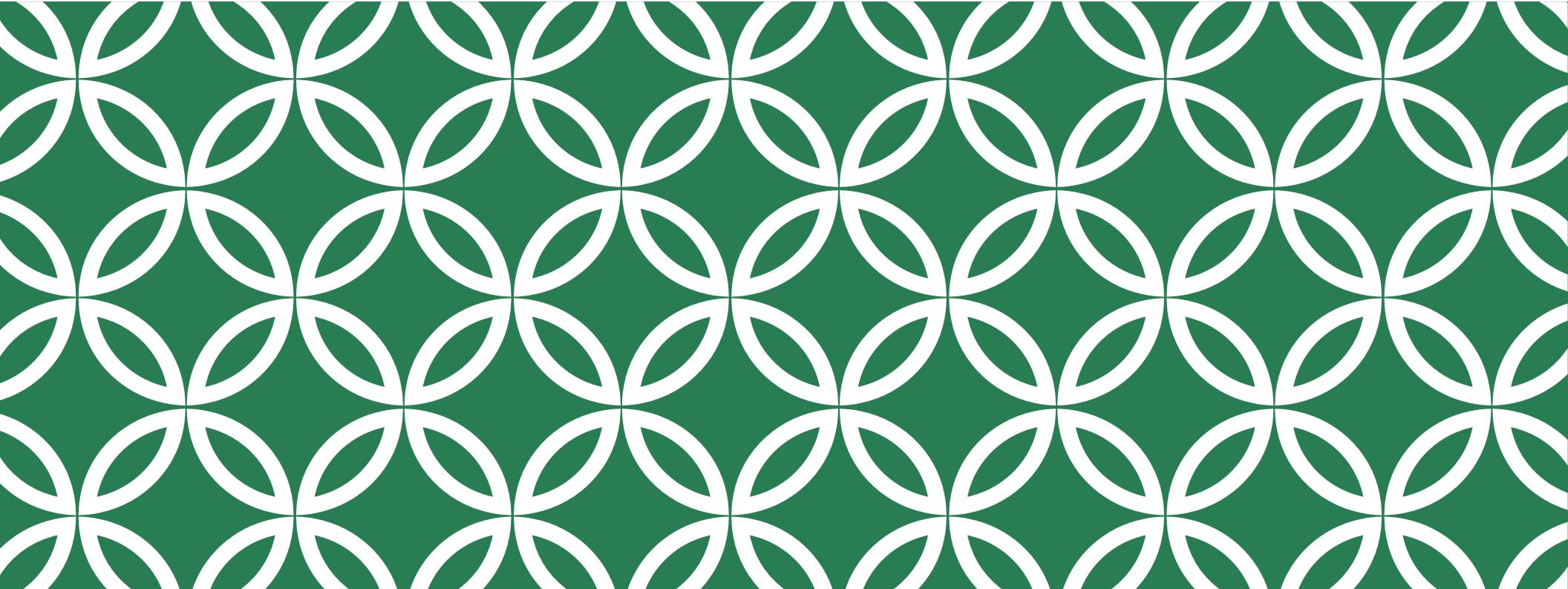
Levi, Yin, Morales, Klem  
Edison; Kalin, Porto, Neel,  
Liu, Datta; Maggi, Plefka;  
Steinhoff, Jakobson;  
Sauer, Xu; Riva,  
Vernizzi; Moussetakos ...

- Eikonal

Di Vecchia, Hessenberg,  
Russo, Vanziano,  
Georgiou, Vazquez-  
Holm; ...

- HEFT

Brandhuber, Travaglini;  
Chen; Ando, Haddad, Hecht  
...

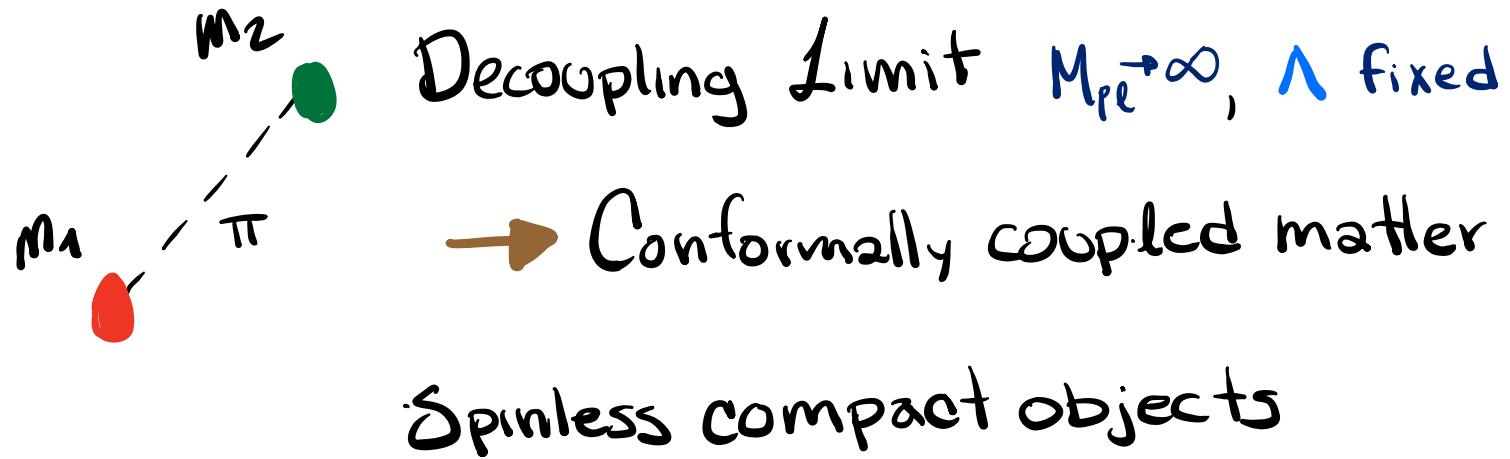


# BINARY SYSTEMS BEYOND GR

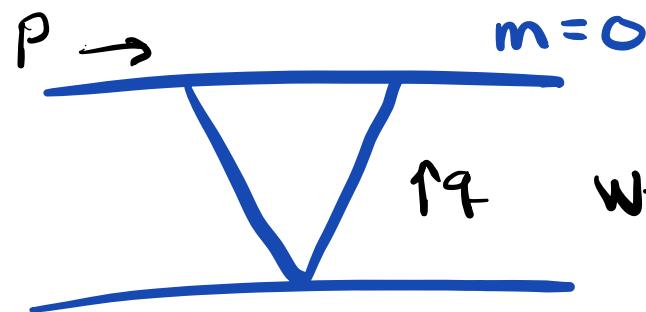
2107.11384:  
MCG, de Rham, Tolley

# Amplitude methods beyond minimal couplings and GR

Test simple scenario: Cubic Galileon



## CLASSICAL REGIME



$$W = \frac{i\hbar}{e} \log Z$$

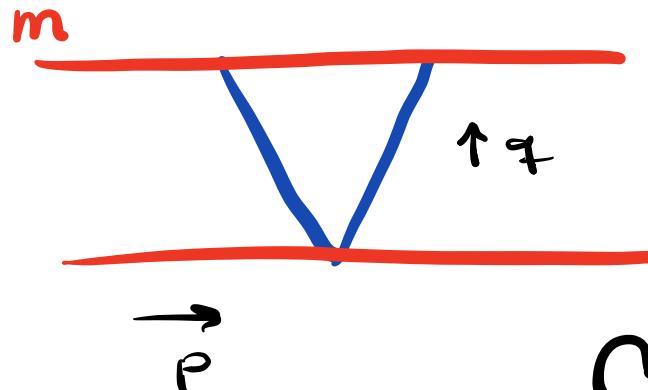
"Usual counting"

$$\hbar^{I-V+1} = \hbar^L$$

loop  
= quantum

$$I: \frac{i\hbar}{q^2 - ie}$$

$$V: e^{\frac{i}{\hbar} \oint L_{int}}$$



$$I_m: \frac{i\hbar}{q^2 + \frac{m^2}{k^2} - ie}$$

$$\hbar \left( \sqrt{\frac{m^2}{q^2 k^2}} \right) \sim \hbar^0$$

Classical contribution from loops!

# 1 Identify Classical Contributions

Restore  $\hbar$ :  $G \rightarrow \frac{G}{\hbar}$   $q \rightarrow \hbar q$

Classical non-linearities

$$\epsilon^{NL} = \frac{\hbar}{M_{Pl}} \propto G m q \quad \hbar \ll 1$$



Quantum corrections

$$\epsilon^Q = \frac{\partial^2}{M_{Pl}^2} \propto q^2 G \quad \hbar \ll \epsilon^{NL}$$

$$\epsilon^{Qm} = \frac{a}{m} \propto \frac{q}{m} \quad \hbar \ll \epsilon^{NL}$$

Post-Minkowskian expansion

Classical limit

# Identify Classical Contributions

wave number  $\rightarrow q \sim 1/b \leftarrow$  impact parameter

Classical non-linearities

$$\epsilon^{NL} = \frac{h}{M_{pl}} \propto \frac{r_s}{b} \ll 1$$

Quantum corrections

$$\epsilon^Q / \epsilon^{NL} \propto \frac{\alpha_c}{b} \ll 1$$

$$\epsilon^{Qm} / \epsilon^{NL} \propto \frac{\alpha_c}{r_s} \ll 1$$

Post-Minkowskian expansion

Classical limit

## Classical non-linearities

$$\epsilon^{NL} = r_s q \hbar^0 \ll 1$$



Post-Minkowskian expansion



$$\epsilon^{NL} = \frac{\partial \partial \pi}{\lambda^3} = r_v q \hbar^0 \ll 1$$

$$\leftarrow = \left( \frac{g m}{m_{pe}} \right)^{1/3} \frac{c}{\lambda}$$

## Quantum corrections

$$\epsilon^Q = q^2 G \hbar \ll \epsilon^{NL}$$

$$\epsilon^{Qm} = \frac{q}{m} \hbar \ll \epsilon^{NL}$$

Classical limit

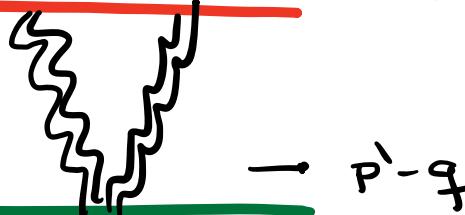


$$\epsilon^Q = \frac{\partial^2}{\lambda^2} = \frac{q^2}{\lambda^2} \hbar^{1/3} \ll \epsilon^{NL}$$

$$\rightarrow \lambda c / r_v \ll 1$$

# Classical physics from loops

2PM

$$\begin{array}{ccc} p & \xrightarrow{\quad} & p+q \\ \text{---} & & \text{---} \\ & \text{---} & \text{---} \\ p' & \xrightarrow{\quad} & p'-q \end{array}$$


$$A \sim \frac{G m_1^2 m_2^2}{q^2} (G m_1 q)$$

↑ Non-analytic

Generic EFT  
classical expansion

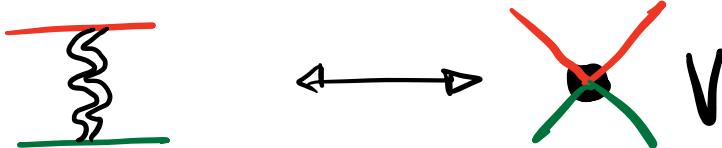
$$A \sim \frac{G m_1^2 m_2^2}{q^2} (r_s q)^n (r_v q)^{3m}$$

GR / Matter coupling / Galileon

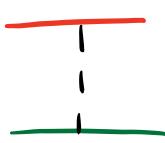
# 2 Matching EFTs

GR: Cheung, Rothstein, Solon

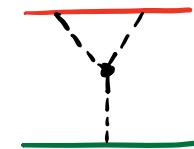
$\mathcal{O}(G)$



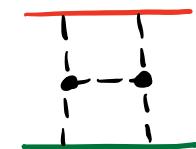
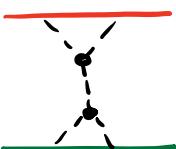
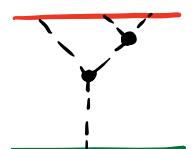
$\mathcal{O}(r_v^0)$



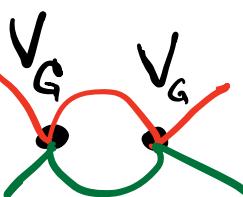
$\mathcal{O}(r_v^3)$



$\mathcal{O}(r_v^4)$



$\mathcal{O}(G^2)$



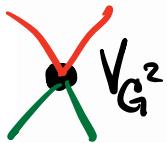
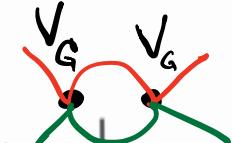
.....

$$\frac{\mathcal{M}^{\text{clas}}}{4E_1 E_2} = \mathcal{M}^{\text{EFT}}$$

$\mathcal{L}^{\text{EFT}} \rightarrow \sqrt{E_p E_{p'} E_{-p} E_{-p'}}$

\* Consider canonically normalized states for matching

$$V(\mathbf{p}, \mathbf{p} + \mathbf{q})|_{\mathcal{O}(g^n)} = -\frac{1}{4E_1(\mathbf{p})E_2(\mathbf{p})}\mathcal{M}_4(\mathbf{p}, \mathbf{p} + \mathbf{q})\Big|_{\mathcal{O}(g^n)}$$


 $\times$ 

 $\cdots$

$$-\frac{1}{4E_1(\mathbf{p})E_2(\mathbf{p})} \int_{\mathbf{k}} \frac{\mathcal{M}_4(\mathbf{p}, \mathbf{k})\mathcal{M}_4(\mathbf{k}, \mathbf{p} + \mathbf{q})}{4E_1(k)E_2(k)(E_p - E_k + i\epsilon)}\Big|_{\mathcal{O}(g^n)}$$

**CLASSICAL POTENTIAL:**  $V(\bar{\mathbf{p}}, \bar{\mathbf{r}}) = \cancel{k^3} \int \frac{d^3 \bar{\mathbf{q}}}{(2\pi)^3} e^{-i\bar{\mathbf{q}} \cdot \bar{\mathbf{r}}} \frac{V(\bar{\mathbf{p}}, \bar{\mathbf{p}} + \bar{\mathbf{q}})}{\cancel{k^3}}$

Same as Lippmann-Schwinger Eq. approach  
 Cristofoli, Bjerrum-Bohr, Damgaard, Vanhove

# 3 Scattering Angle

Conserved energy  $\rightarrow \mathcal{H} = E \rightarrow |P|^2 = |P_\infty|^2 - V_{\text{eff}}(E, r)$

Implicit function theorem  $= |P_\infty|^2 - \frac{1}{2E} M_{\text{ce}}(|P_\infty|, r)$

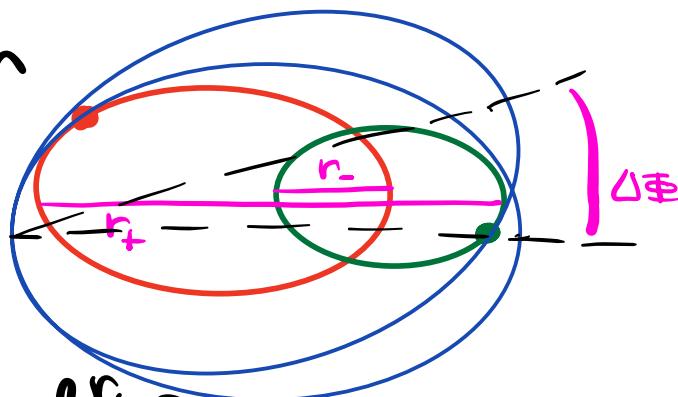
$$\chi = \sum_{k=1}^{\infty} \frac{2b}{k!} \int_0^\infty du \left(\frac{d}{db^2}\right)^k \frac{M_{\text{cl.}}^k (U^2 + b^2)^{k-1}}{(2E)^k |P_\infty|^{2k}}$$

← for any EFT

Bjerrum-Bohr, Cristofoli, Damgaardj  
Kälin, Porto j, Damour j ---

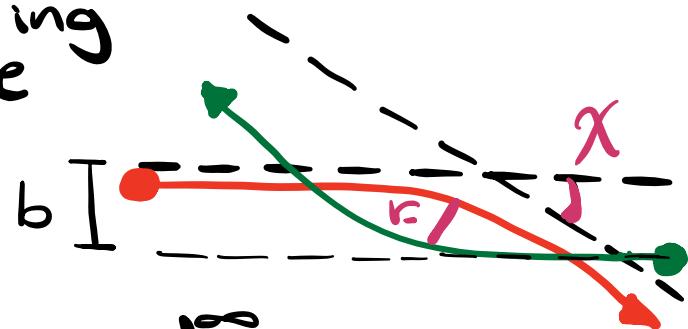
# BINARY SYSTEMS AND SCATTERINGS: BOUNDED VS UNBOUNDED

Perihastron  
advance



$$\Delta\Phi = -2 \int_{r_-}^{r_+} \frac{\partial p_r}{\partial J} - \pi$$

Scattering  
angle



$$\chi = -2 \int_{r_-}^{\infty} \frac{\partial p_r}{\partial J} - \pi$$

$$\Delta\Phi = \chi(J, E) + \chi(-J, E)$$

\* without spin and for aligned spins

Kälin, Porto  
See also : van de Meent

# New contribution to $\Delta \Phi$

$$\Delta \Phi = C_{Gr_v^3} \left( \frac{gGM^2}{J} \right) \left( \frac{r_v^3 |P_{\text{soft}} M|}{J^3} \right) \underline{\underline{1PM}}$$

$$+ C_G^2 \left( \frac{gGM^2}{J} \right)^2 2PM + \dots$$

+ gravitons contribution  $2PM$

Resummation?  
Access screened  
region.

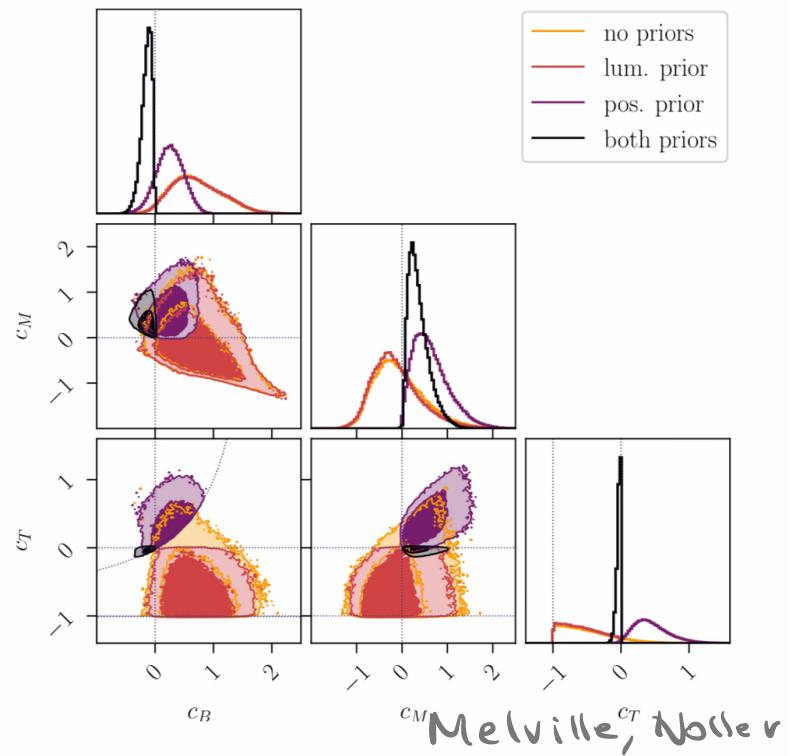
Leading contribution  
depends on the size of

$$g C_{Gr_v^3} \left( \frac{r_v^3 |P_{\text{soft}} M|}{J^3} \right)$$

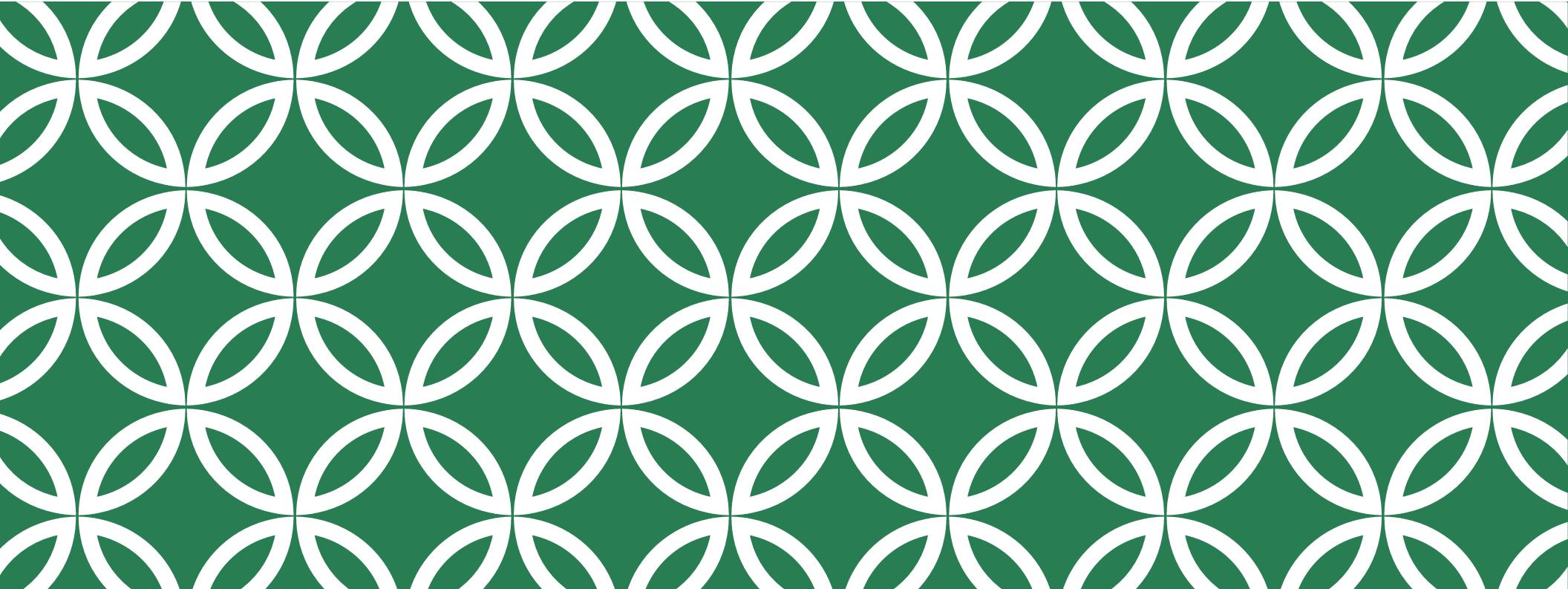
Bounds on  $C_{Gr_v^3}$ ?

# Importance of bounds on Wilson coeffs.

Theoretical priors  
can drastically change  
estimations of parameters



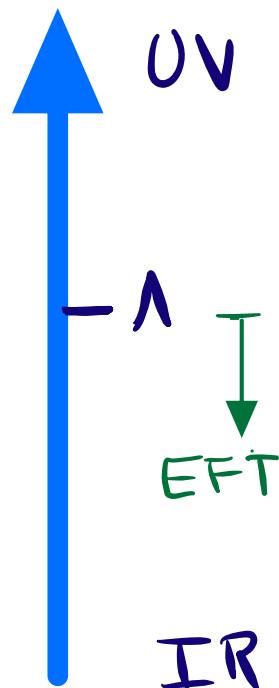
Melville, Noller



# CAUSALITY BOUNDS ON EFTS

---

# CAUSAL EFFECTIVE FIELD THEORIES



$$\mathcal{L} = \Lambda^4 \sum_n c_n \Lambda^{-n} \partial_n$$

What are the allowed  $c_n$ ?

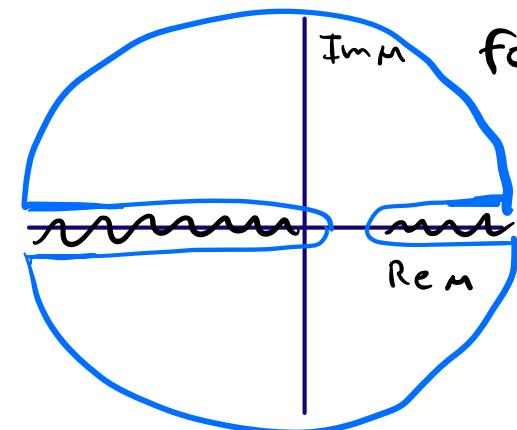
1) UV = string theory  $\rightarrow$  Swampland conjectures

# CAUSAL EFFECTIVE FIELD THEORIES

$$\mathcal{L} = \Lambda^4 \sum_n c_n \Lambda^{-n} O_n$$

What are the allowed  $c_n$ ?

2) UV = local, unitary, causal, Lorentz invariant  $\rightarrow$  Positivity bounds



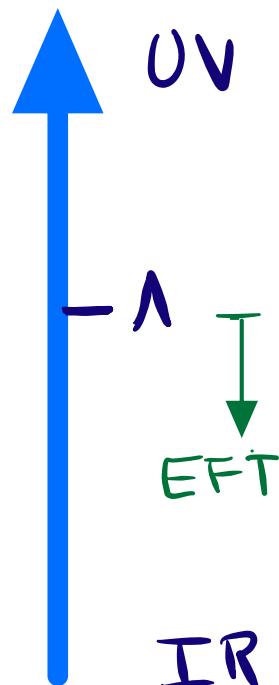
forward lim. ( $+ \rightarrow 0$ )

$$A''(s) = \int \frac{dN}{2\pi i} \frac{A(m)}{(m-s)^3} \stackrel{1,3}{=} \left( \int_{-\infty}^0 + \int_+^\infty \right) \frac{\text{Im } A}{cm^2} \frac{\text{Im } A}{(m-s)^3} > 0$$

UV

related by 4

# CAUSAL EFFECTIVE FIELD THEORIES



$$\mathcal{L} = \Lambda^4 \sum_n c_n \Lambda^{-n} \partial_n$$

What are the allowed  $c_n$ ?

1) UV = string theory  $\rightarrow$  Swampland

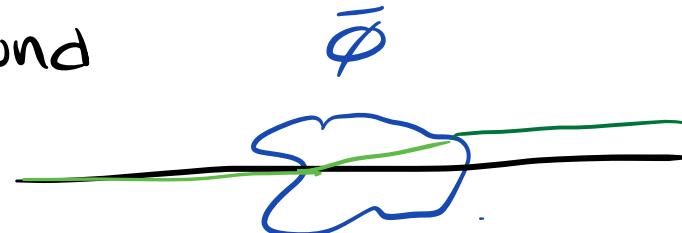
2) UV = local, unitary, causal, Lorentz invariant  
 $\rightarrow$  Positivity bounds

3) Causal IR propagation

**CAUSALITY BOUNDS**

# CAUSALITY

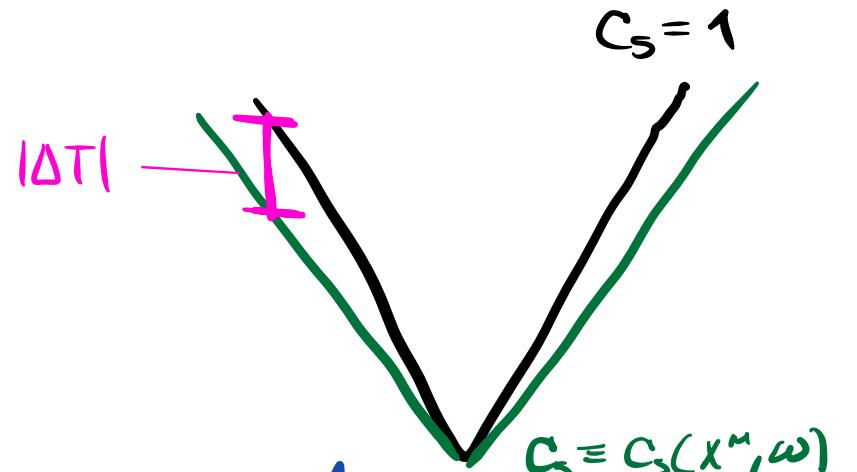
- Consider local propagation of information around a fixed background



Diagnose acausality by looking at time delay

$$\Delta T = -i \langle \text{in} | \hat{S}^+ \frac{\partial}{\partial \omega} \hat{S} | \text{in} \rangle$$

$$G_R(x-y) = 0 \text{ for } (x-y)^2 > 0$$



$$|\Delta T| \lesssim \lambda \sim \frac{1}{\omega}$$

Unresolvable



$$\Delta T \geq -\frac{1}{\omega}$$

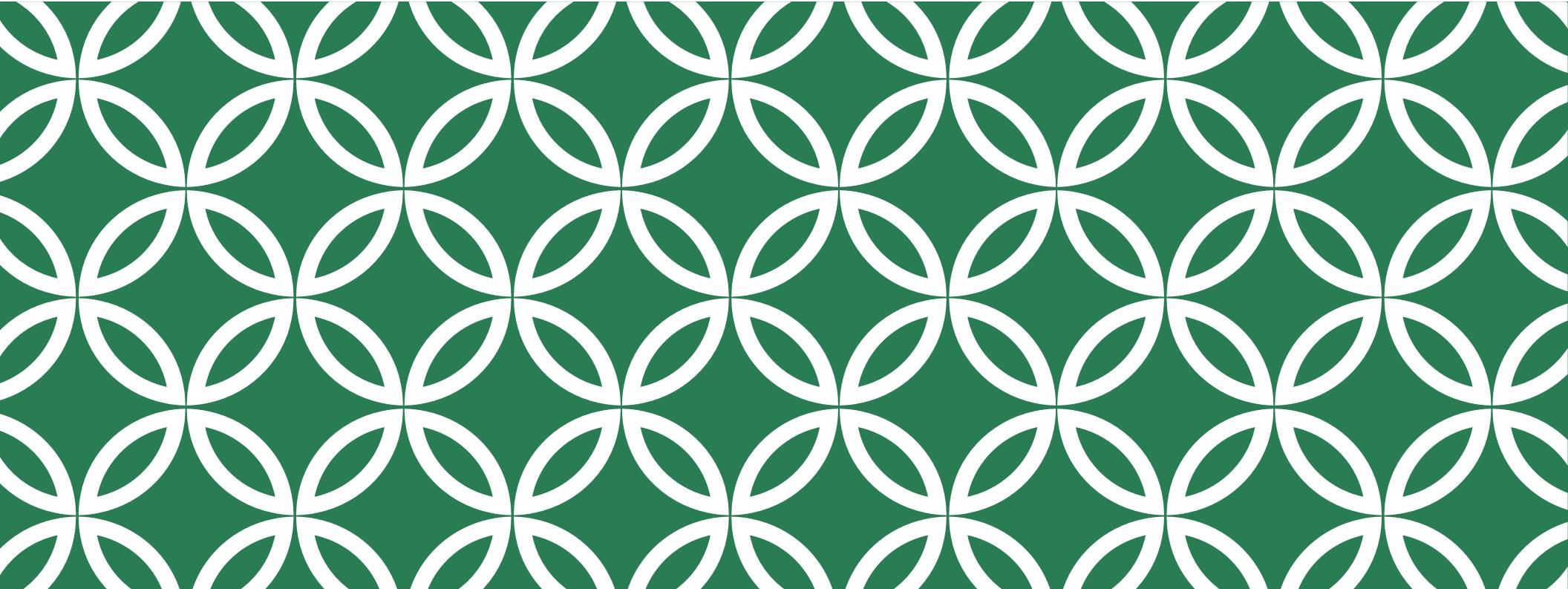
## CAUSALITY + WKB + EFT

Consider  $\phi = \bar{\phi} + \varphi$ . Find  $\Delta T$  experienced by  $\varphi$ .  
Solve linearized  $\varphi$  com using WKB approximation.

$$\frac{\partial}{\partial \lambda_{\text{perturbation}}} \int_{X \in \mathbb{R}^{1+3}} \left( 1 - C_S(\lambda^{\text{pert.}}) \right) \geq -1$$

$\omega \Delta T \geq -1$

$\gg 1$  WKB       $= -\epsilon$  EFT  $\rightarrow |\epsilon| \ll 1$



# CAUSALITY BOUNDS ON SCALAR EFTS

2207.03491:  
MCG, de Rham, Pozsgay,  
Tolley

# CAUSALITY INSIGHTS ON SCALAR EFTS

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 + \frac{g_8}{\Lambda^4}((\partial\phi)^2)^2$$

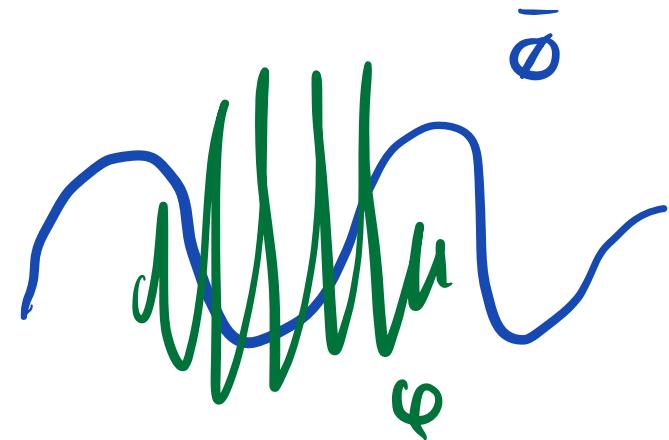
$$+ \frac{g_{10}}{\Lambda^6}(\partial\phi)^2((\partial\partial\phi)^2 - (\square\phi)^2) + \frac{g_{12}}{\Lambda^8}((\partial\partial\phi)^2)^2 - g\phi J$$

↑  
quartic galileon

↑  
matter external source

# CAUSALITY INSIGHTS ON SCALAR EFTS

$$\phi = \bar{\phi} + \varphi \quad \partial_\mu \varphi = i k_m \varphi \quad \text{plane waves}$$



from Disp. rel  $\rightarrow C_S$

$$C_S^2 \simeq \left(1 - \# \frac{g_8}{\Lambda^4} \underbrace{\frac{(K \cdot \partial \bar{\phi})^2}{(|K|^2)}_{> 0}}_{\text{Adams et al}}$$



Work with

$$g_8 = \begin{cases} 1 \\ 0 \end{cases}$$

$$+ \# \frac{g_{10}}{|K|^2 \Lambda^6} \left( (K \cdot \partial \bar{\phi})^2 - \square \phi (K_m K_n \partial^\mu \partial^\nu \phi) \right) - \# \frac{g_{12}}{\Lambda^8} \underbrace{\frac{(K \cdot \partial \bar{\phi})^2}{(|K|^2)^2}_{> 0}}_{\text{Adams et al}}$$

# PROPAGATION AROUND SPHERICALLY SYMMETRIC BACKGROUNDS

Spherically-symmetric background  $\bar{\phi} \equiv \Phi_0 f(r/r_0)$

Perturbations

$$\rightarrow \chi_e^{(1)}(r/r_0) + \underbrace{(\omega r_0)^2}_{\gg 1 \text{ WKB}} \underbrace{\frac{1}{C_s^2(\omega, r)} \left( 1 - \frac{V_e^{\text{eff}}(r)}{(\omega r_0)^2} \right)}_{\equiv W_e(\omega, r)} \chi_e(r/r_0) = 0$$

$\rightarrow \text{Find } \delta \rightarrow \Delta T$

$\ll 1 \text{ EFT}$

WKB :  $\frac{\lambda_{\text{bckgrnd}}}{\lambda_{\text{pert}}} = \omega r_0 \gg 1$

EFT :  $\frac{\partial \phi}{\Lambda}, \frac{\partial^{p+1} \phi}{\Lambda^{p+2}} \ll 1$

$$\frac{\dot{\phi}_0}{r_0 \Lambda}, \frac{1}{r_0 \Lambda} \sim \mathcal{O}(\epsilon) \quad \frac{\omega}{r_0 \Lambda^2} \sim \mathcal{O}(\epsilon \frac{\omega}{\Lambda})$$

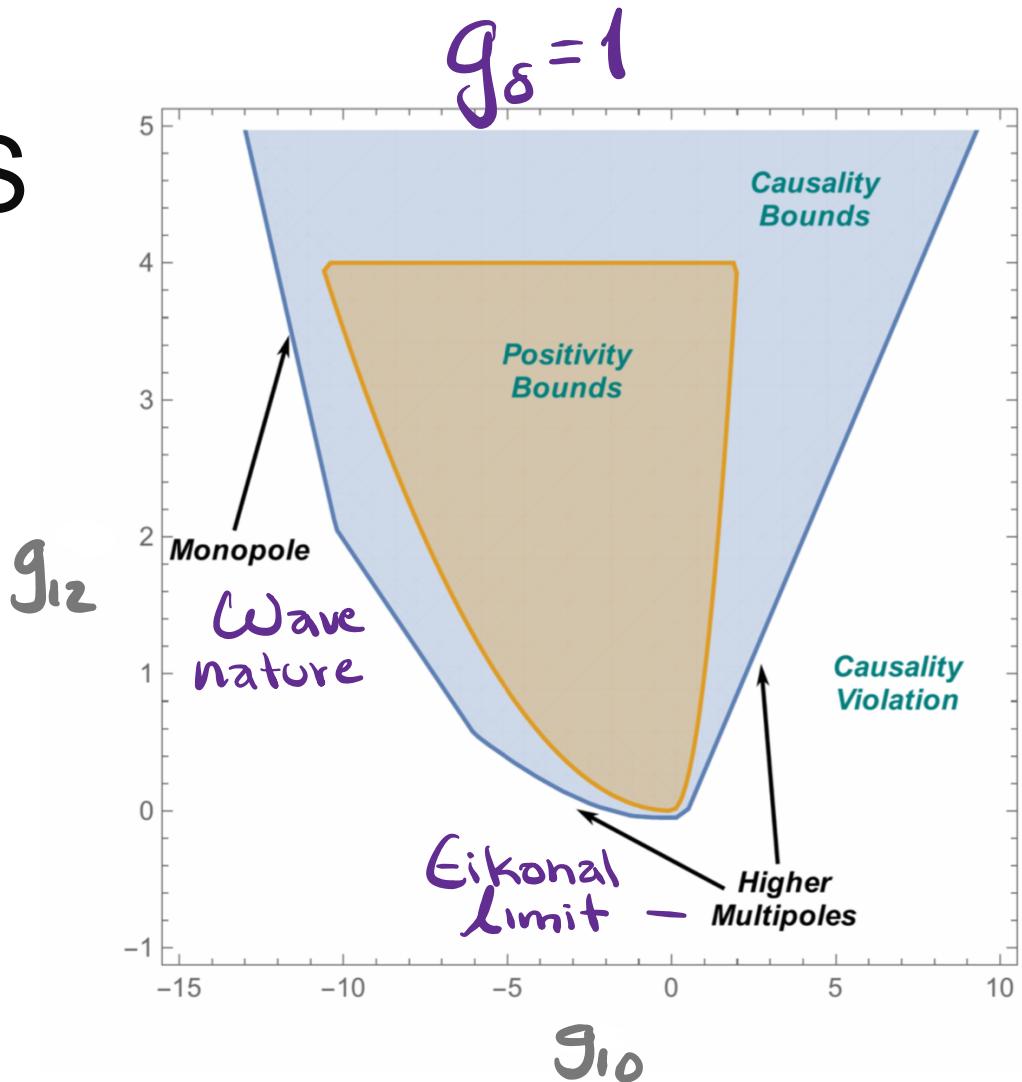
Note bound  
on  $\omega$   $\rightarrow \epsilon < \frac{\omega}{\Lambda} < \frac{1}{\epsilon}$

WKB      EFT

# CAUSALITY VS POSITIVITY



No upper bound on  $g_{12}$  from causality due to WKB technical issues.



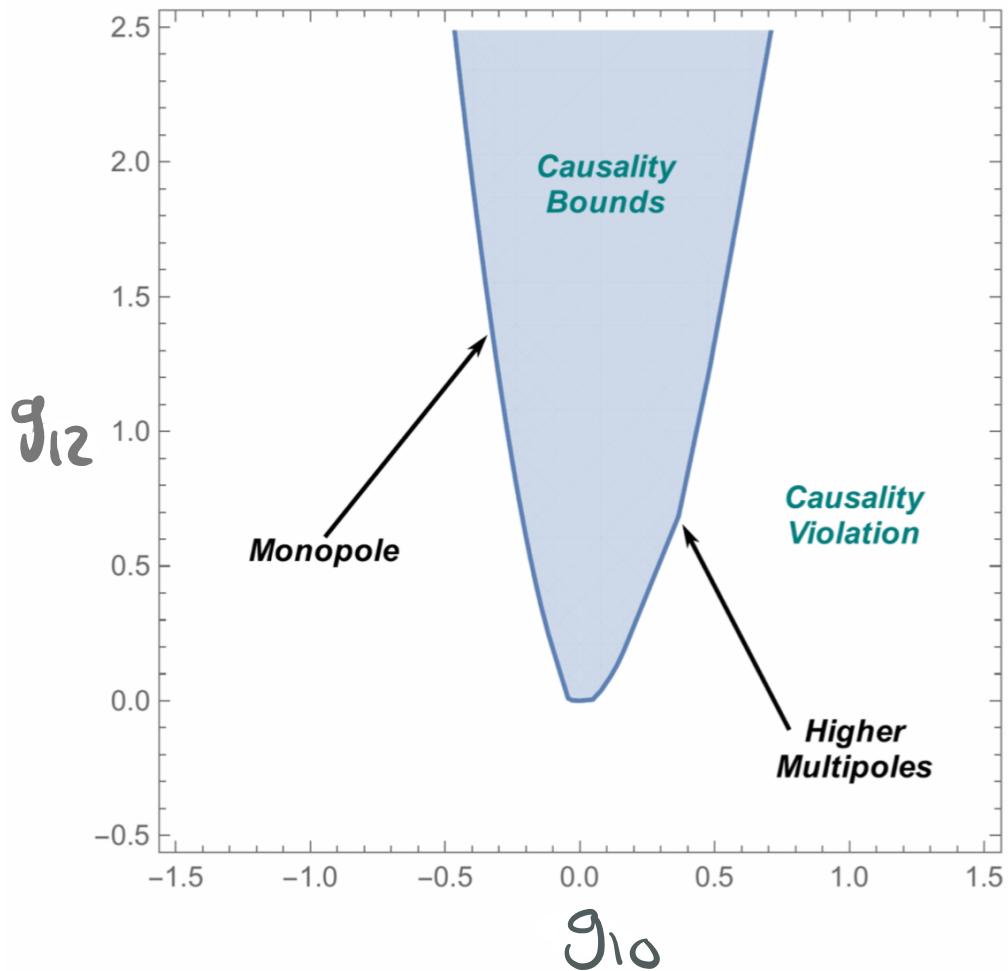
# ADDITIONAL SYMMETRIES: GALILEONS

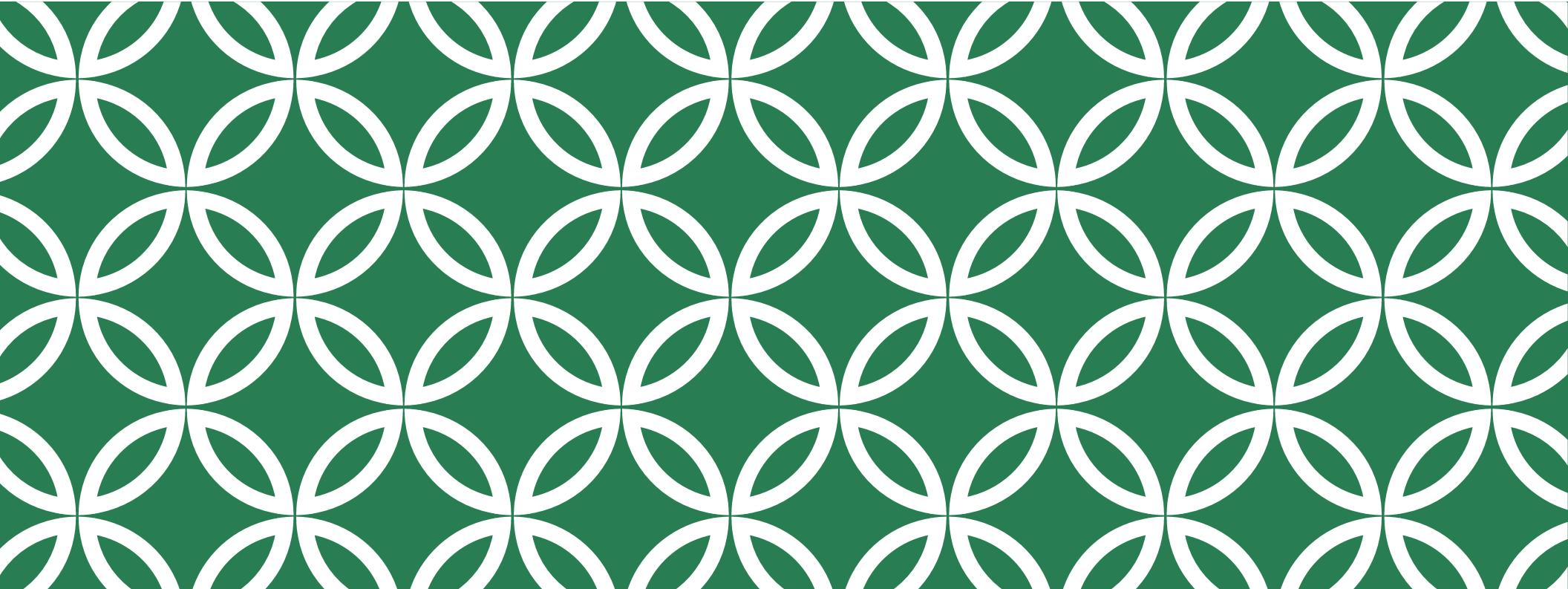
From positivity:

$$g_8 = 0 \Rightarrow g_{10} = g_{12} = 0$$

valid around  $\langle \phi \rangle = 0$

$$g_8 = 0$$





# CAUSALITY AND POSITIVITY BOUNDS ON PHOTON EFTS

2307.04784

MCG, de Rham, Jaitly,  
Pozsgay, Tokareva

# PHOTON EFT

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \\
 & + \frac{c_1}{\Lambda^4}F^{\mu\nu}F_{\mu\nu}F^{\alpha\beta}F_{\alpha\beta} + \frac{c_2}{\Lambda^4}F^{\mu\nu}F^{\alpha\beta}F_{\mu\alpha}F_{\nu\beta} \\
 & + \frac{c_3}{\Lambda^6}F^{\alpha\mu}F^{\nu\beta}\partial_\mu F_{\beta\gamma}\partial_\nu F_\alpha{}^\gamma + \frac{c_4}{\Lambda^6}F^{\alpha\mu}F^{\nu\beta}\partial_\beta F_{\mu\gamma}\partial^\gamma F_{\alpha\nu} + \frac{c_5}{\Lambda^6}F^{\alpha\mu}F^{\nu\beta}\partial_\beta F_{\nu\gamma}\partial^\gamma F_{\alpha\mu} \\
 & + \frac{c_6}{\Lambda^8}F^{\mu\nu}\partial_\mu F_{\nu\rho}\partial^\rho\partial^\alpha F^{\beta\gamma}\partial_\alpha F_{\beta\gamma} + \frac{c_7}{\Lambda^8}F^\mu{}_\gamma\partial_\mu F_{\nu\rho}\partial^\nu F_{\alpha\beta}\partial^\rho\partial^\gamma F^{\alpha\beta} \\
 & + \frac{c_8}{\Lambda^8}F^{\mu\gamma}\partial_\mu F_{\nu\rho}\partial^\rho\partial^\beta F_{\alpha\gamma}\partial^\alpha F^\nu{}_\beta.
 \end{aligned}$$

Dim. 8

$$F^4 : \underline{c_1, c_2}$$

Dim. 10

$$\partial^2 F^4 : \underline{c_3, c_4, c_5}$$

Dim. 12

$$\partial^4 F^4 : \underline{c_6, c_7, c_8}$$

# PHOTON EFT

$$\begin{aligned}\mathcal{A}^{++++} &= \underline{f_2} (s^2 + t^2 + u^2) + \underline{f_3 stu} + \underline{f_4} (s^2 + t^2 + u^2)^2 + \mathcal{O}(s^5) \\ \mathcal{A}^{++--} &= \underline{g_2 s^2} + \underline{g_3 s^3} + \underline{g_4 s^4} + \underline{g'_4 s^2 tu} \\ \mathcal{A}^{+++-} &= \underline{h_3 stu}\end{aligned}$$

Dim. 8

$$F^4 : f_2, g_2 \quad \downarrow = 1,0$$

Dim. 10

$$\partial^2 F^4 : \underline{f_3, g_3, h_3}$$

Dim 12

$$\partial^4 F^4 : \underline{f_4, g_4, g'_4}$$

# PROPAGATION AROUND SPHERICALLY SYMMETRIC BACKGROUNDS

$$A = A + \delta A, \quad \bar{A} = \Xi_0 f(r/r_0) dt +$$

$$\chi_e^{\text{even}} + (\omega r_0)^2 W_e^{\text{even}}(\omega, r) \chi_e^{\text{even}} = 0$$

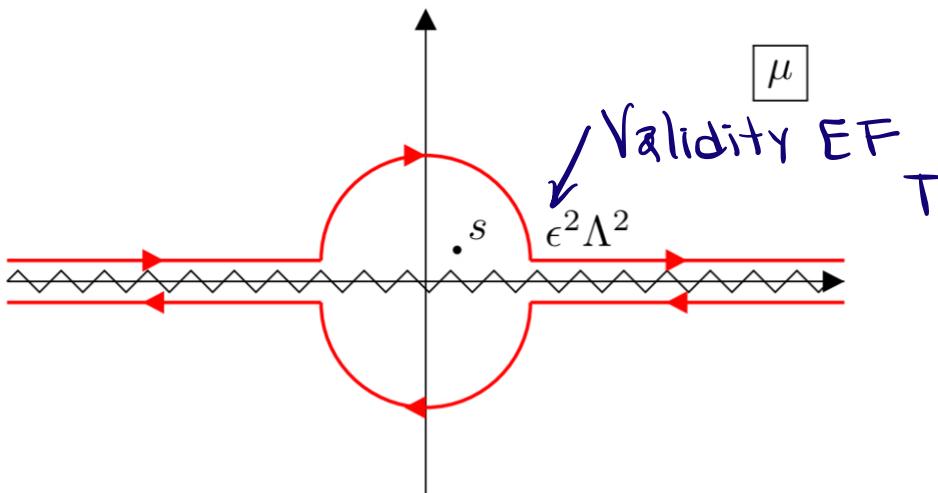
$$\Delta T^{\text{even}} > -\frac{1}{\omega}$$

$$\chi_e^{\text{odd}} + (\omega r_0)^2 W_e^{\text{odd}}(\omega, r) \chi_e^{\text{odd}} = 0$$



$$\Delta T^{\text{odd}} > -\frac{1}{\omega}$$

# Positivity Bounds



Other explorations:  
 Henriksson, McPeak, Russo,  
 Vichi; Häring, Hebbbar, Kavaler,  
 Heinri, Penedones

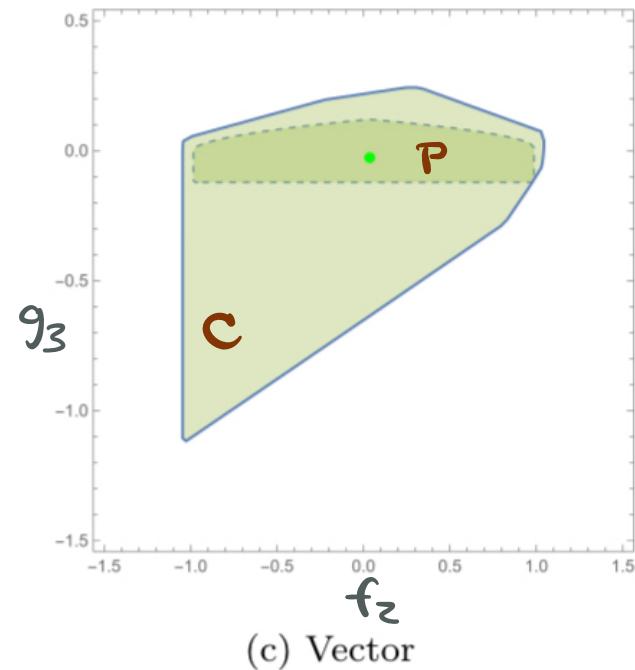
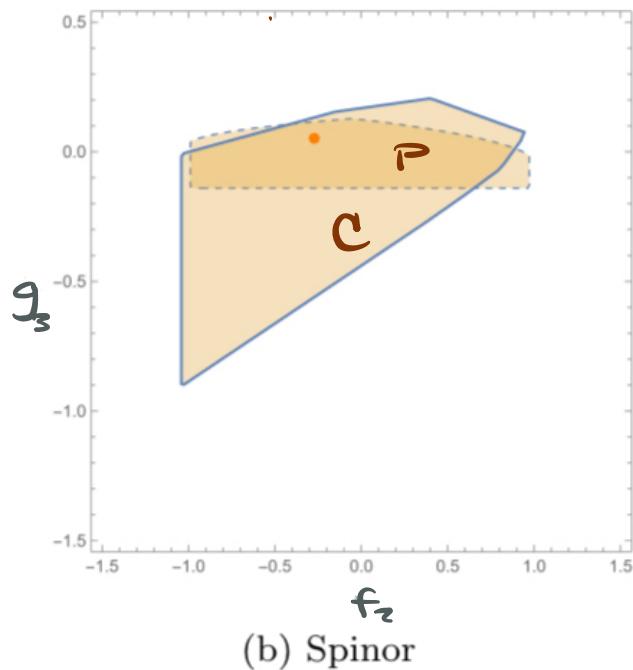
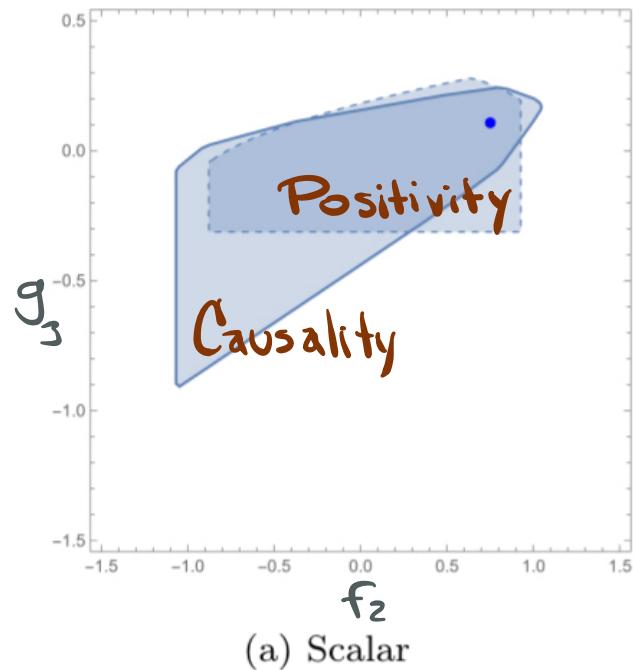
$$\frac{1}{2\pi i} \int_{\delta} \frac{\mathcal{A}(\mu, t)}{(\mu - s)^3} d\mu = \int_{\epsilon^2 \Lambda^2}^{\infty} \frac{d\mu}{\pi} \frac{\text{Disc}_s \mathcal{A}_s(\mu, t)}{(\mu - s)^3} + \int_{\epsilon^2 \Lambda^2 - t}^{\infty} \frac{d\mu}{\pi} \frac{\text{Disc}_s \mathcal{A}_u(\mu, t)}{(\mu - u)^3}$$

Simple example: forward limit ( $t \rightarrow 0$ )

$$\frac{1}{2\pi i} \int_{\delta} \frac{\mathcal{A}(\mu, 0)}{\mu^3} d\mu > 0 \quad \text{For } \mathcal{A}(\mu) = g_2 \mu^2 + \alpha \mu^4 \log(\mu)$$

$$\Rightarrow g_2 + \frac{\alpha}{2} \epsilon^4 > 0$$

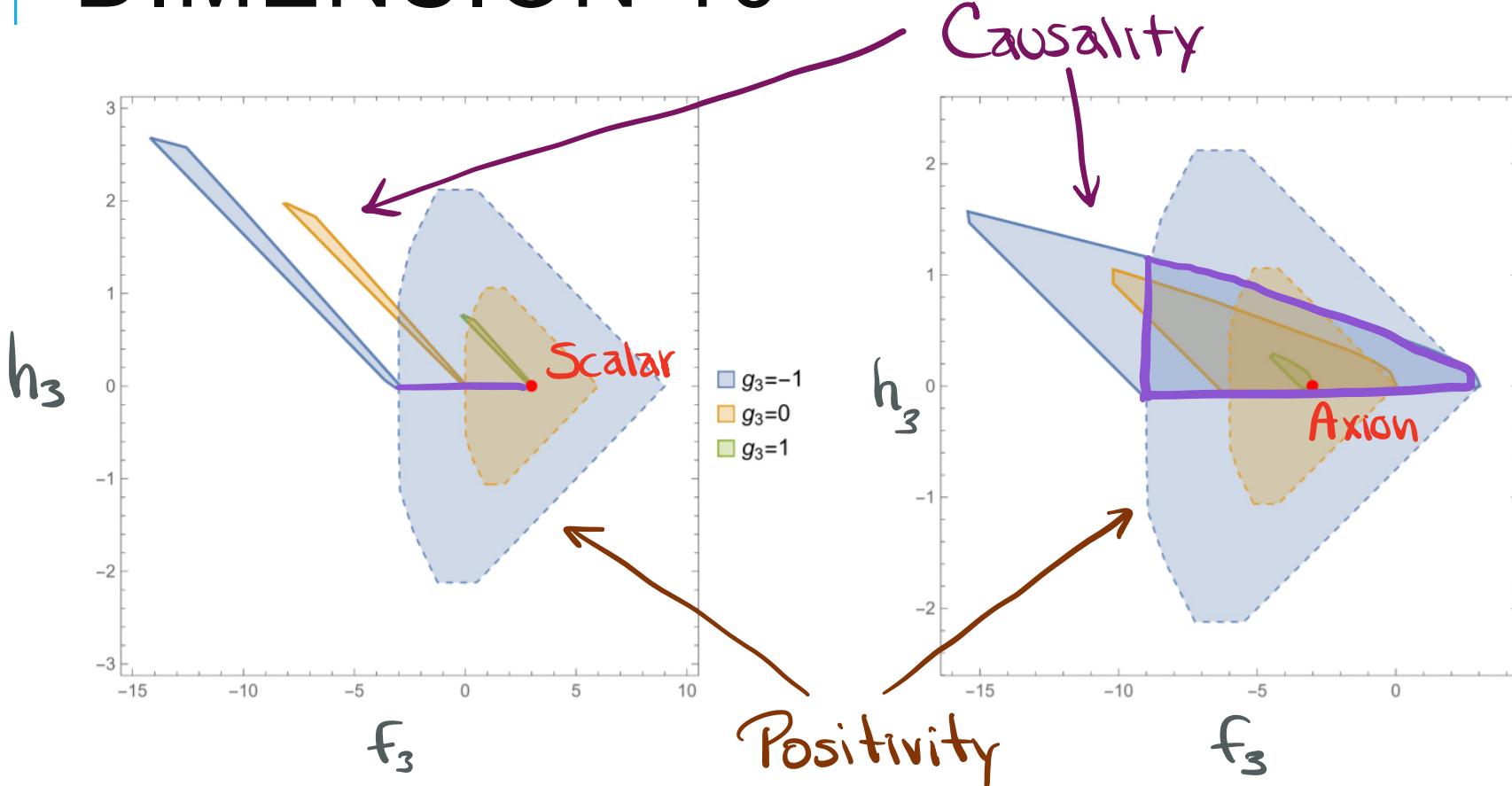
# QED-like partial UV completions

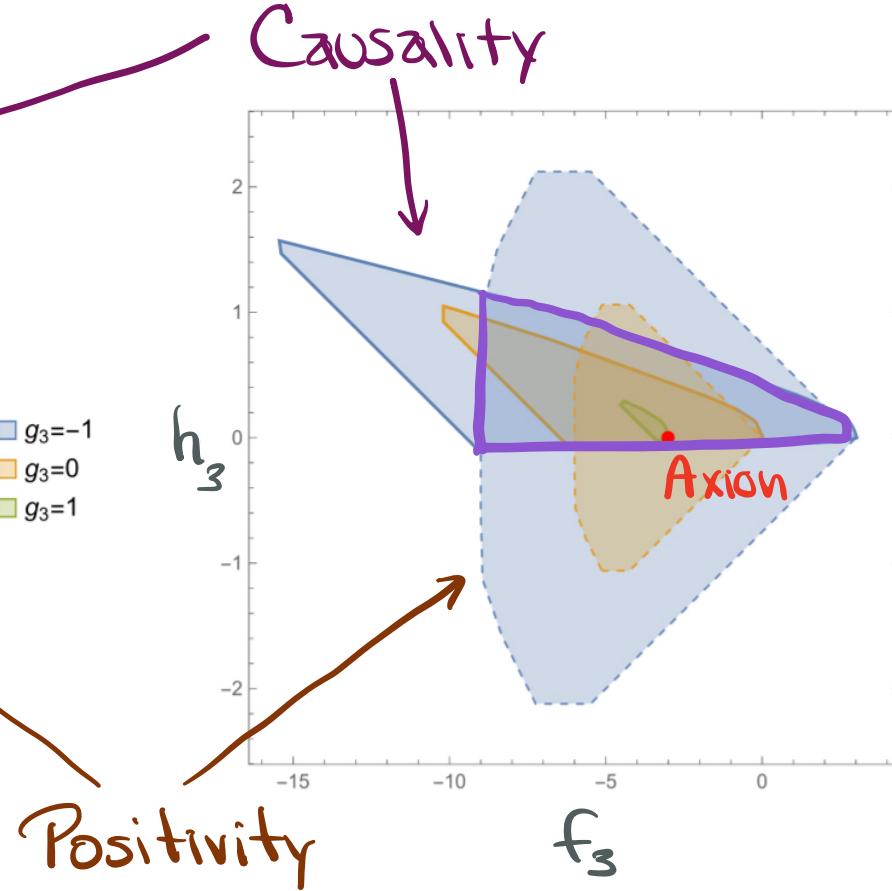
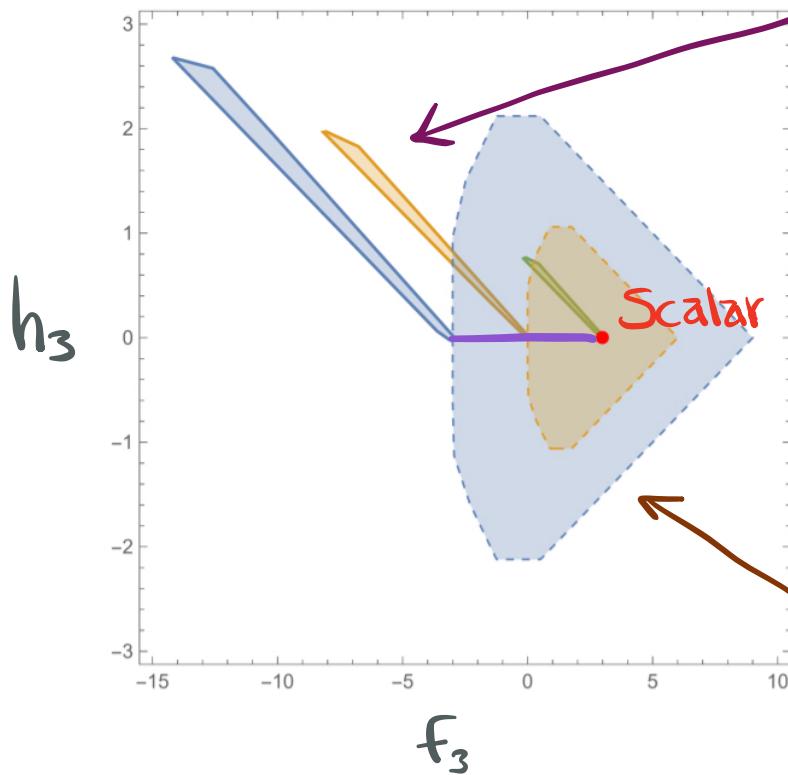


# DIMENSION 10

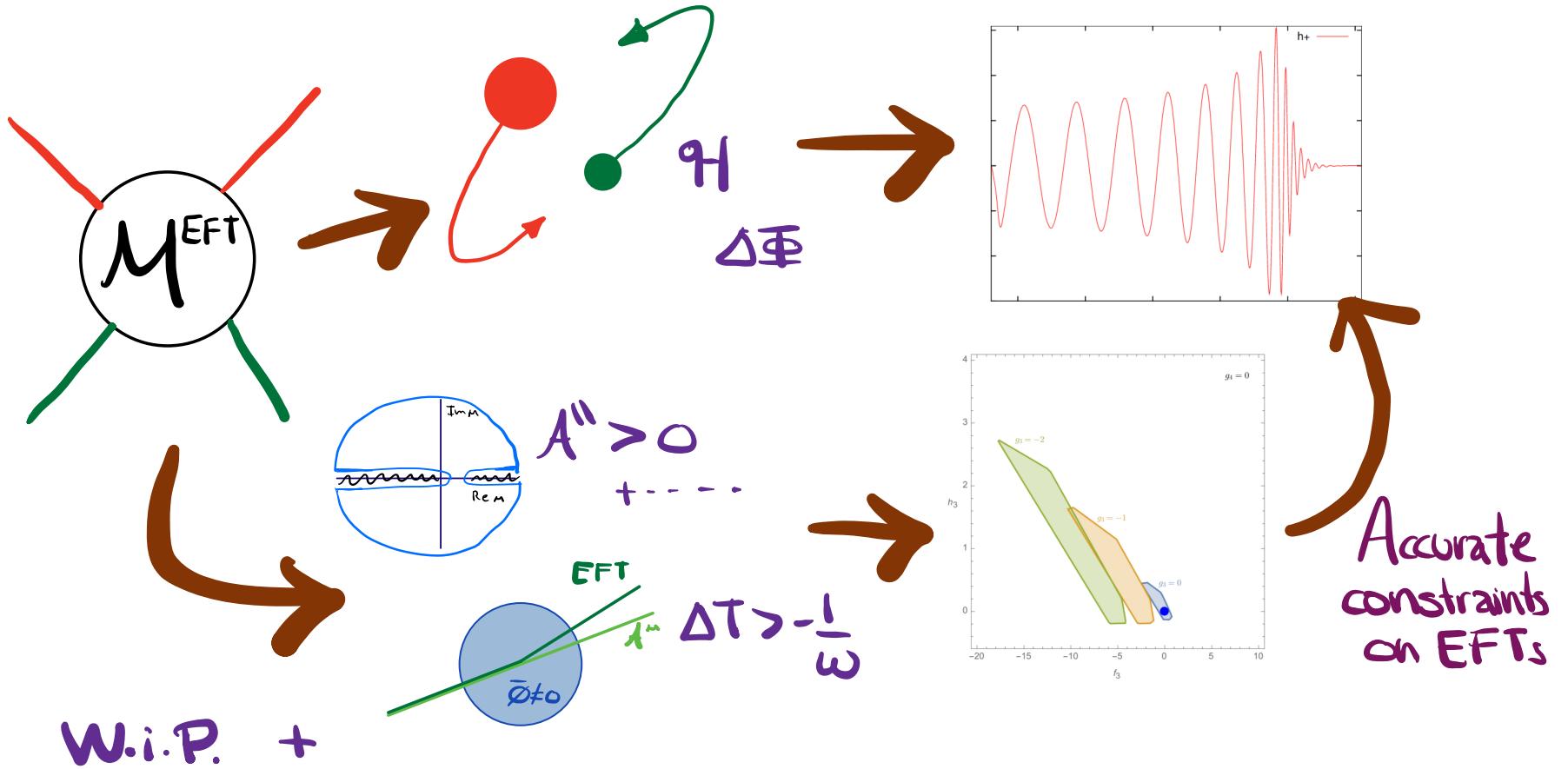
$$(f_3 + 3g_3) < X^{\text{even}} + \epsilon h_3$$

$$-X^{\text{odd}} - \epsilon h_3 < f_3 - 3g_3 + 4h_3 < X^{\text{odd}} + \epsilon h_3$$





Positivity and Causality test different regions of the parameter space, NOT in conflict w/ each other



- Future directions
- Test higher dimensional gravitational operators
  - Cosmological backgrounds ✓ + EFT of inflation