

EXTREME VALUE THEORY AND LOCALIZATION IN RANDOM SPIN CHAINS



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arXiv:2305.10574

Entanglement in strongly correlated systems
Benasque, 2023



EXTREME VALUE THEORY



Extreme floods



Market risks



Athletic records



Large wildfires



E. J. Gumbel, *Statistics of Extremes*, Dover, (1958, 2004)

S. N. Majumdar, A. Pal, G. Schehr, *Physics Reports*, **840**, 1 (2020)

EXTREME VALUE THEORY



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Condensed matter



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EXTREME VALUE THEORY



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Disordered spin chains

R. Juhász, Y.C. Lin, and F. Iglói, “S Phys. Rev. B 73, 224206 (2006)

N. Pancotti, M. Knap, D. A. Huse, J. I. Cirac, and M. C. Bañuls, Phys. Rev. B 97, 094206 (2018)

I. A. Kovács, T. Pető, and F. Iglói, Phys. Rev. Res. 3, 033140 (2021)

W.-H. Kao and N. B. Perkins, Phys. Rev. B 106, L100402 (2022)

J. C., N. Laflorencie, arXiv:2305.10574



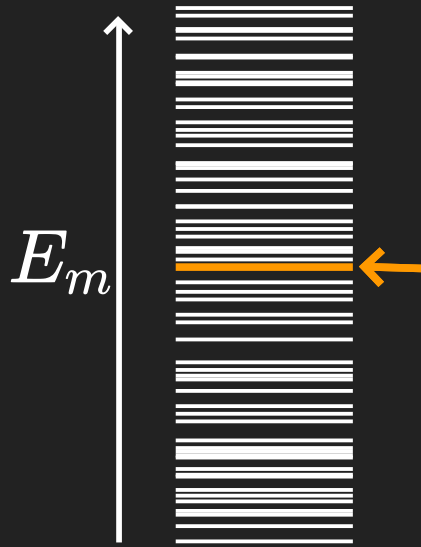
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"TRAILER"



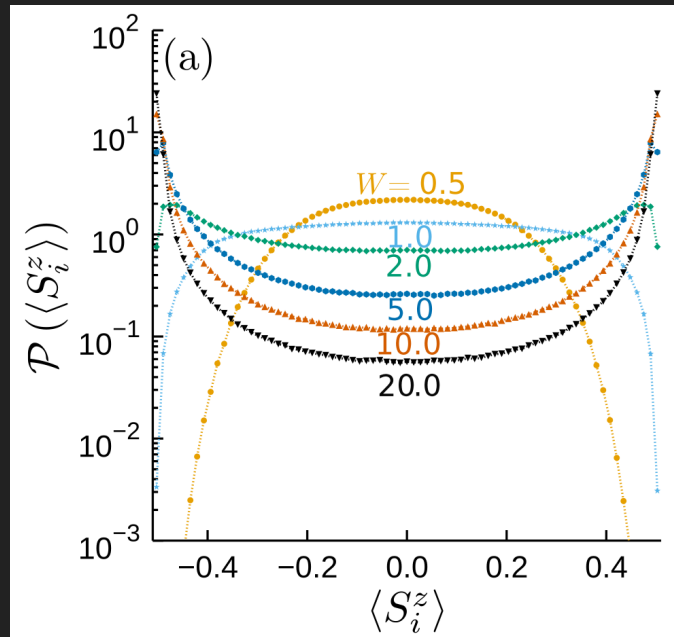
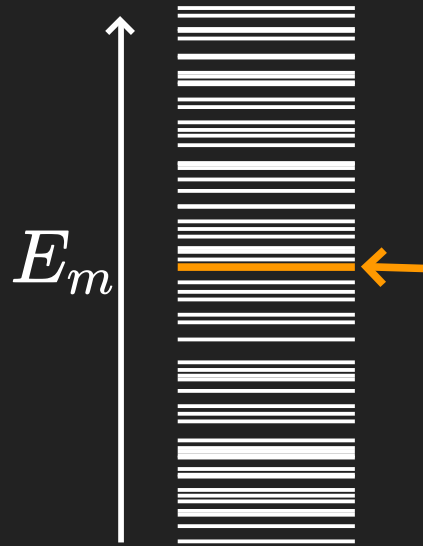
Spin-1/2 $W = h$ = disorder strength for random fields



"TRAILER"



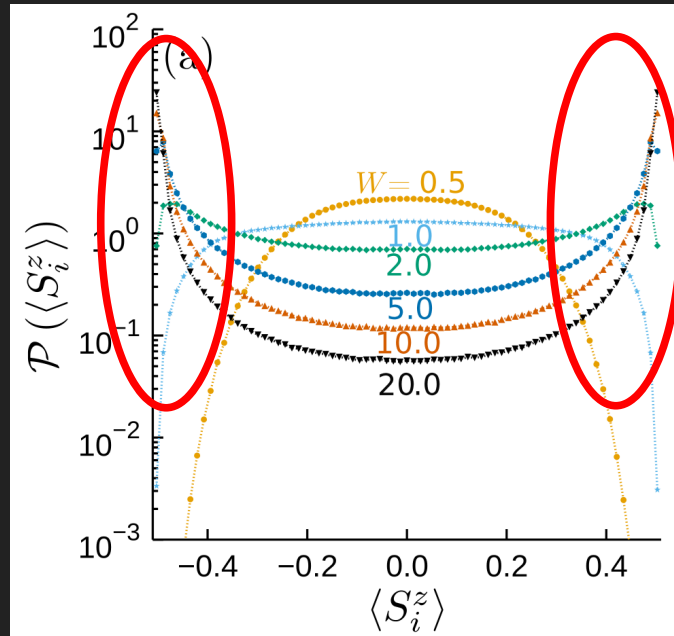
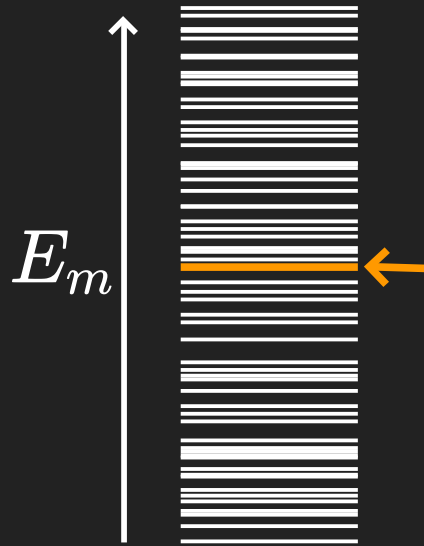
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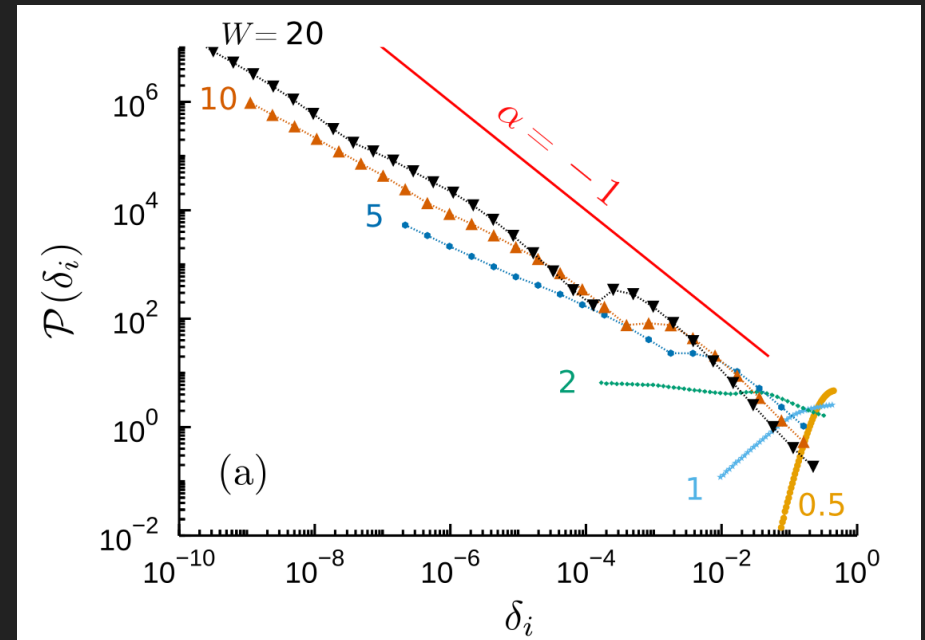
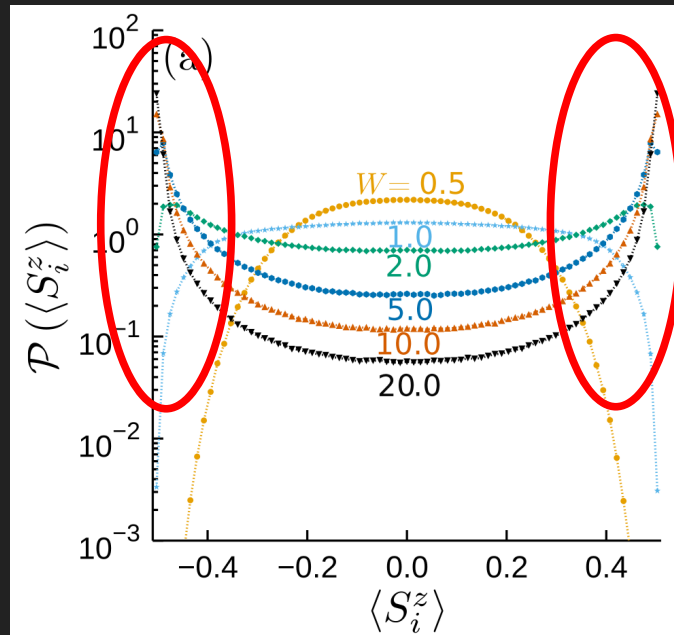
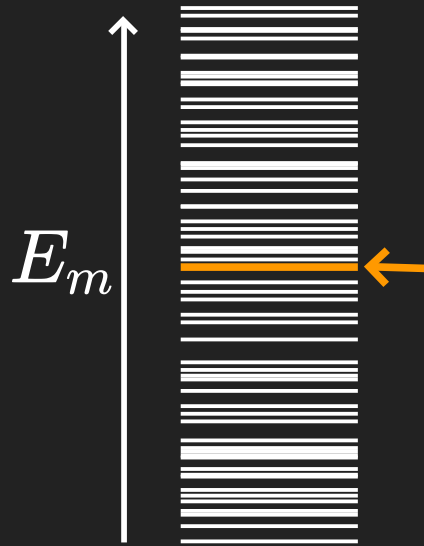
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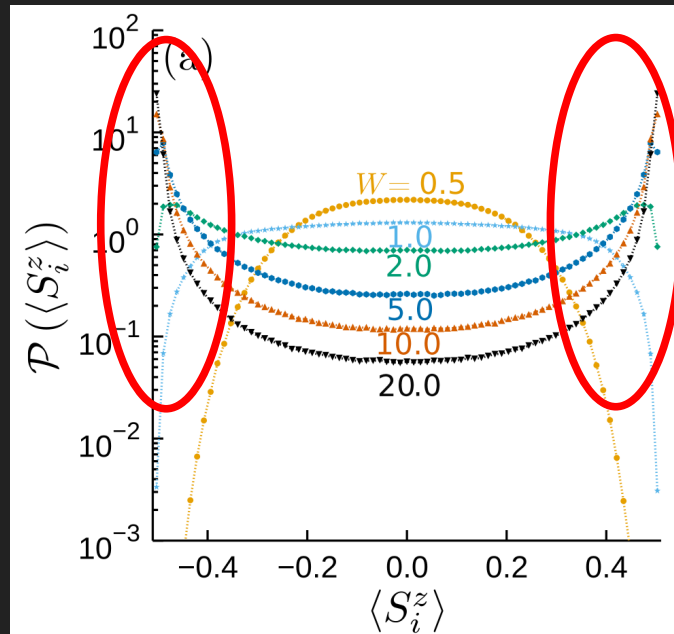
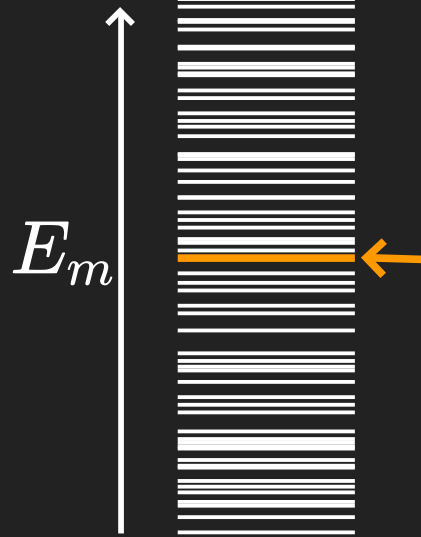


$$\delta_i = 1/2 - |\langle S_i^z \rangle|$$

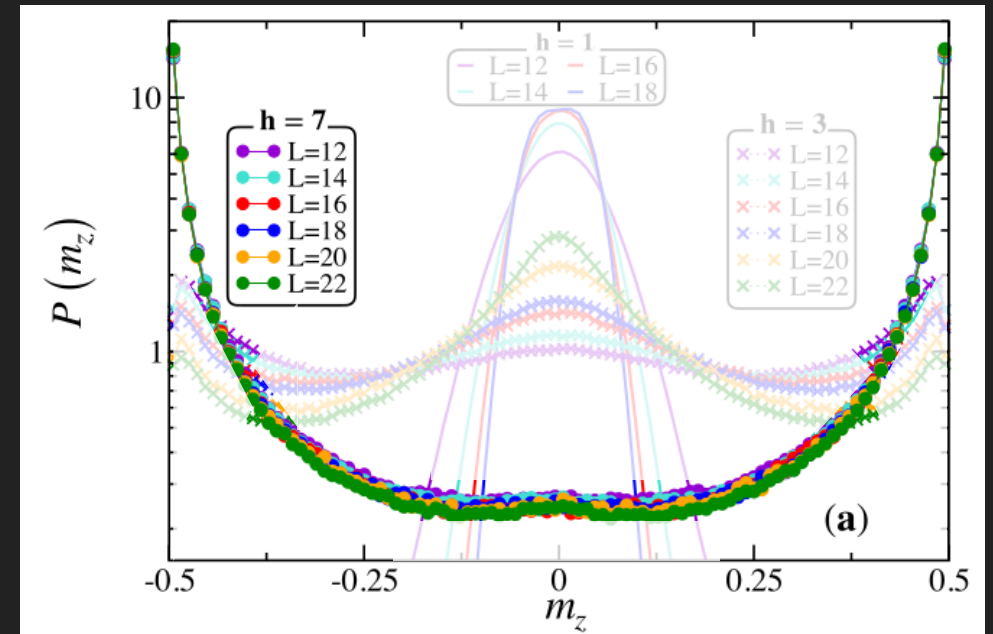
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Spin-1/2 $W = h =$ disorder strength for random fields



↓ **XX** ↑

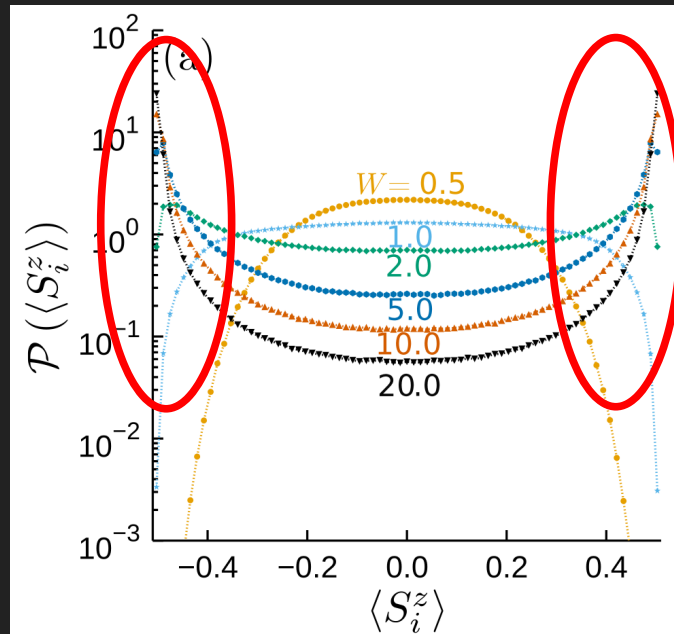
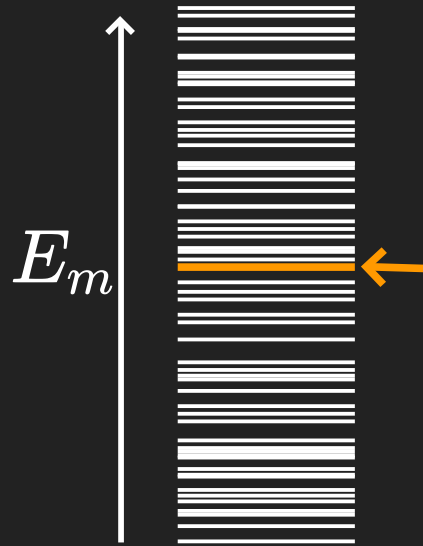


↓ **Heisenberg (XXX)** ↑

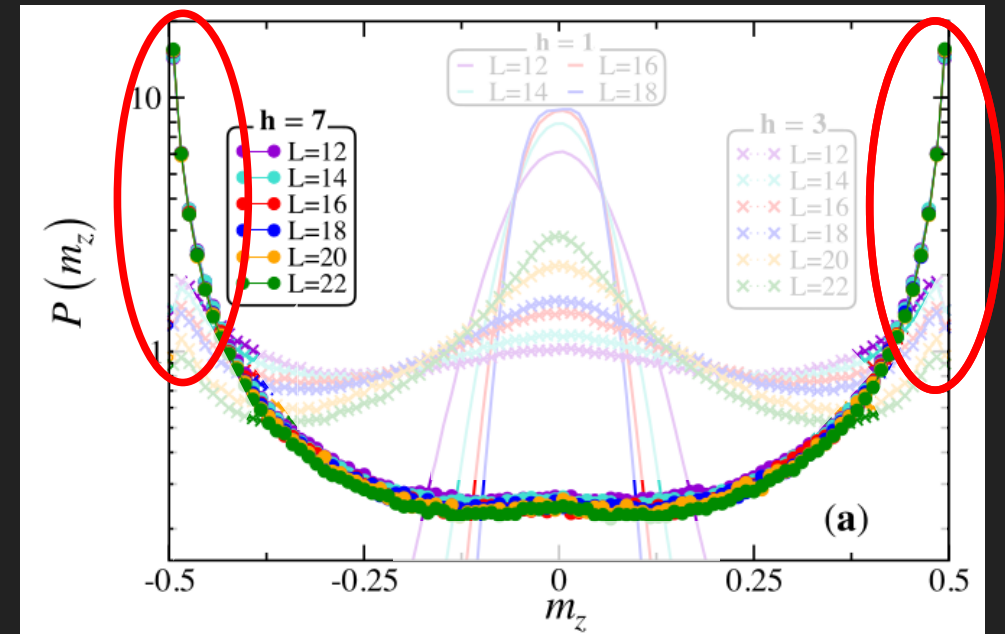
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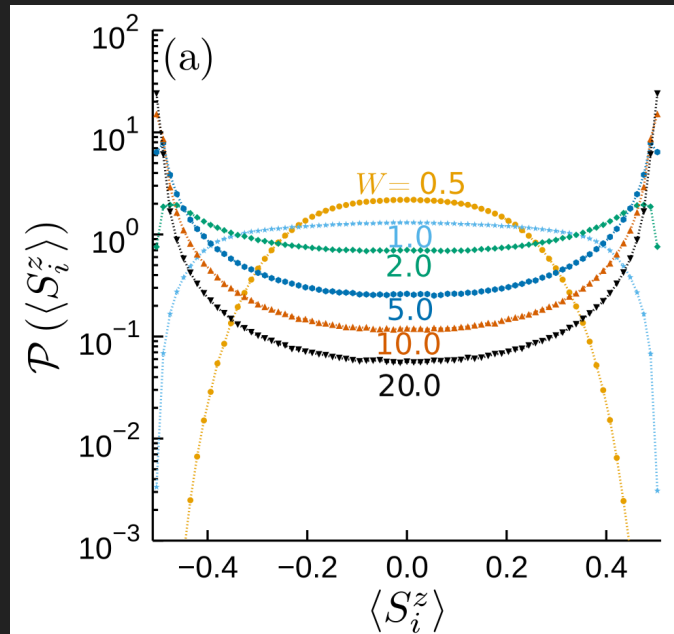
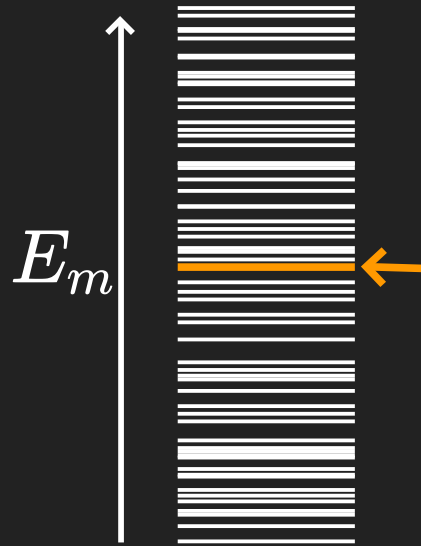


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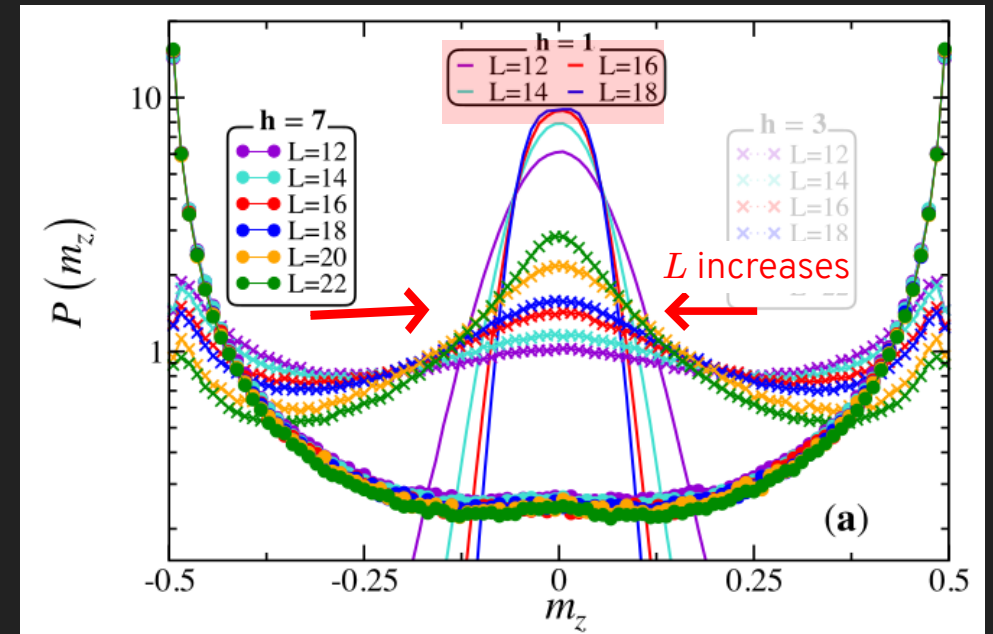
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↓ **Heisenberg (XXX)** ↑

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Spin-1/2 $W = h$ = disorder strength for random fields



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Spin-1/2 $W = h =$ disorder strength for random fields



Toy model?
Quantitative description?
Consequences?

SCOPE

1. Spin chains in random field and localization
2. Toy model : chain breaking!
3. Quantitative analysis: extreme value theory

SPIN CHAINS IN RANDOM FIELD AND LOCALIZATION

SPIN-1/2 CHAIN IN A RANDOM FIELD

$$\mathcal{H} = \sum_i \frac{J}{2} (S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+ + 2\Delta S_i^z S_{i+1}^z) - \sum_i h_i S_i^z$$

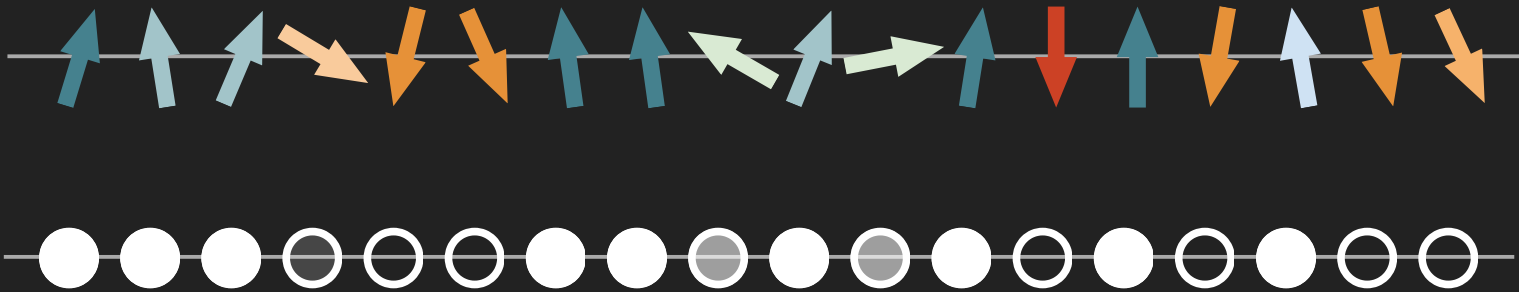
$$S^{x,y,z} = \frac{1}{2} \sigma^{x,y,z}$$



SPIN-1/2 CHAIN IN A RANDOM FIELD

$$\mathcal{H} = \sum_i \frac{J}{2} (S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+ + 2\Delta S_i^z S_{i+1}^z) - \sum_i h_i S_i^z$$

$$S^{x,y,z} = \frac{1}{2} \sigma^{x,y,z}$$



$$\mathcal{H}_f = \sum_i \left[\frac{J}{2} (c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i + 2\Delta n_i n_{i+1}) - h_i n_i \right]$$

Jordan-Wigner

Spinless fermions

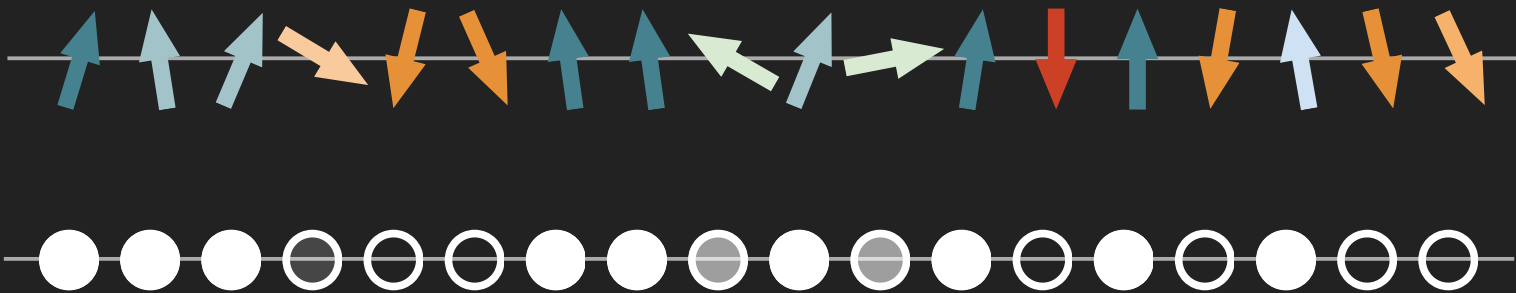
SPIN-1/2 CHAIN IN A RANDOM FIELD

$$\mathcal{H} = \sum_i \frac{J}{2} (S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+ + 2\Delta S_i^z S_{i+1}^z) - \sum_i h_i S_i^z \quad S^{x,y,z} = \frac{1}{2} \sigma^{x,y,z}$$

Spin-flip

Ising interaction

Magnetic field



Attraction/repulsion

$$\mathcal{H}_f = \sum_i \left[\frac{J}{2} (c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i + 2\Delta n_i n_{i+1}) - h_i n_i \right]$$

Jump
On-site energy

ANDERSON LOCALIZATION

$$\mathcal{H} = \sum_i \frac{J}{2} \left(\underbrace{S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+}_{\text{Spin-flip}} + \underbrace{2\Delta S_i^z S_{i+1}^z}_{\text{Ising interaction}} \right) - \sum_i h_i S_i^z \quad \text{1 particle}$$

Spin-flip Ising interaction Magnetic field

ANDERSON LOCALIZATION

$$\mathcal{H}_f = \sum_i \left[\frac{J}{2} \left(c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i \right) - h_i n_i \right]$$

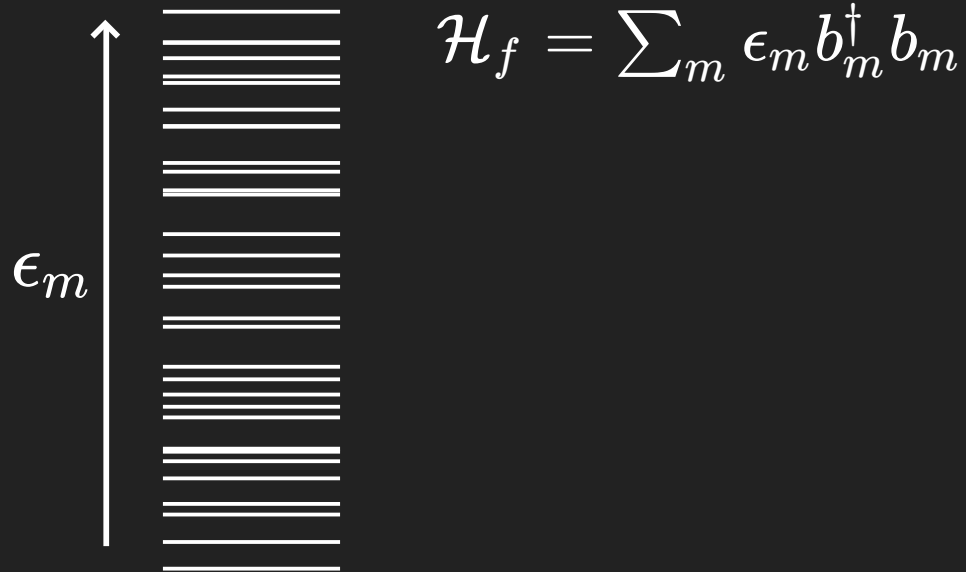
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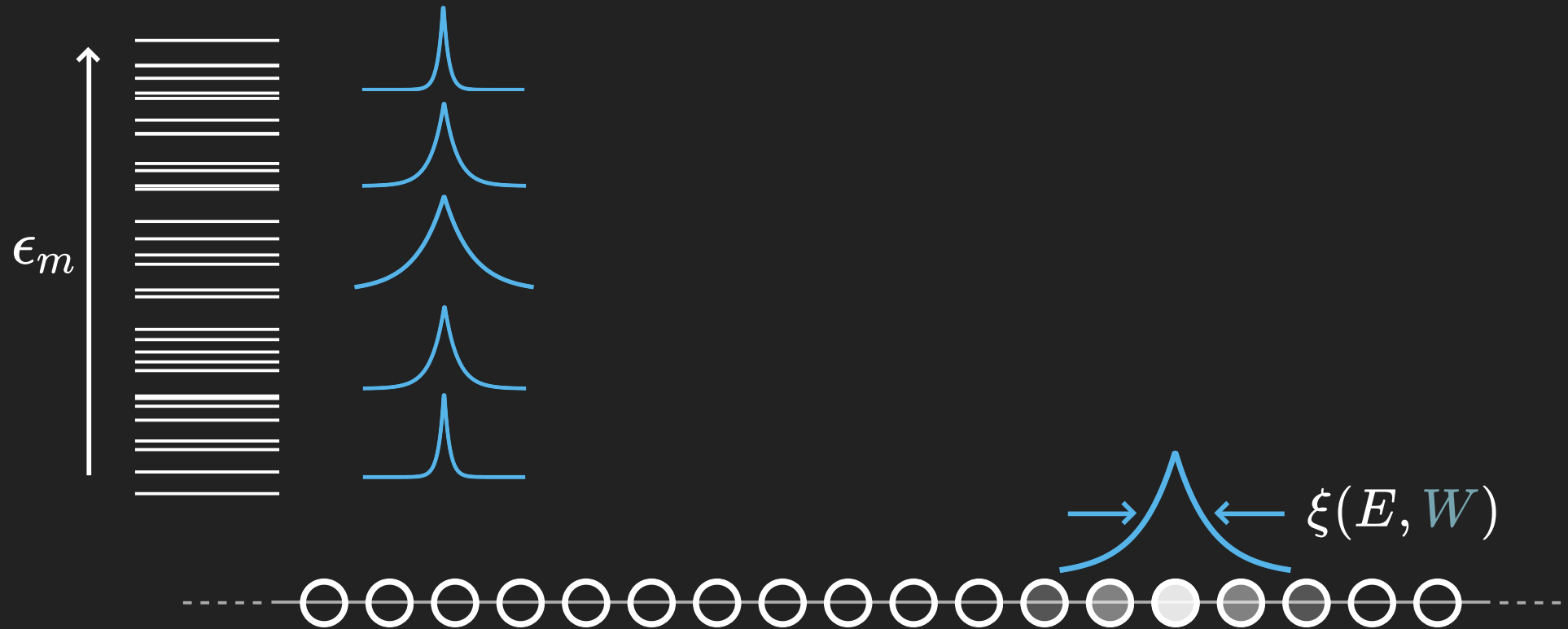
P. W. Anderson, Phys. Rev. **109**, 1492 (1958)

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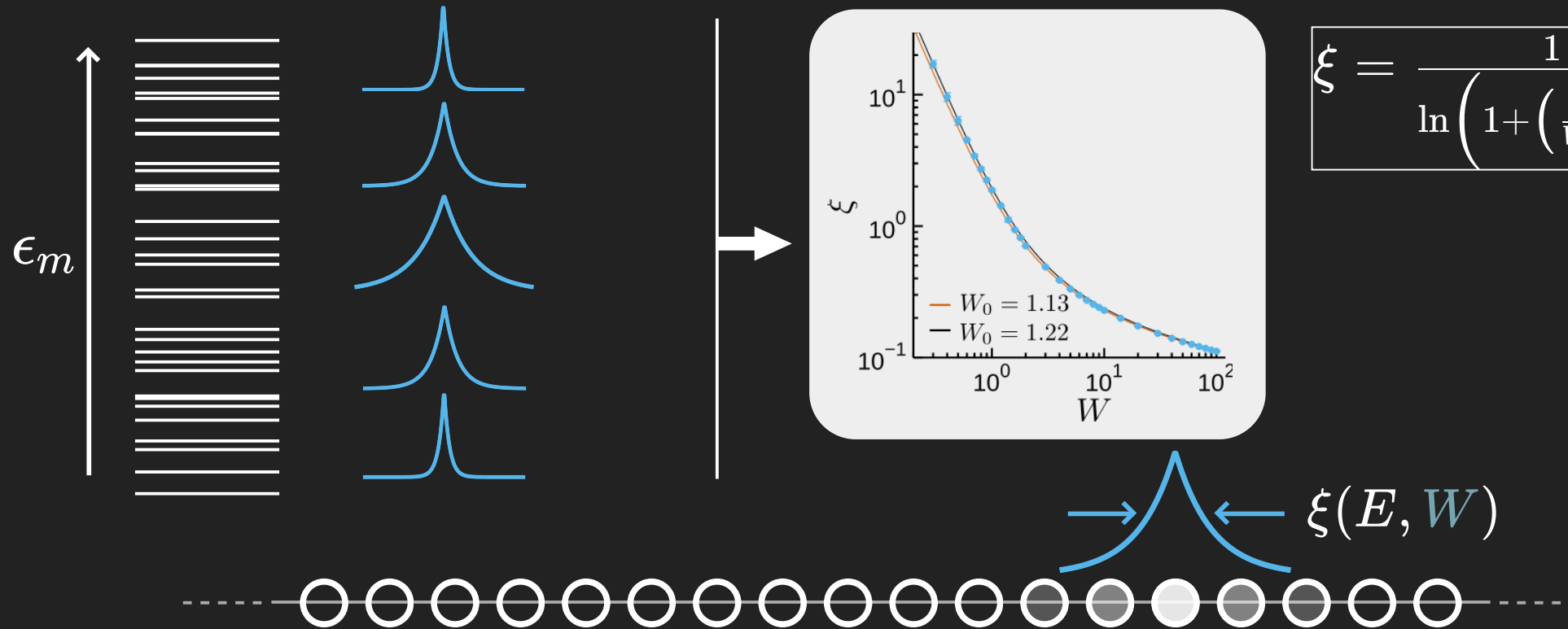
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1 particle



$$\xi = \frac{1}{\ln \left(1 + \left(\frac{W}{W_0} \right)^2 \right)}$$



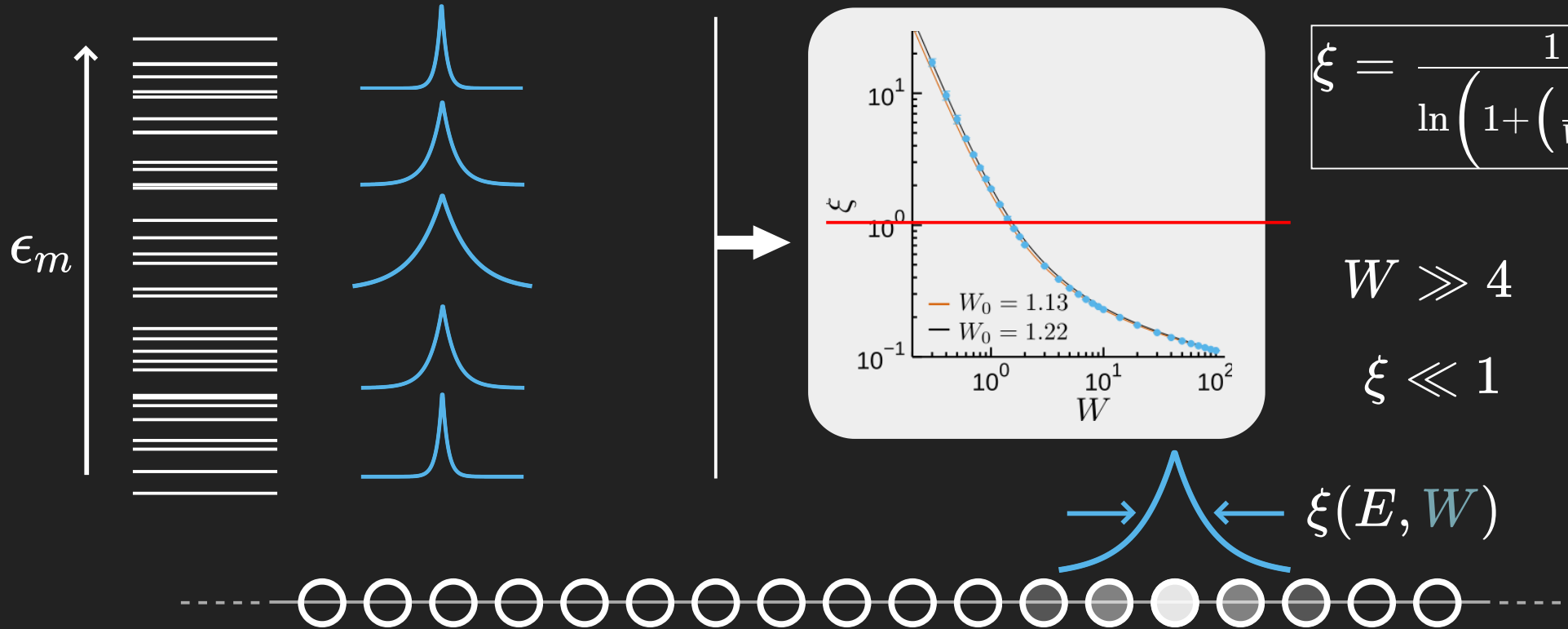
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1 particle



$$\xi = \frac{1}{\ln \left(1 + \left(\frac{W}{W_0} \right)^2 \right)}$$

$$W \gg 4$$

$$\xi \ll 1$$



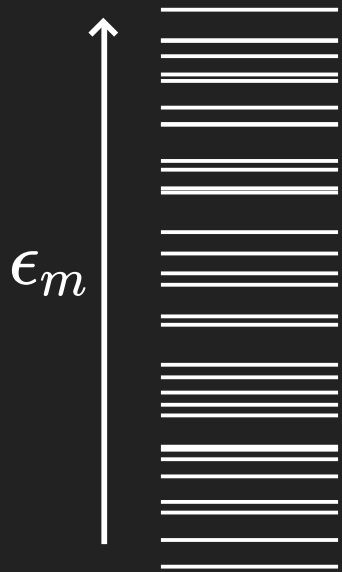
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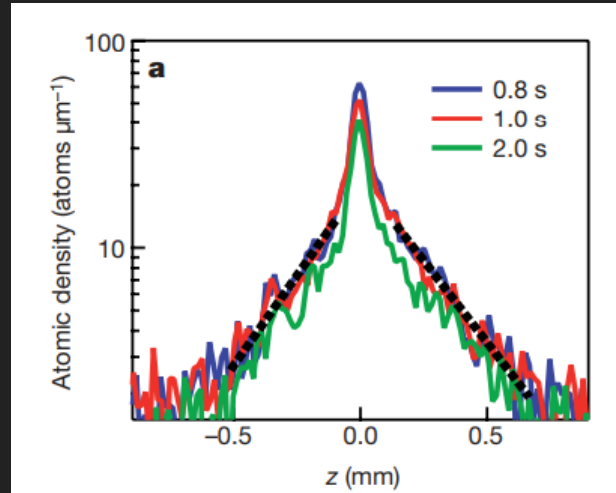
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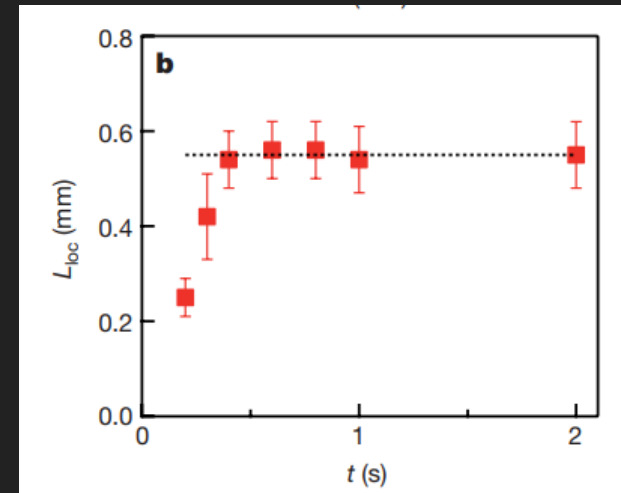
1 particle



$$\mathcal{H}_f = \sum_m \epsilon_m b_m^\dagger b_m$$



Billy J, et al. Direct observation of Anderson localization of matter waves in a controlled disorder. Nature. 2008



P. W. Anderson, Phys. Rev. **109**, 1492 (1958)

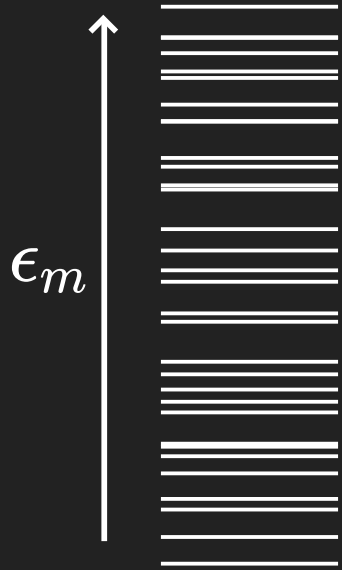
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MANY-BODY ANDERSON INSULATOR

$$\mathcal{H}_f = \sum_i \left[\frac{J}{2} \left(c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i \right) - h_i n_i \right]$$

$L/2$ fermions

$$S_z = 0$$



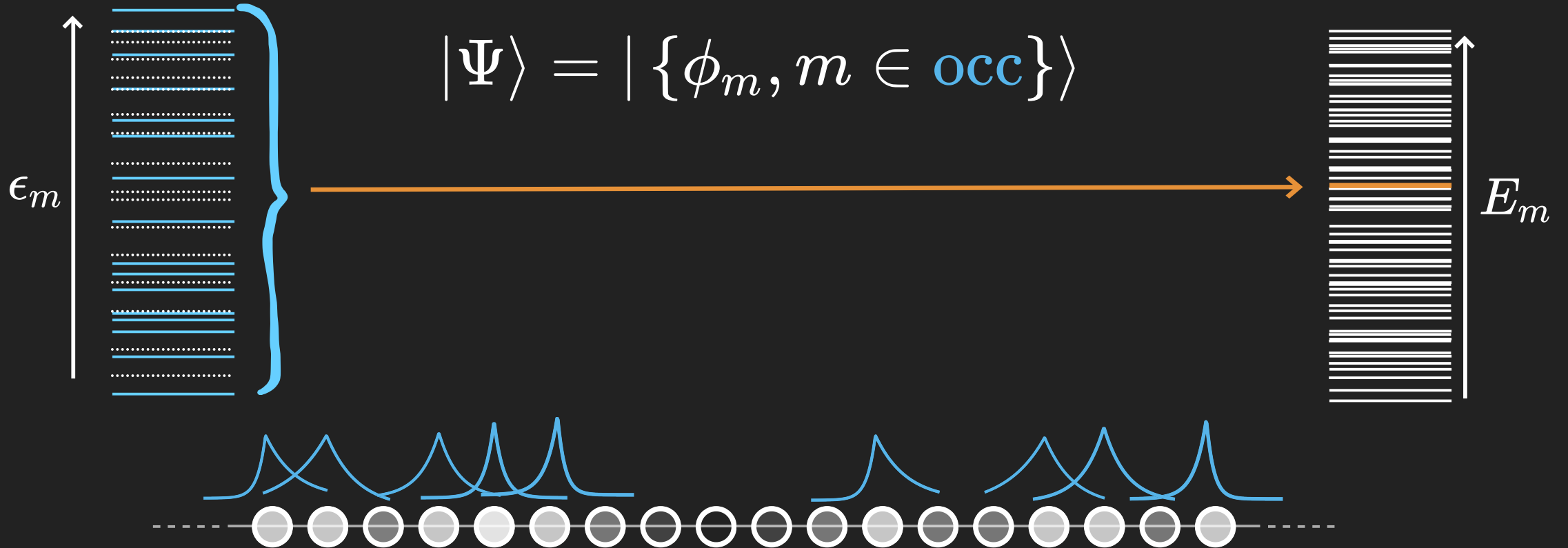
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$L/2$ fermions

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$$|\Psi\rangle = |\{\phi_m, m \in \text{occ}\}\rangle$$



HEISENBERG: INTRODUCING INTERACTIONS

$$\mathcal{H} = \sum_i \frac{J}{2} \left(S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+ + 2\Delta S_i^z S_{i+1}^z \right) - \sum_i h_i S_i^z$$

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Fate of isolated quantum **many-body** systems ?

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In the Anderson basis:

$$\mathcal{H} = \sum_m \epsilon_m b_m^\dagger b_m + \sum_{j,k,l,m} V_{j,k,l,m} b_j^\dagger b_k^\dagger b_l b_m$$



Anderson
orbitals m

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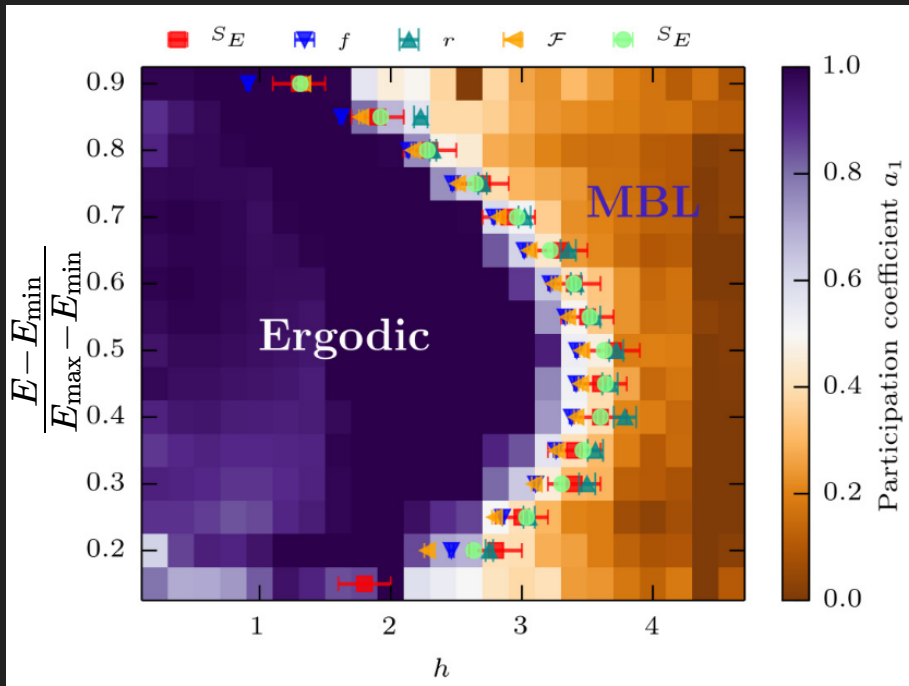
Anderson
orbitals m

Challenge : high-energy eigenstate, lack of symmetries. Shift-invert ED $\rightarrow L = 22$

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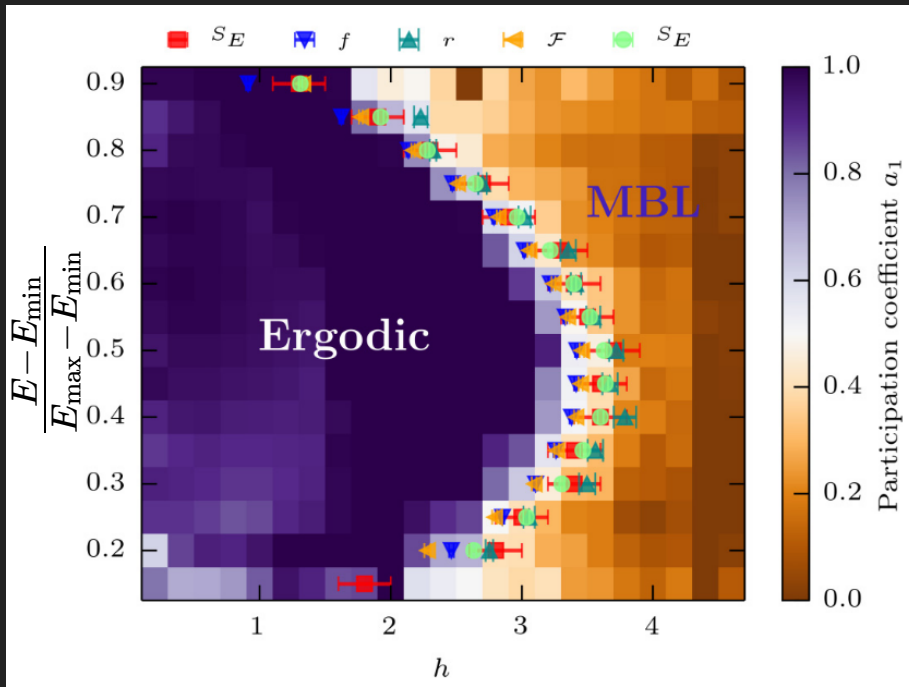
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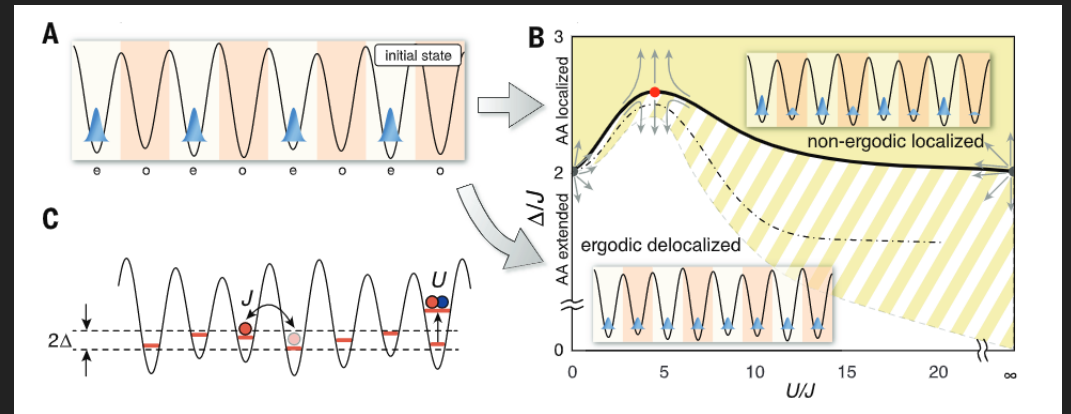
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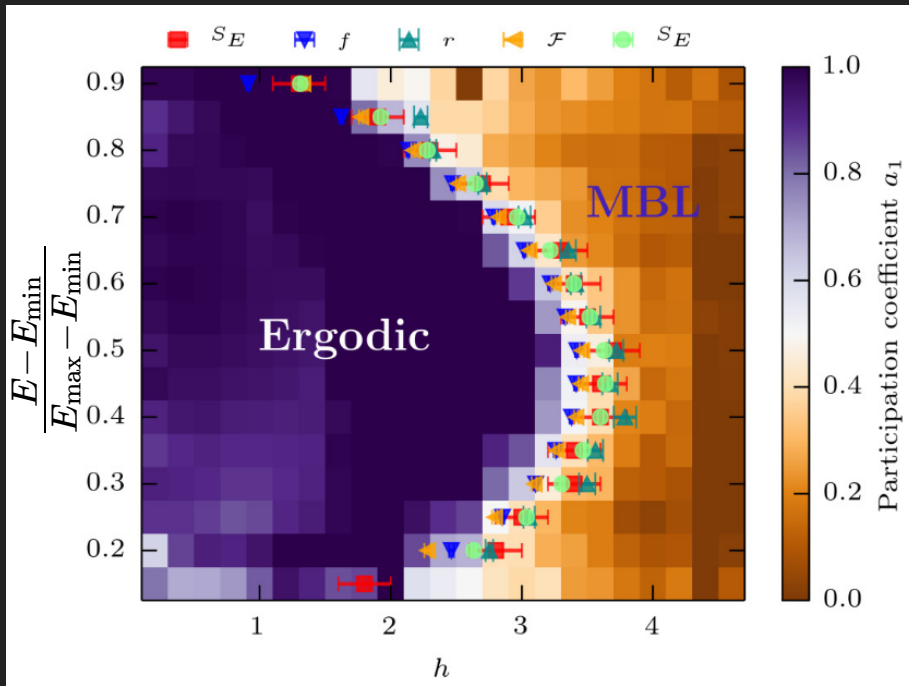
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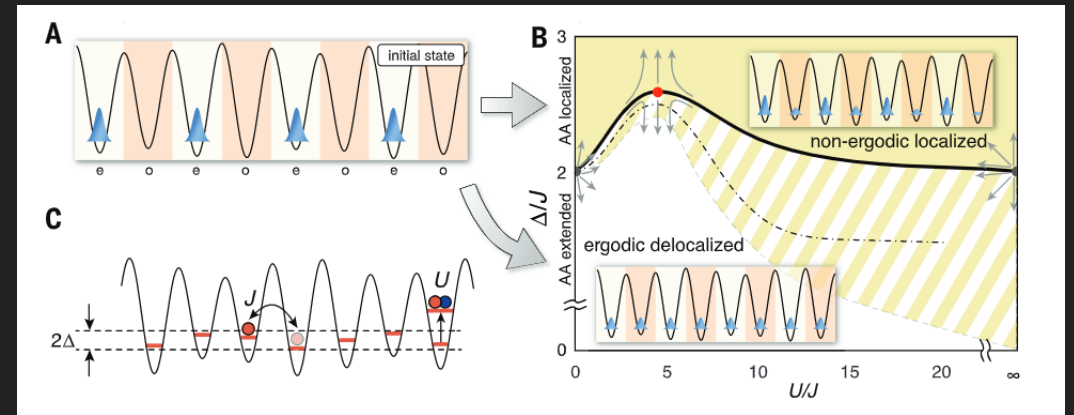
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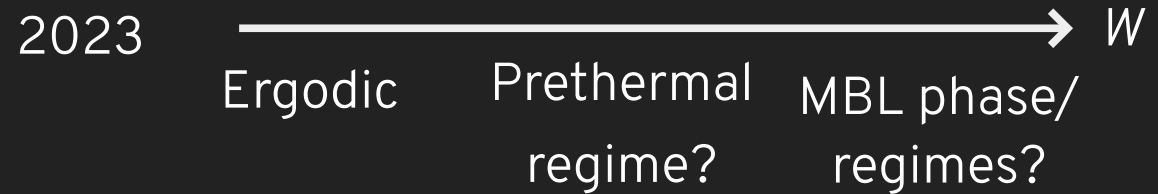
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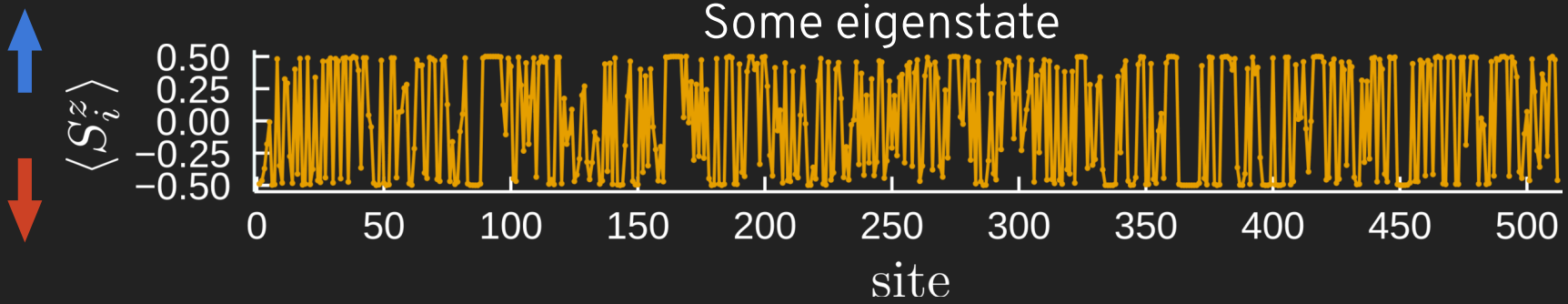


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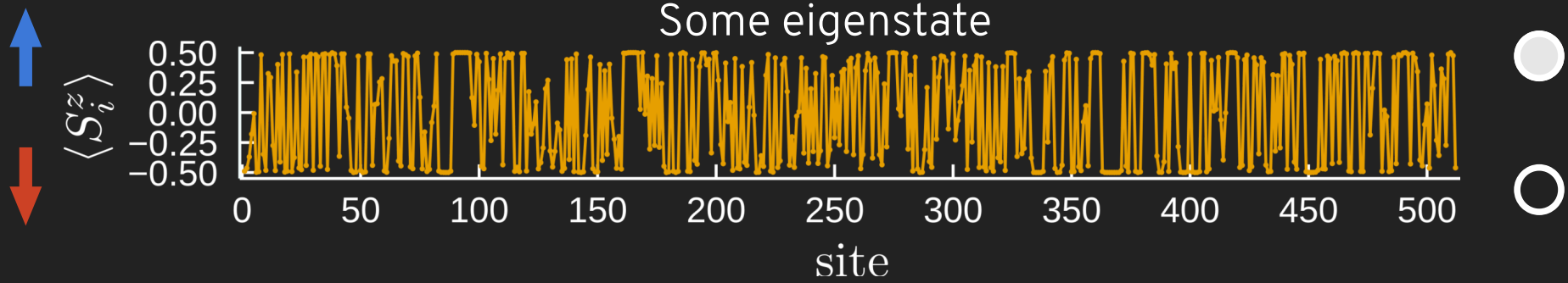


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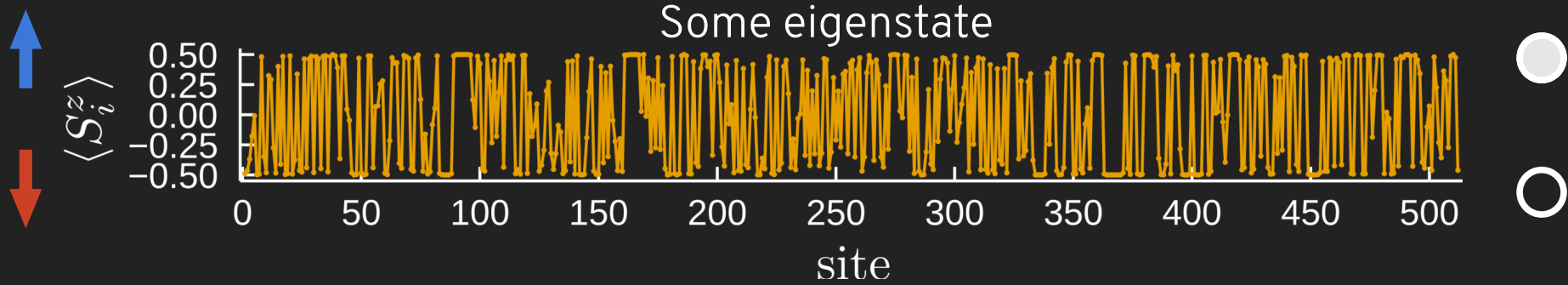
ED ON ONE SAMPLE - XX (ANDERSON) CHAIN



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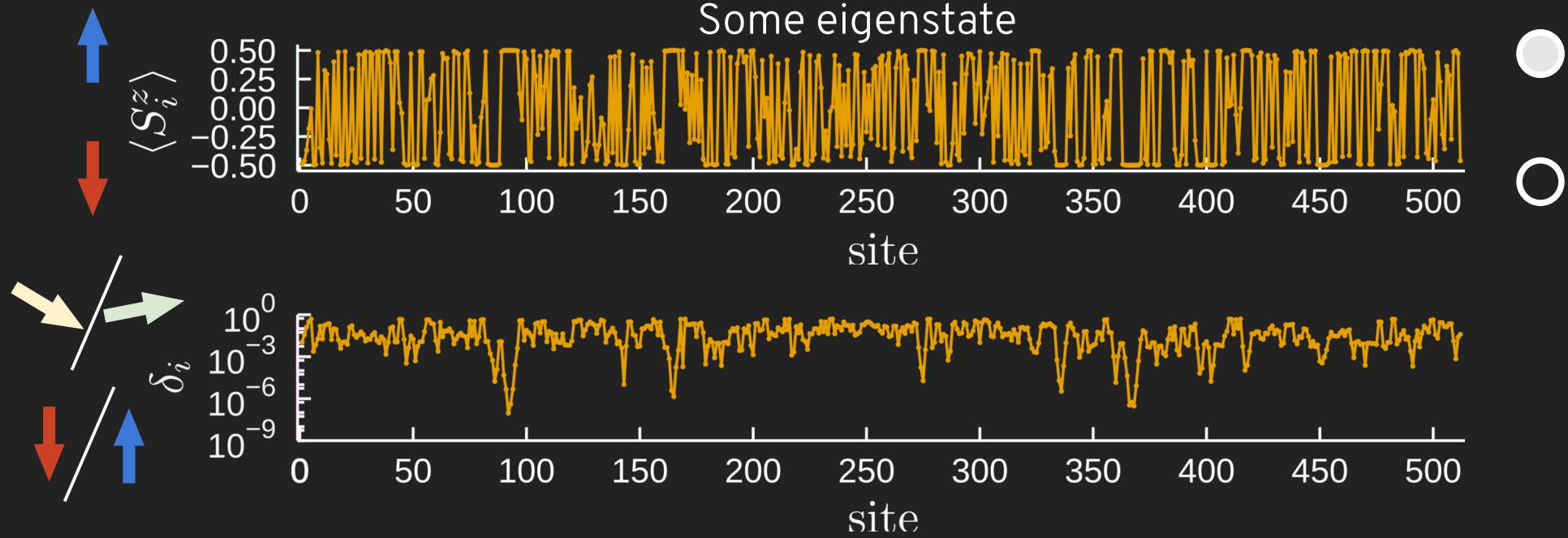


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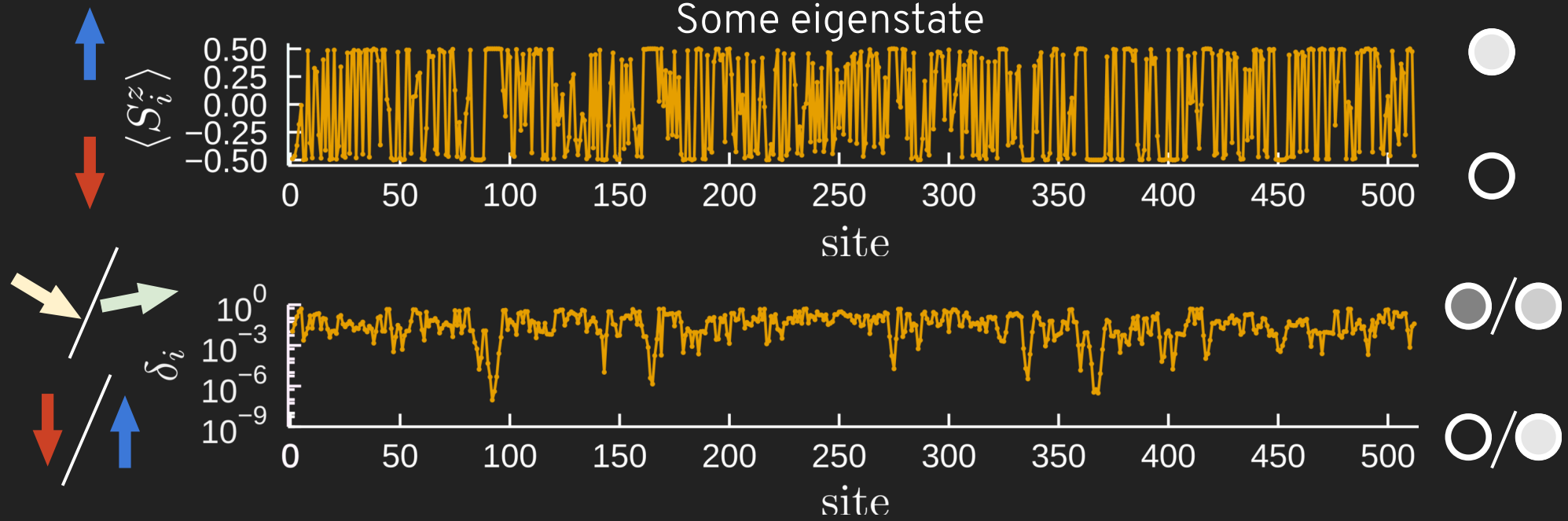
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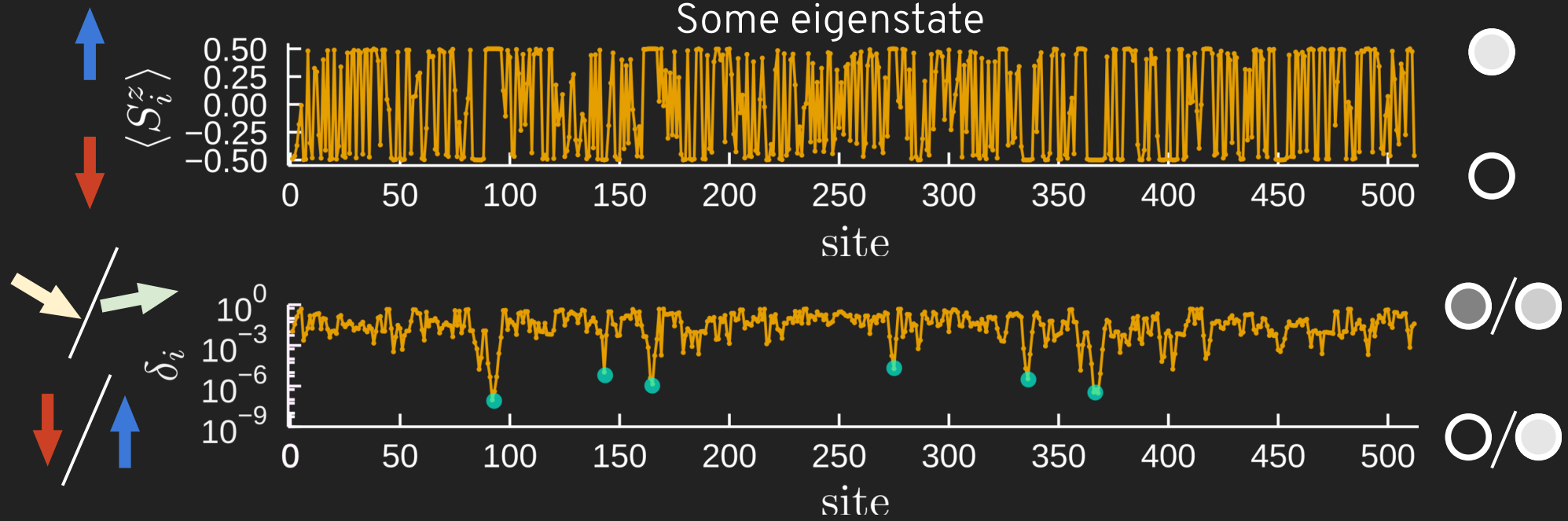
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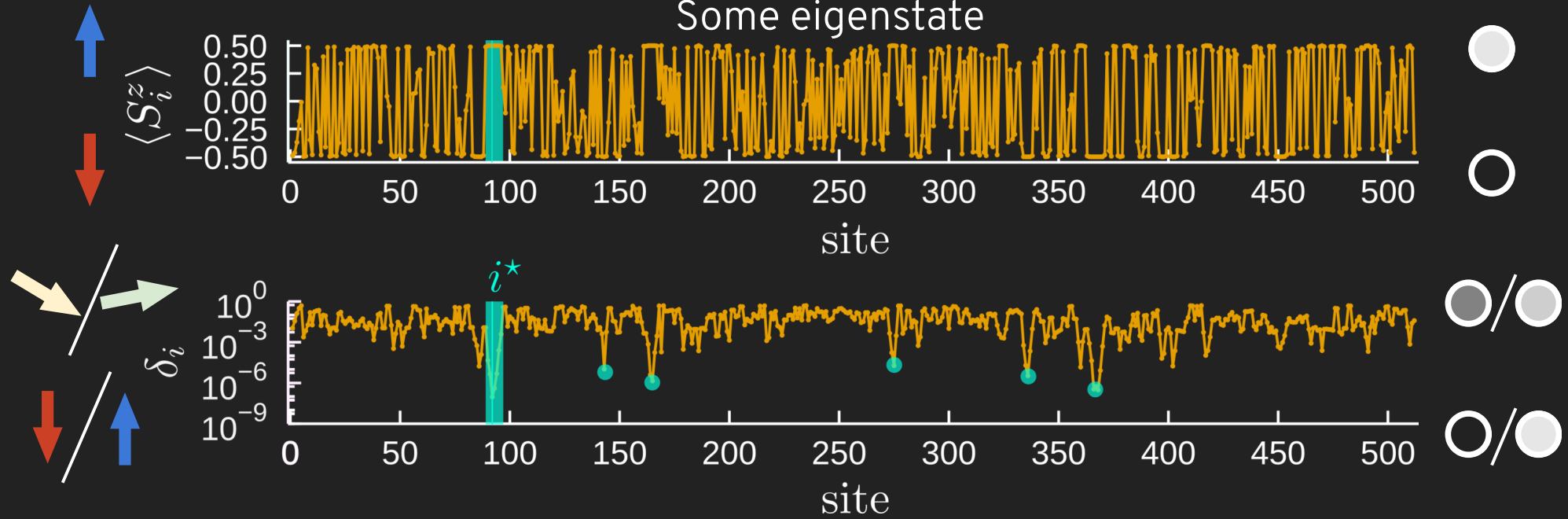
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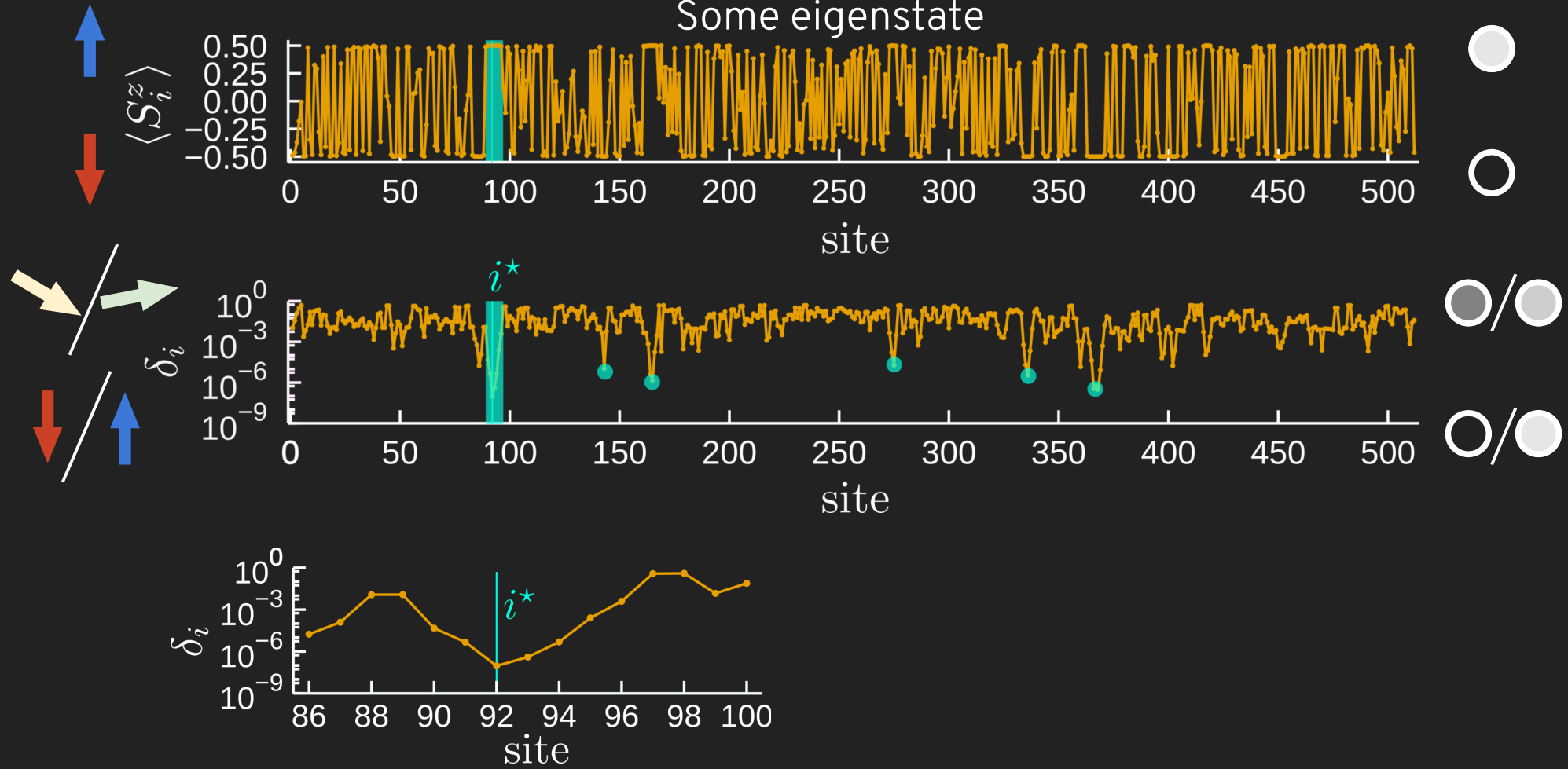
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SPIN FREEZING



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SPIN FREEZING



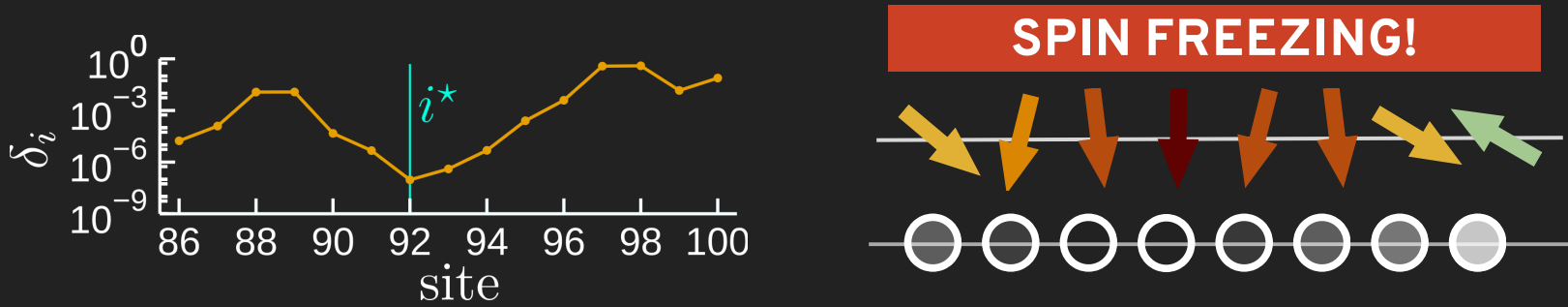
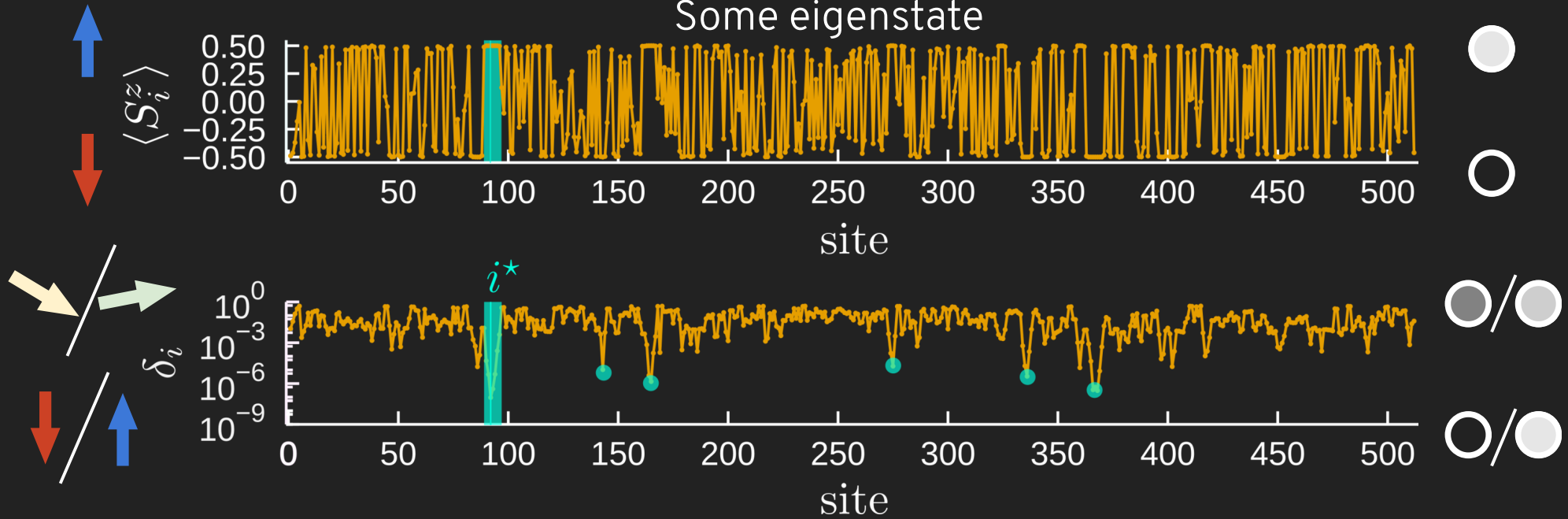
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TOY MODEL : CHAIN BREAKING!

TOY MODEL



Dupont, Macé, Laflorencie, PRB **100**, 134201, (2019)
Laflorencie, Lemarié, Macé, PRR **2**, 042033(R), (2020)
JC, N. Laflorencie, arXiv:2305.10574

TOY MODEL

$$|\phi_m(i)|^2 \propto \exp\left(-\frac{|i-i_0^m|}{\xi}\right)$$



TOY MODEL

$$|\phi_m(i)|^2 \propto \exp\left(-\frac{|i-i_0^m|}{\xi}\right)$$

$$\Rightarrow \langle n_i \rangle = \langle S_i^z \rangle + 1/2 = \sum_{m \in \text{occ}} |\phi_m(i)|^2$$



MINIMAL DEVIATION?

$$\delta_i = 1/2 - |\langle n_i \rangle - 1/2|$$

$$|\phi_m(i)|^2 \propto \exp\left(-\frac{|i-i_0^m|}{\xi}\right)$$

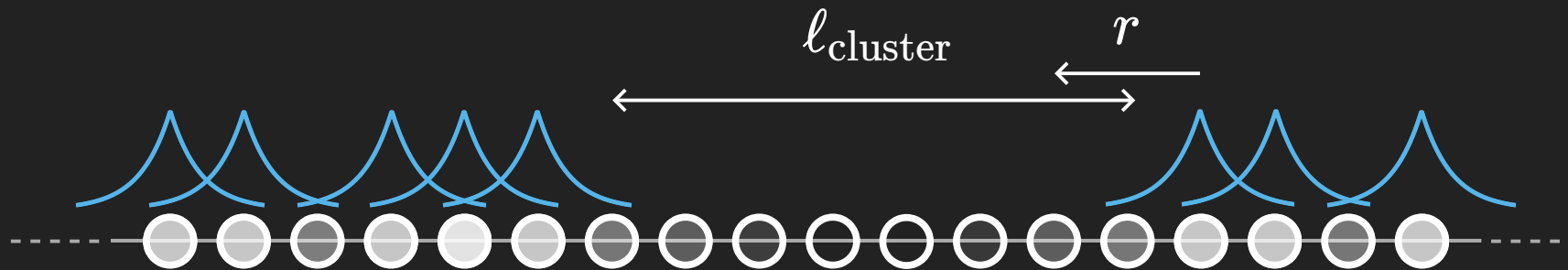


MINIMAL DEVIATION?

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$$\bigcirc : \delta_i = \langle n_i \rangle$$

$$|\phi_m(i)|^2 \propto \exp\left(-\frac{|i-i_0^m|}{\xi}\right)$$

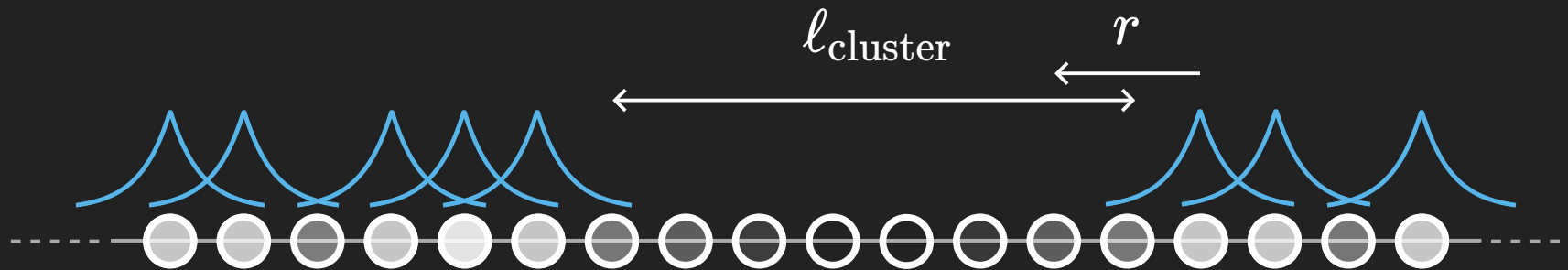


MINIMAL DEVIATION?

$$\delta_i = 1/2 - |\langle n_i \rangle - 1/2|$$

$$\bigcirc : \delta_i = \langle n_i \rangle \approx e^{-\frac{r}{\xi}} + \dots$$

$$|\phi_m(i)|^2 \propto \exp\left(-\frac{|i-i_0^m|}{\xi}\right)$$

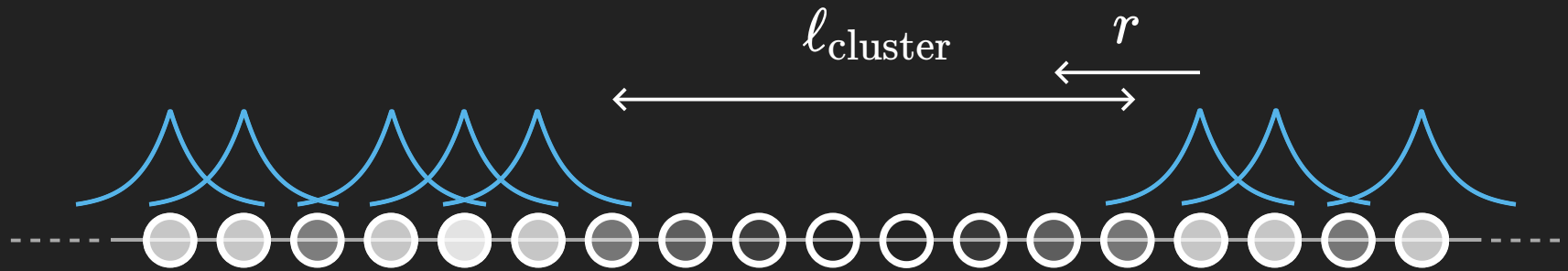


MINIMAL DEVIATION?

$$\delta_i = 1/2 - |\langle n_i \rangle - 1/2|$$

$$\text{O} : \delta_i = \langle n_i \rangle \approx e^{-\frac{r}{\xi}} + \dots$$

$$|\phi_m(i)|^2 \propto \exp\left(-\frac{|i-i_0^m|}{\xi}\right) \Rightarrow \delta_{\min} \approx e^{-\frac{\ell_{\text{cluster}}}{2\xi}}$$

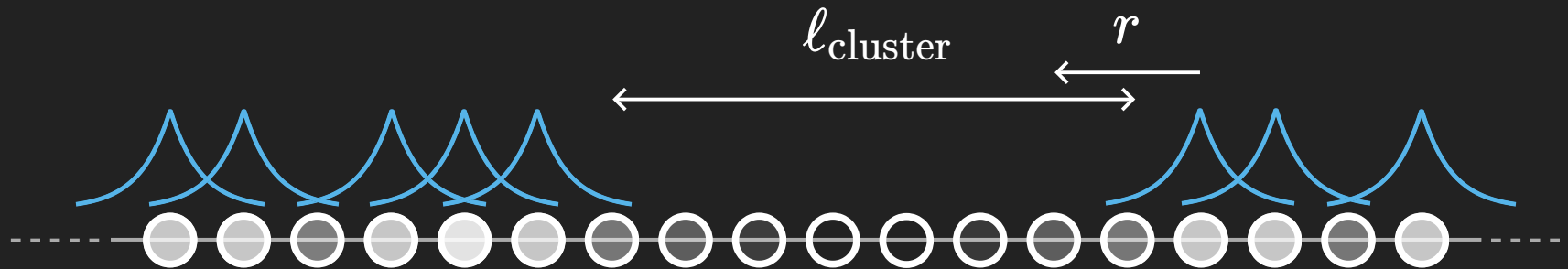
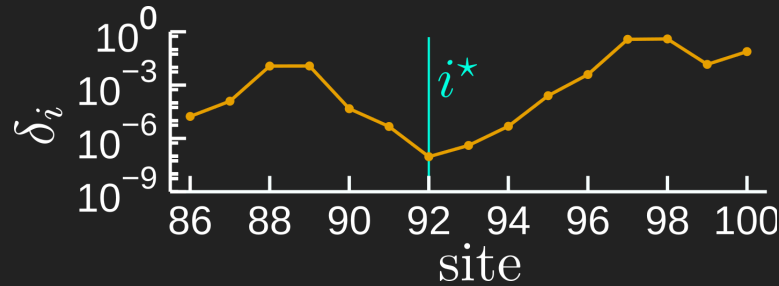


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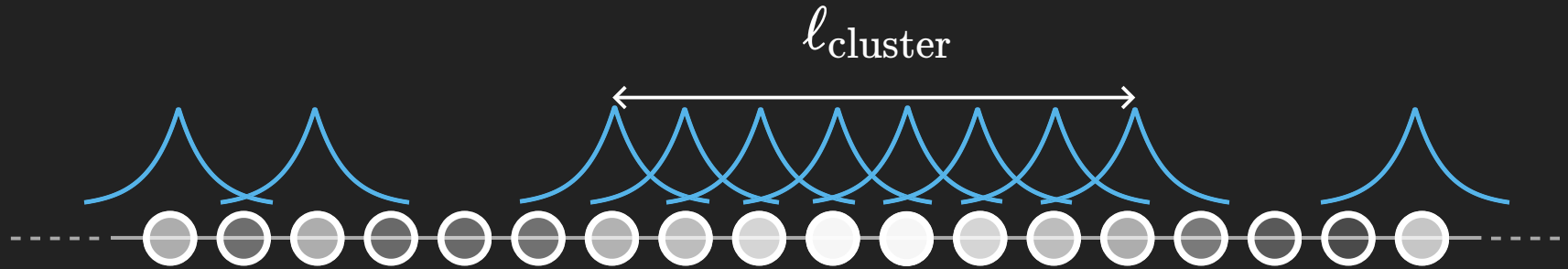
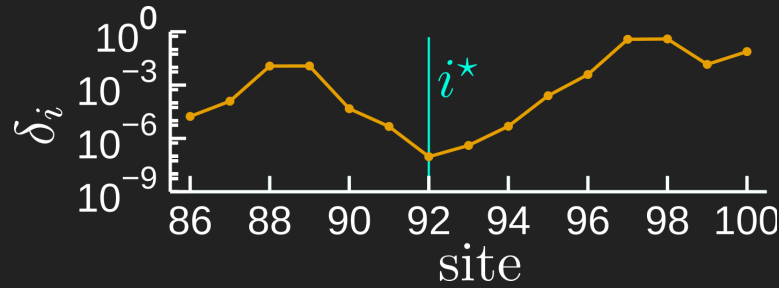


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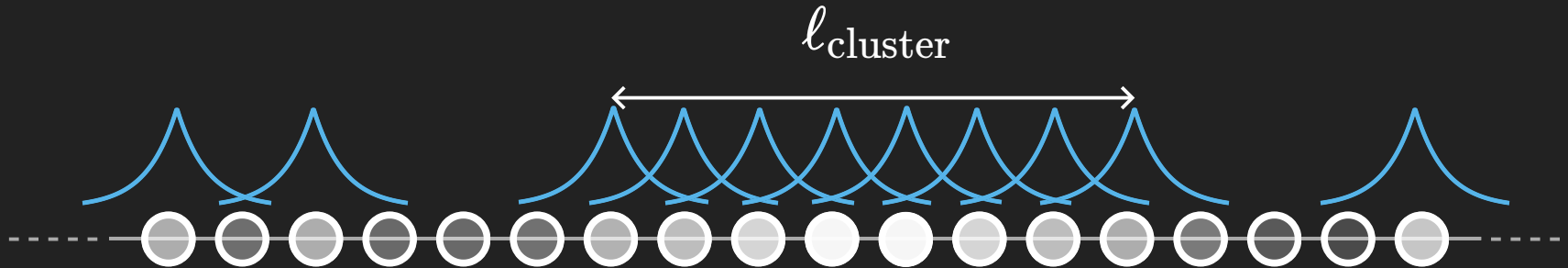
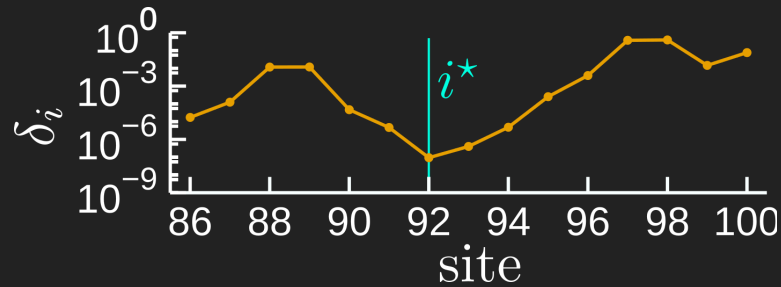


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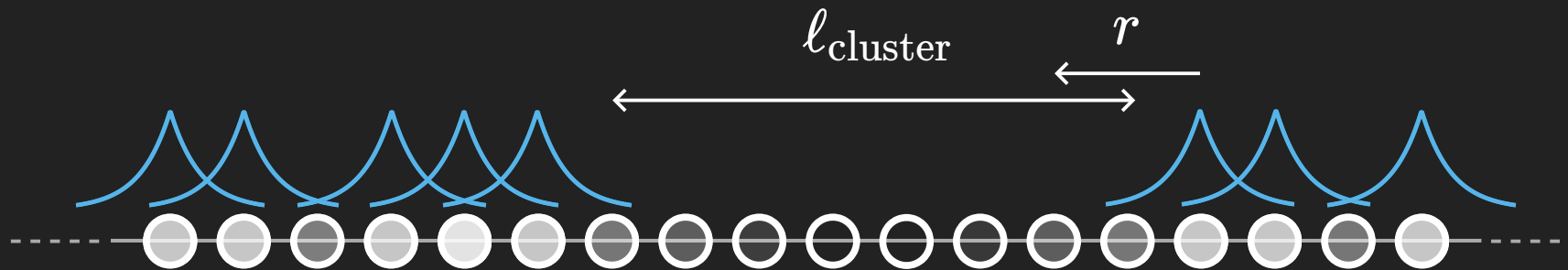
$$\Rightarrow \delta_{\min}^{\text{typ}} \approx e^{-\frac{\ell_{\text{cluster}}}{2\xi}}$$



Dupont, Macé, Laflorencie, PRB **100**, 134201, (2019)
 Laflorencie, Lemarié, Macé, PRR **2**, 042033(R), (2020)
 JC, N. Laflorencie, arXiv:2305.10574

MINIMAL DEVIATION?

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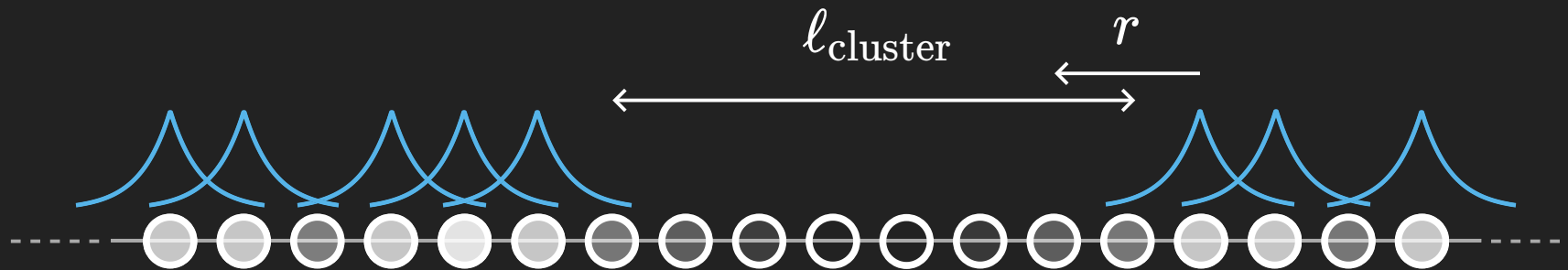
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MINIMAL DEVIATION?

$$\delta_{\min}^{\text{typ}} \approx e^{-\frac{\overline{l_{\text{cluster}}}}{2\xi}}$$



$$\overline{l_{\text{cluster}}} \approx \frac{\ln L}{\ln 2}$$



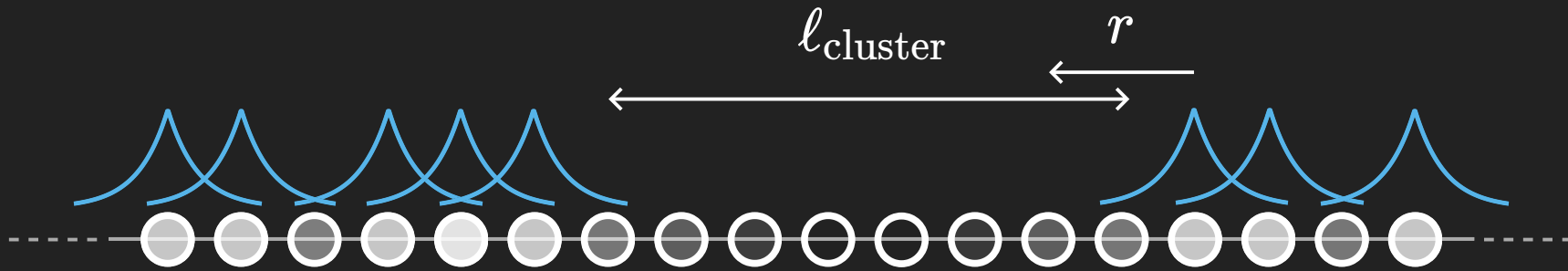
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$$\delta_{\min}^{\text{typ}} \approx L^{-\frac{1}{2\xi \ln 2}}$$



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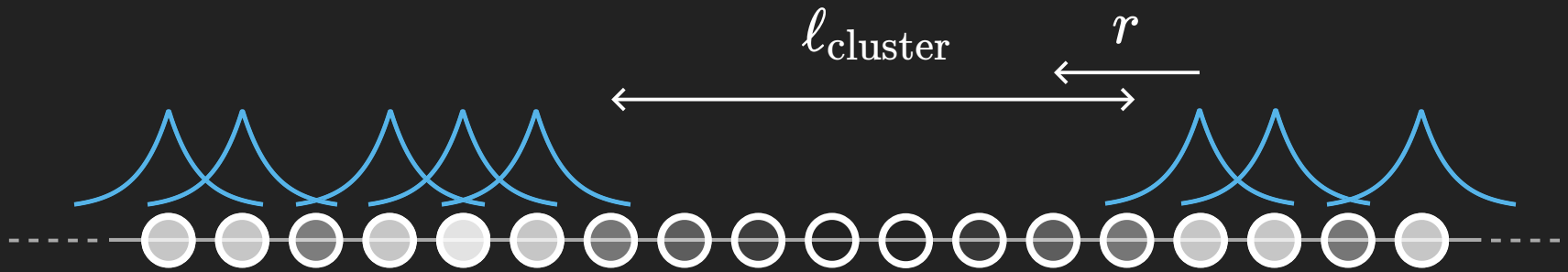
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EXPONENT : TOY MODEL

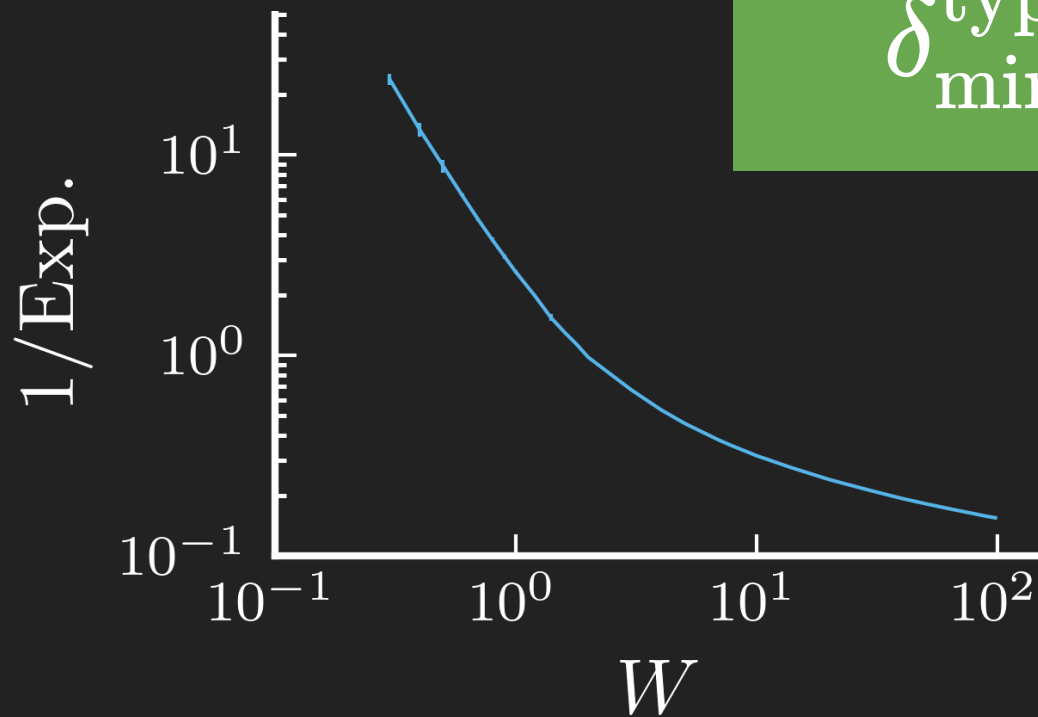
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$$\xi = \frac{1}{\ln\left(1 + \left(\frac{w}{w_0}\right)^2\right)}$$

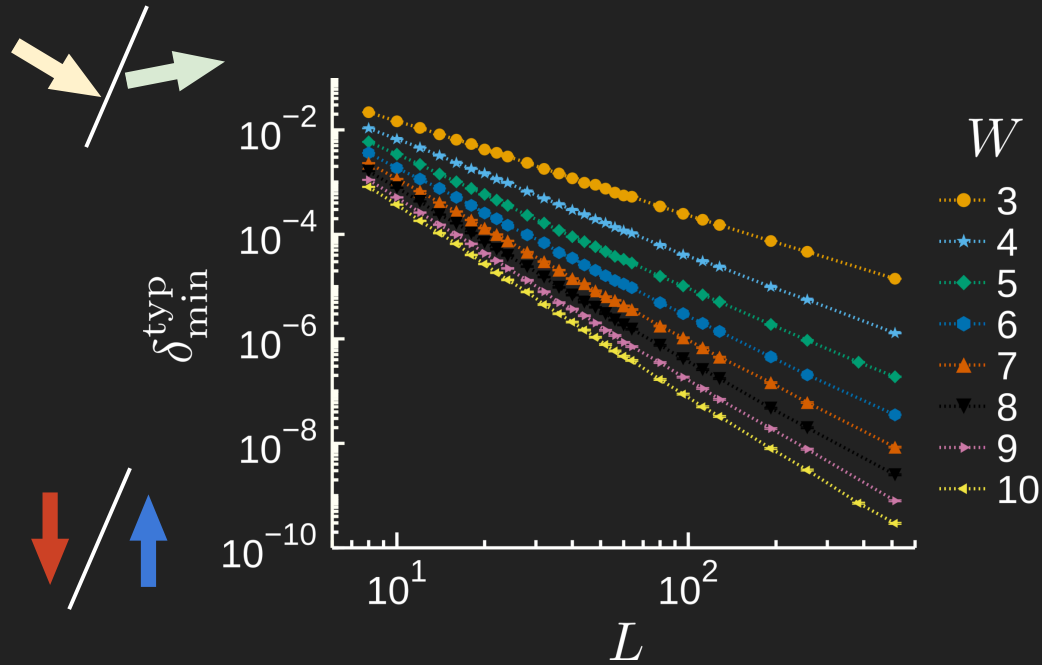
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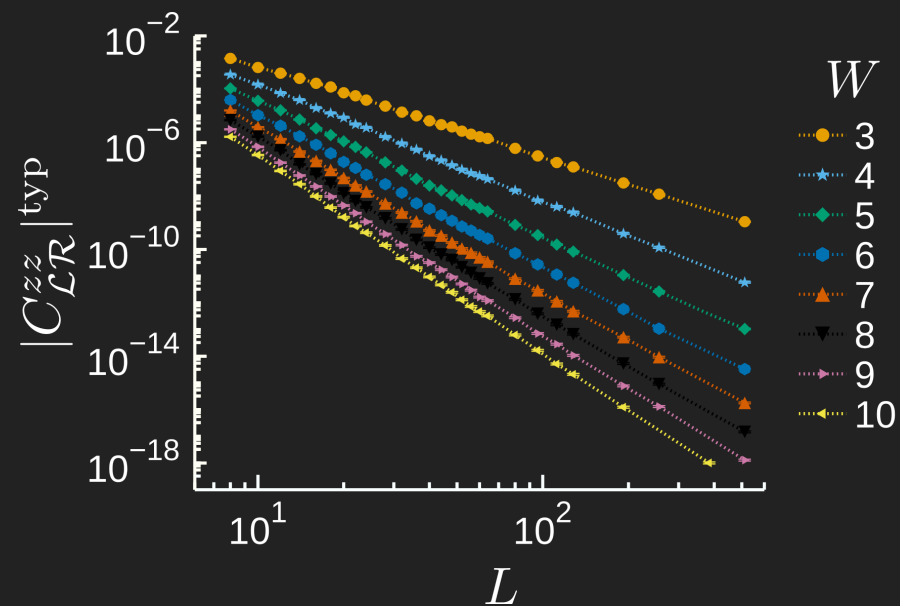
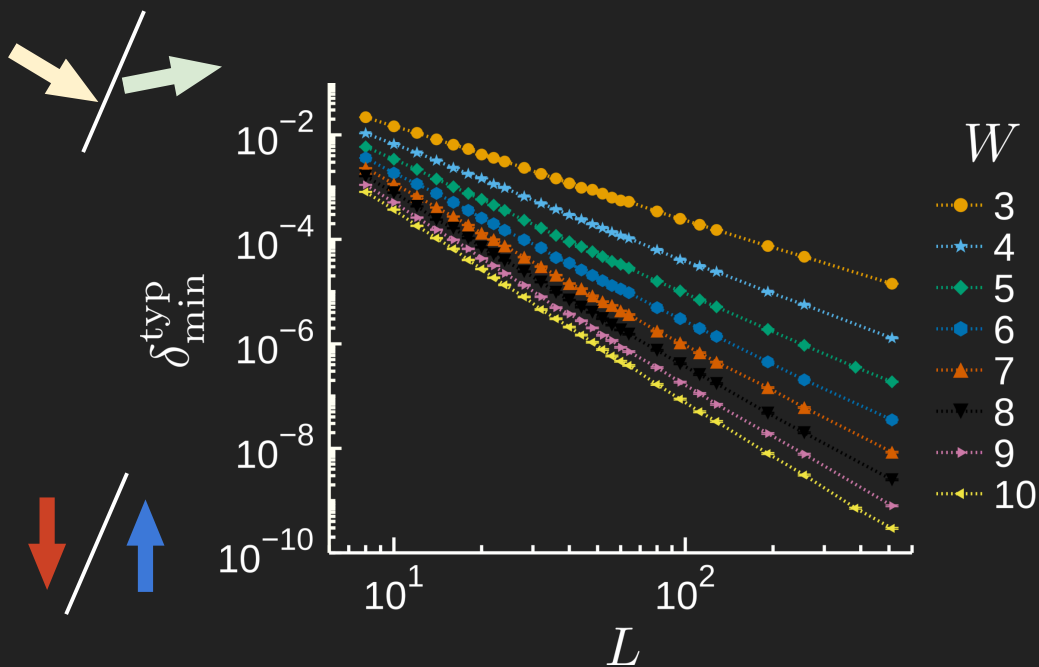
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CHAIN BREAKING



$$\delta_{\min}^{\text{typ}} = e^{\overline{\ln \delta_{\min}}} \approx L^{-\gamma(W)}$$

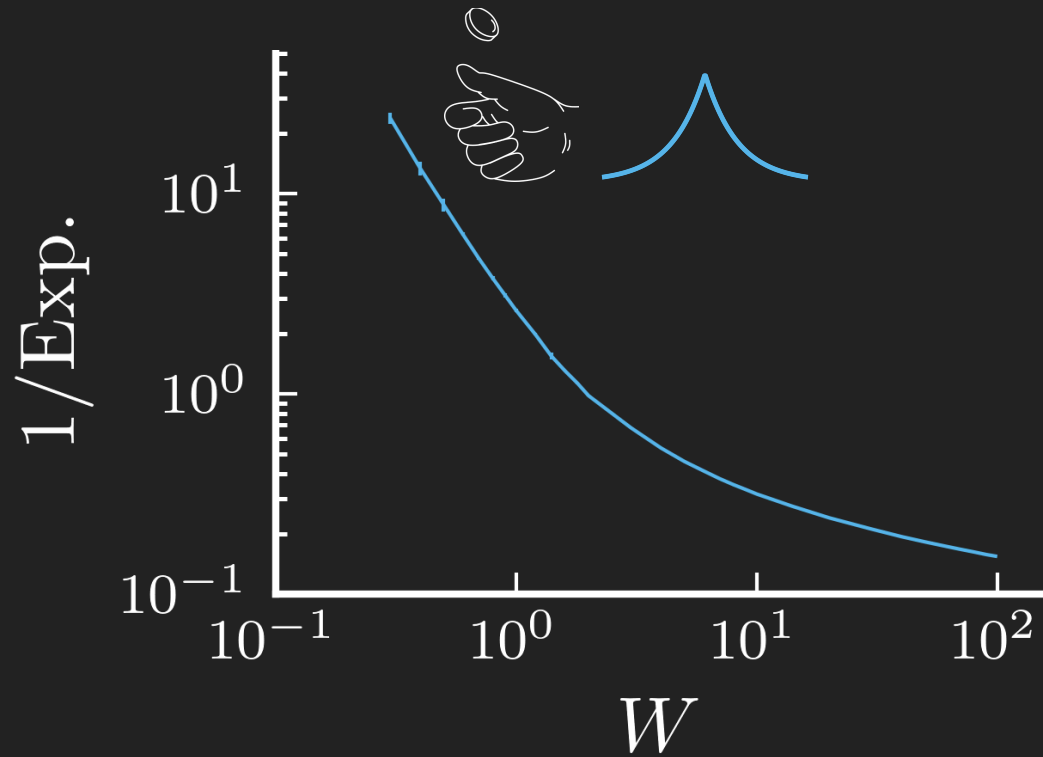
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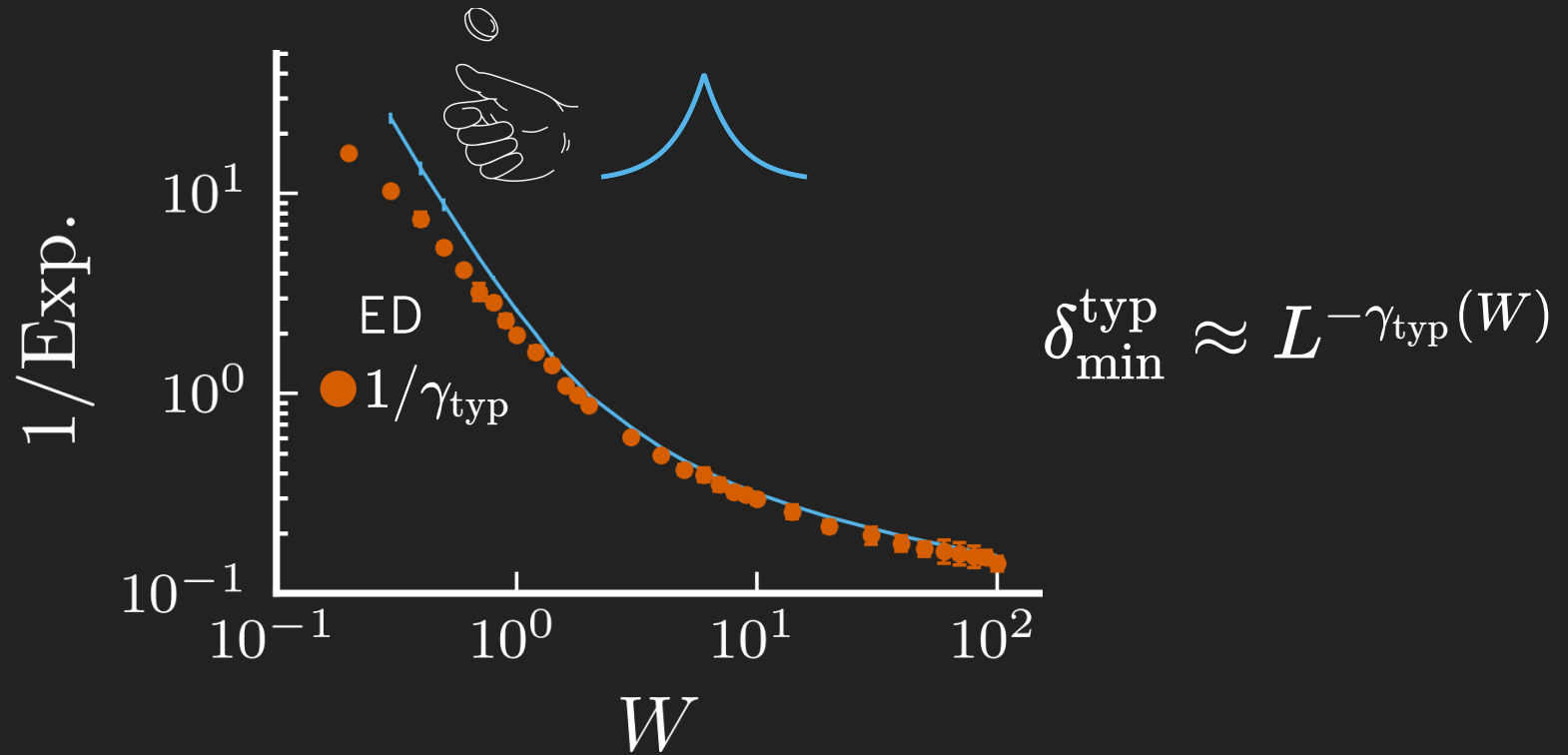
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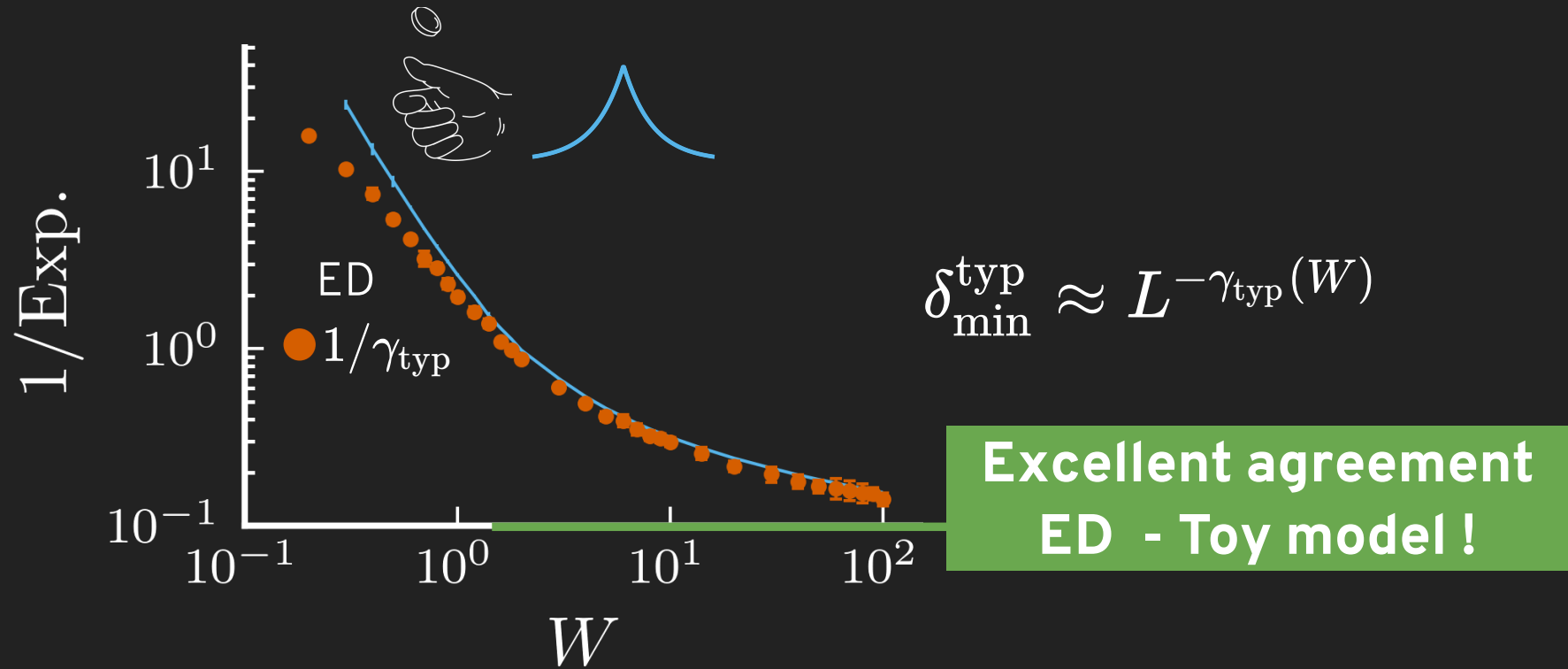
EXPONENTS: XX CHAIN



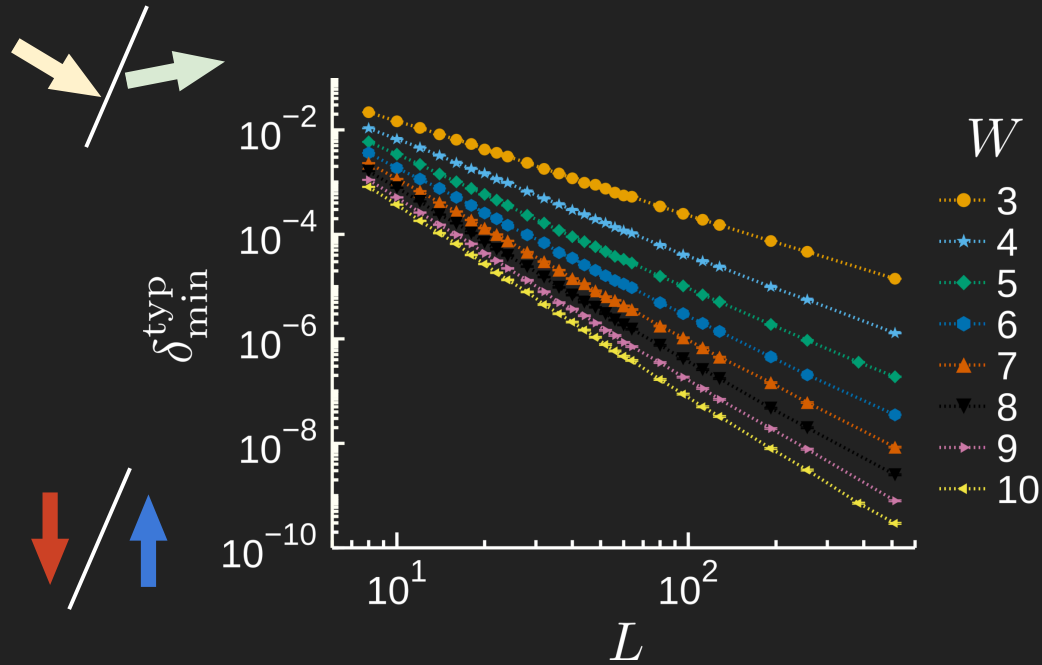
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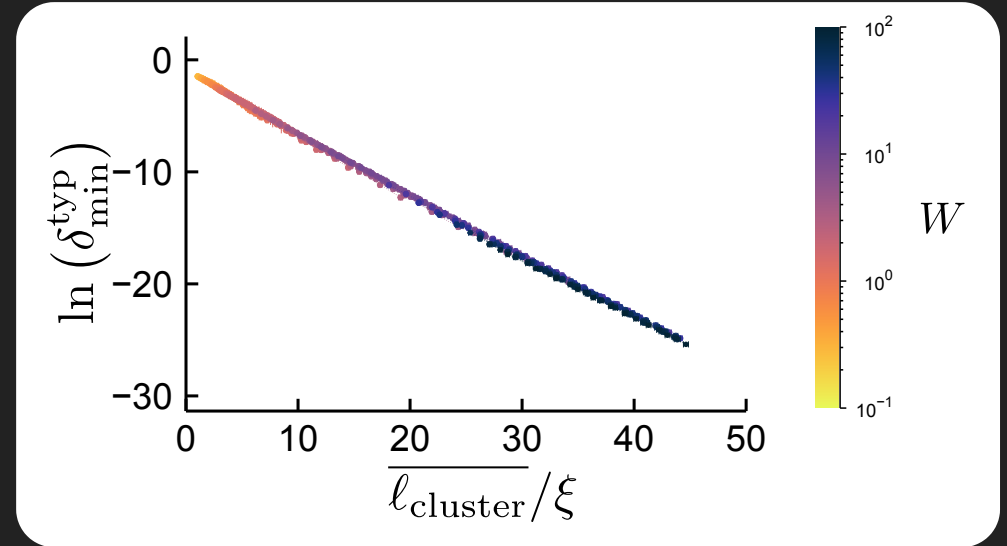
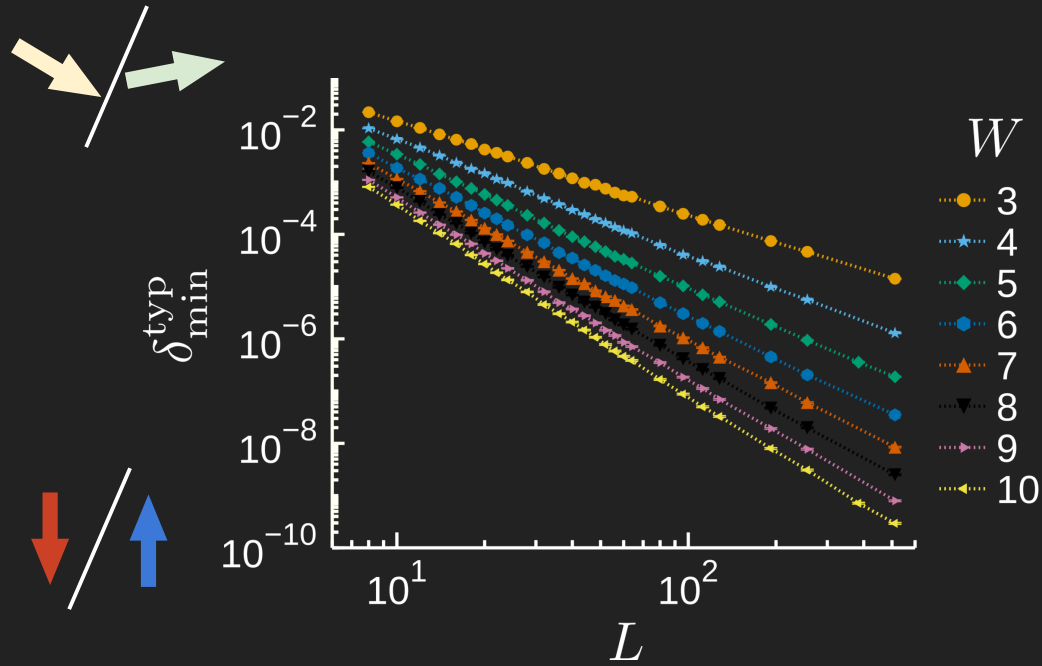


CHAIN BREAKING



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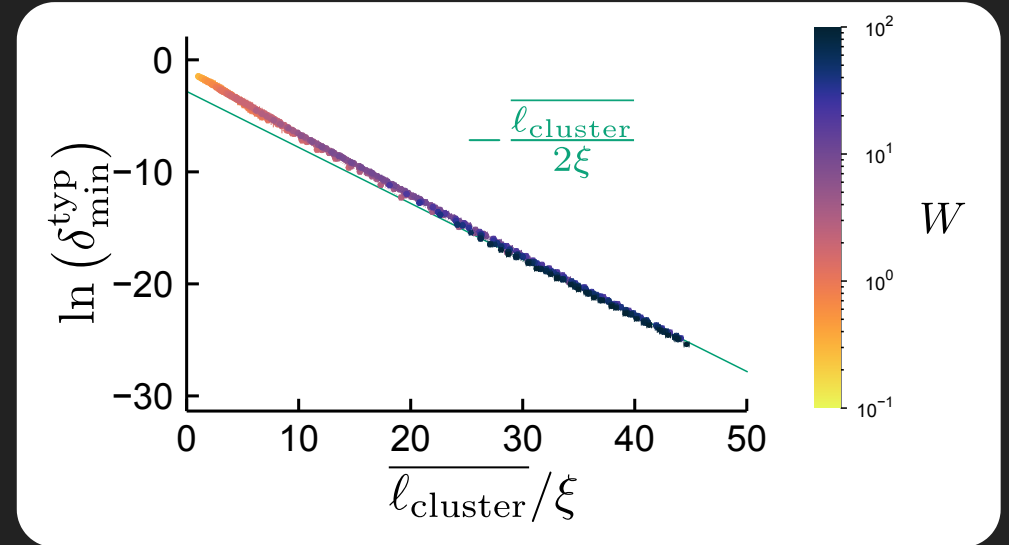
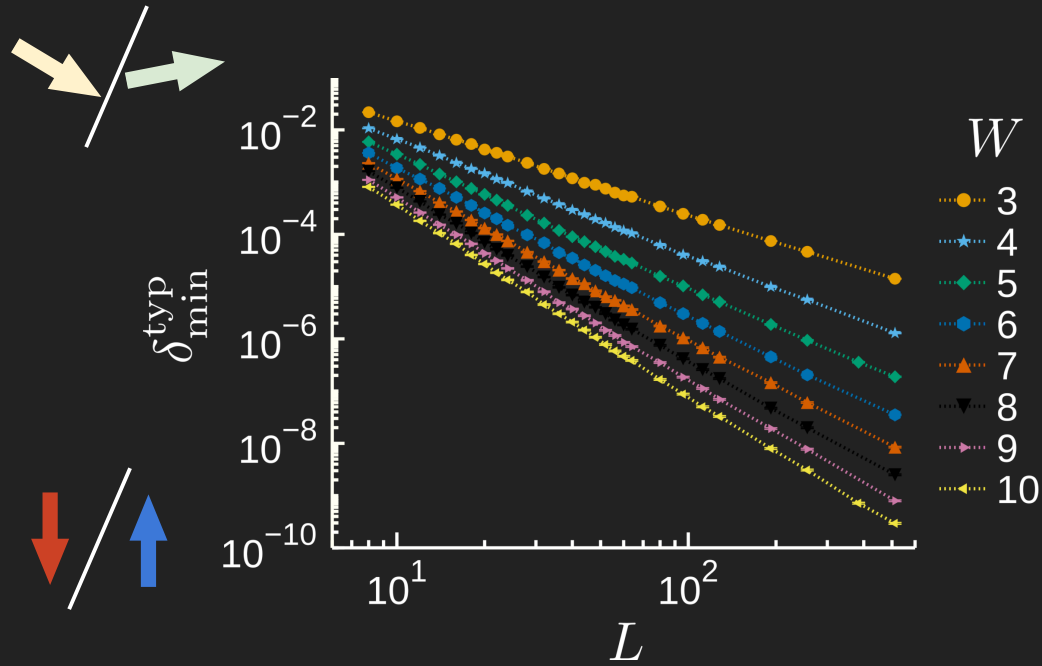
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$$L \gg \xi$$

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POWER-LAW TAILS?

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$$\Rightarrow \mathcal{P}_L(\ln(\delta)) \propto \exp(2\xi \ln 2 \times \ln(\delta))$$

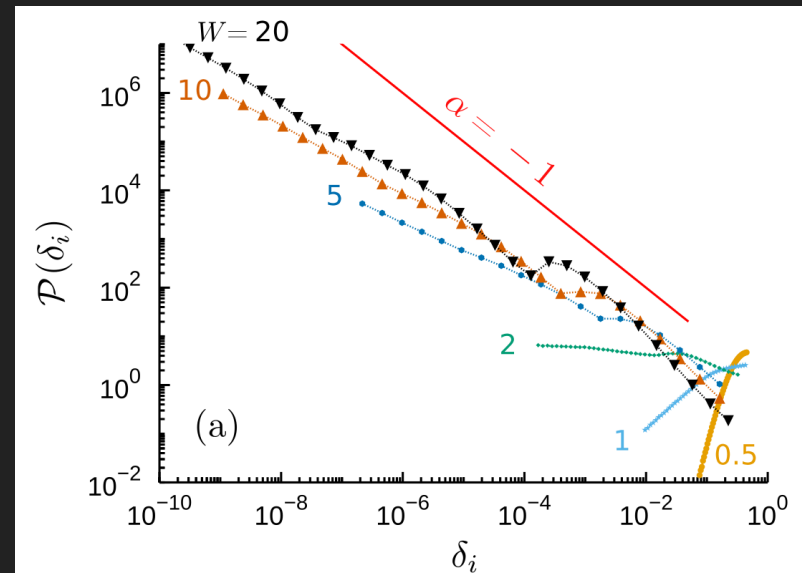
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QUANTITATIVE DESCRIPTION: EXTREME VALUE THEORY

TAILS AND EXTREMES

TAILS AND EXTREMES

$$\{X_i\}_{i=1,2,\dots,L} \sim p(x) \longrightarrow Y = \max(X_i)$$



E. J. Gumbel, *Statistics of Extremes*, Dover, (1958, 2004)

S. N. Majumdar, A. Pal, G. Schehr, *Physics Reports*, **840**, 1 (2020)

TAILS AND EXTREMES

Tails

$$\{X_i\}_{i=1,2,\dots,L} \sim p(x) \longrightarrow Y = \max(X_i)$$



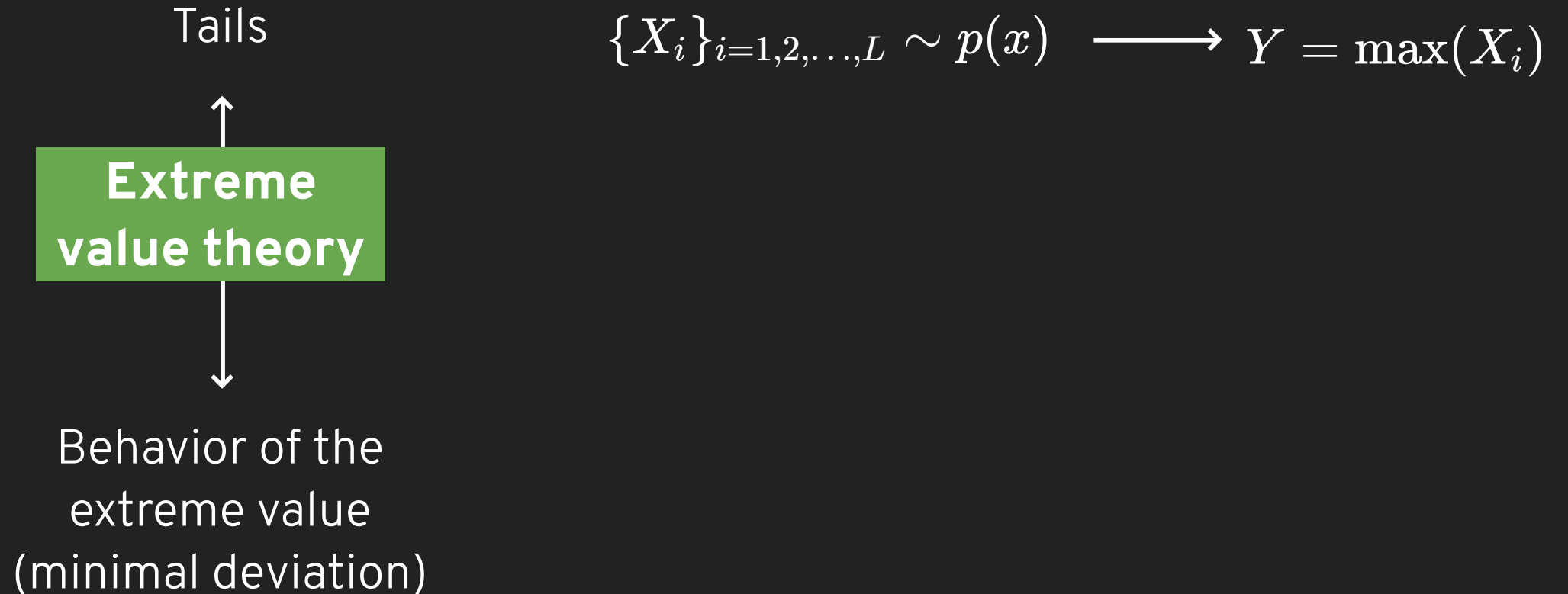
Behavior of the
extreme value
(minimal deviation)



E. J. Gumbel, *Statistics of Extremes*, Dover, (1958, 2004)

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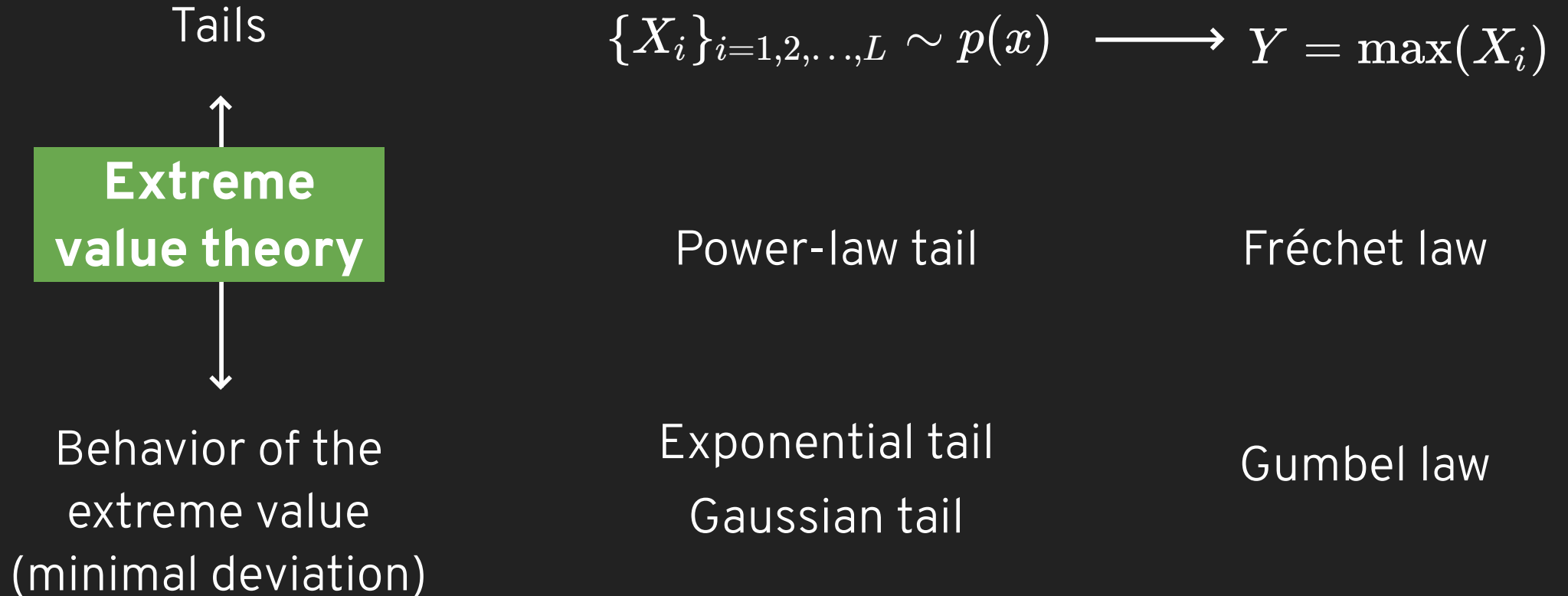
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TAILS AND EXTREMES



EXTREME VALUE THEORY - XX CHAIN

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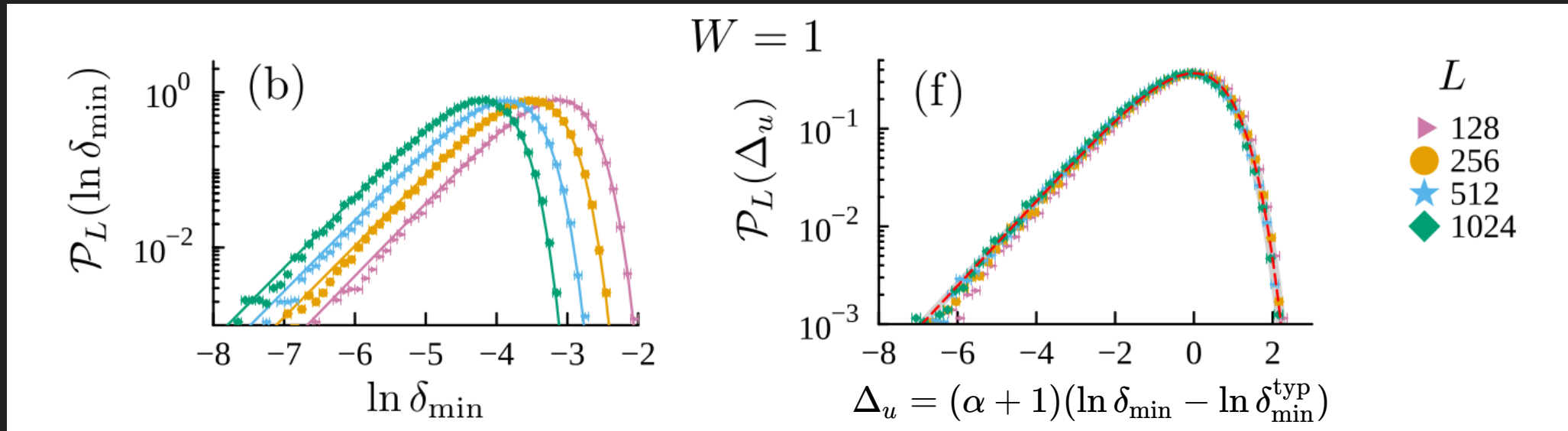
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10^5 samples



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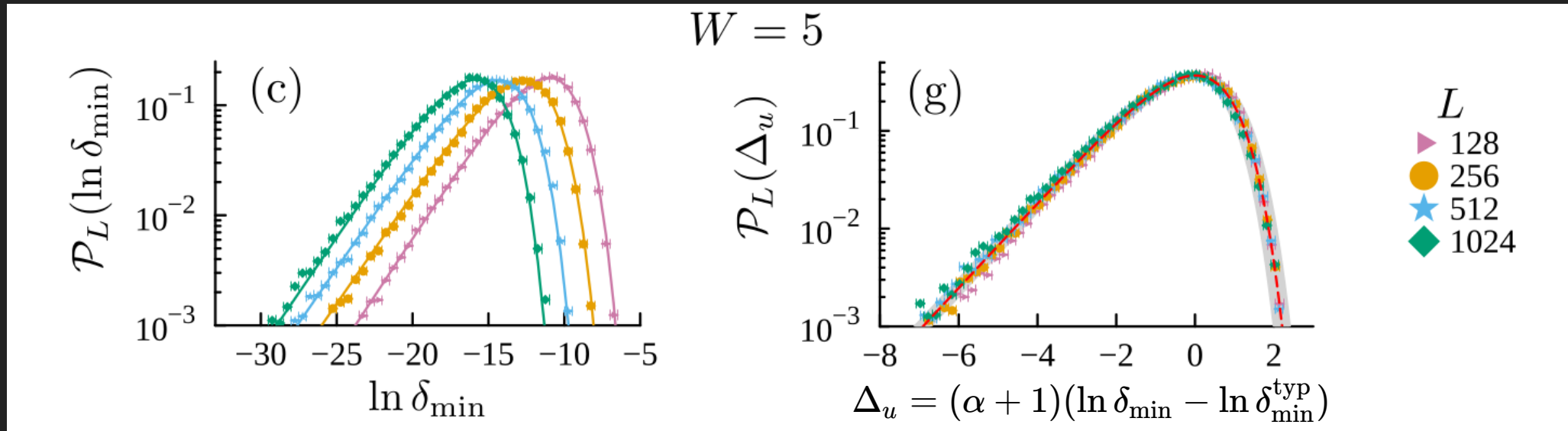
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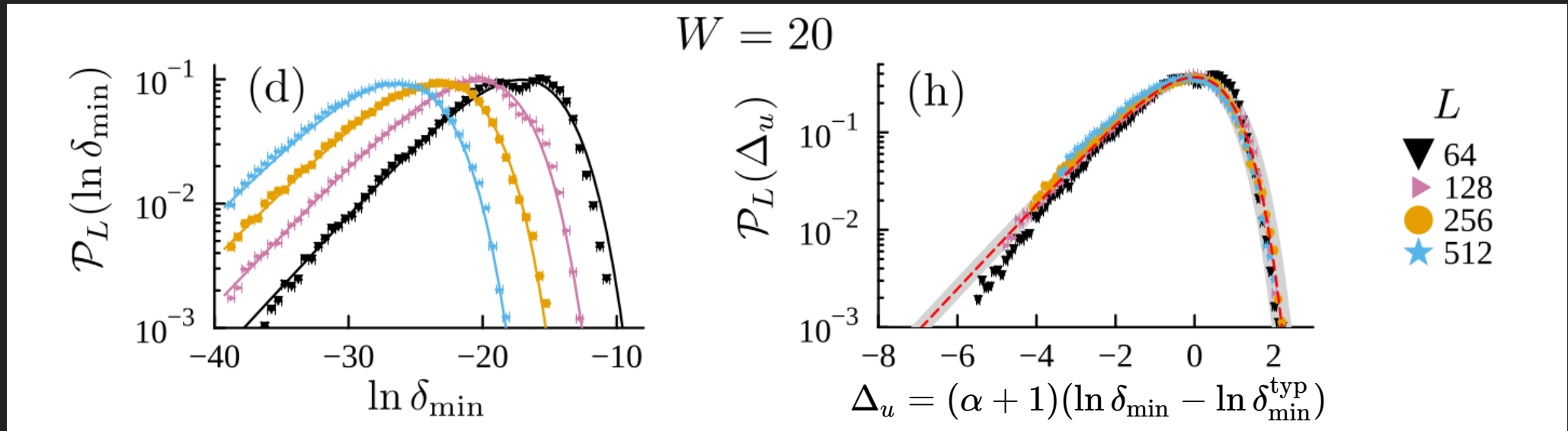
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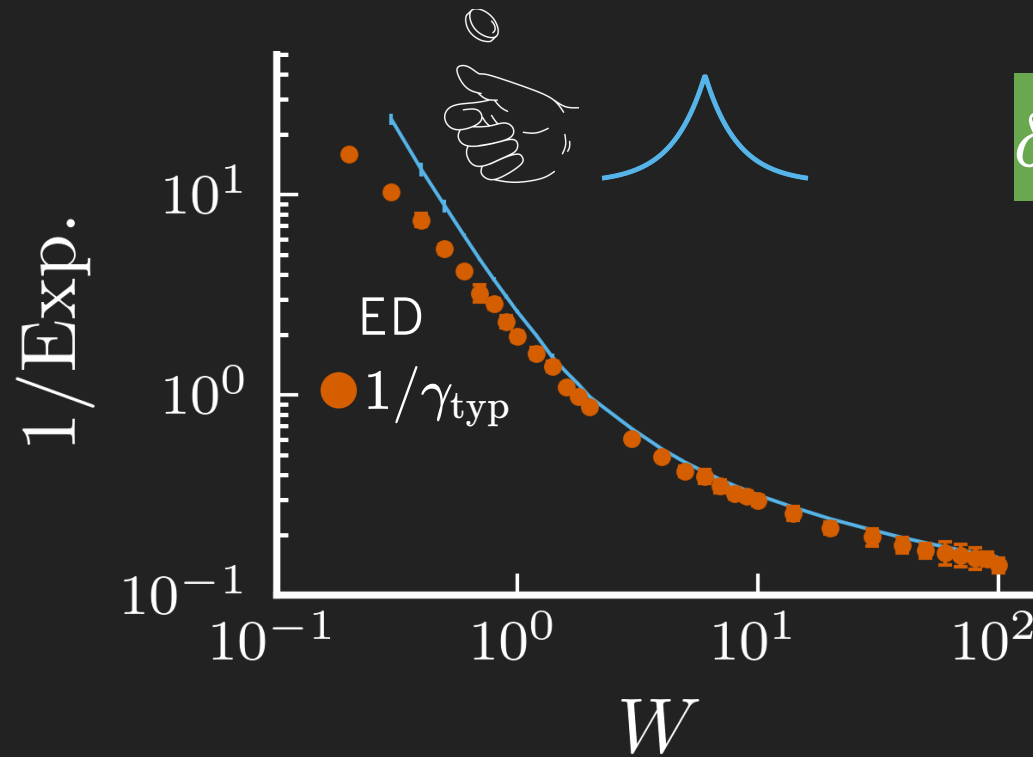


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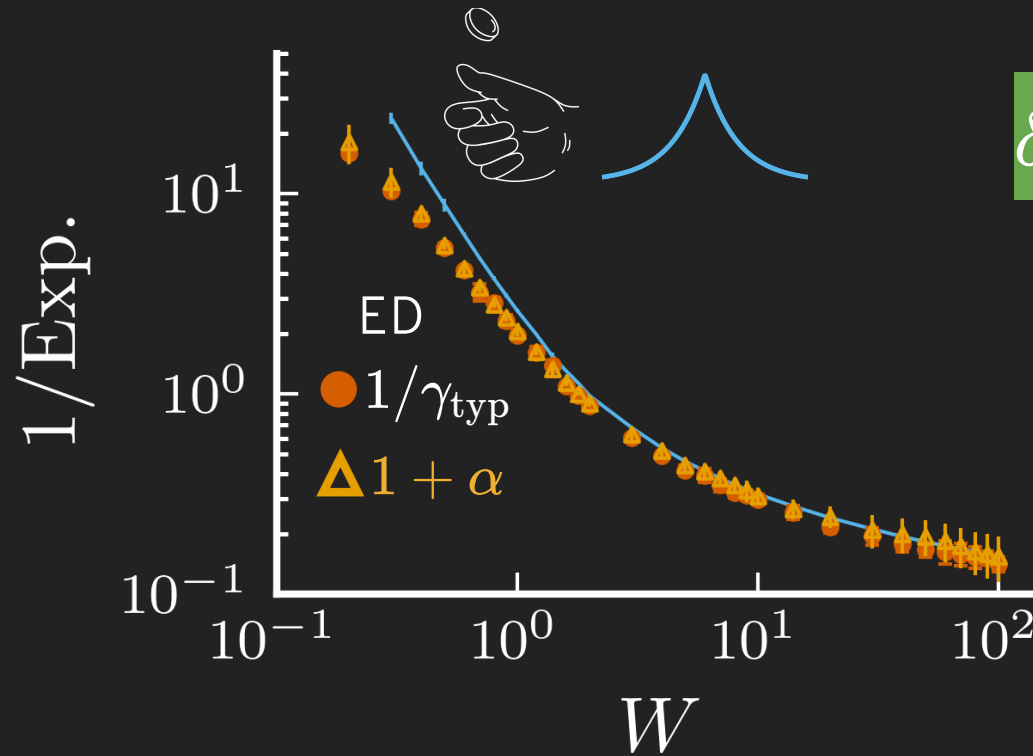
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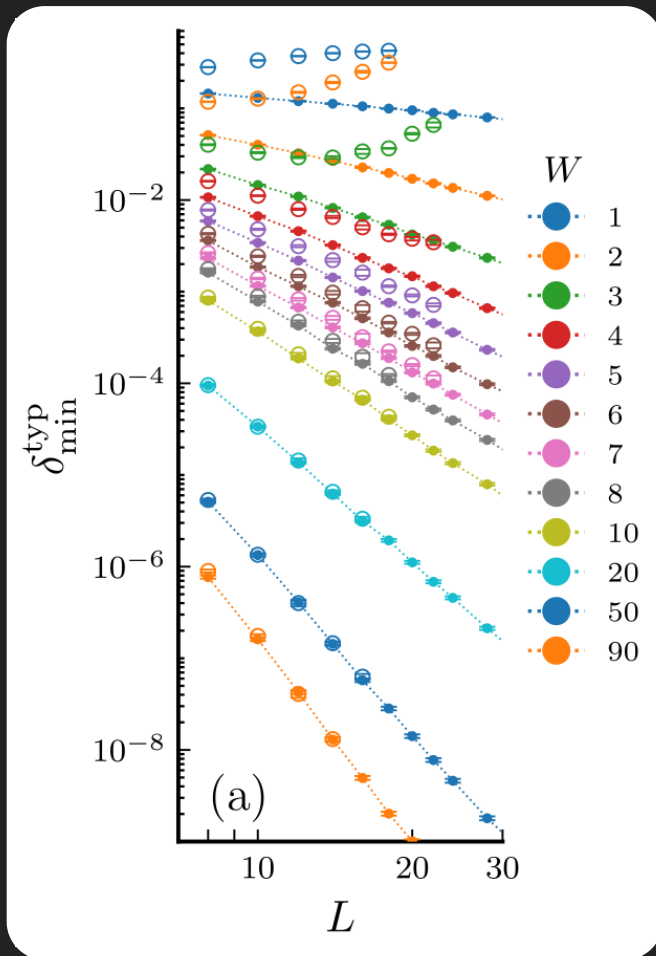
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EFFECT OF INTERACTIONS?

ED: shift-invert, see eg.
F. Pietracaprina, N. Macé, D. J. Luitz, and F. Alet,
SciPost Physics 5, 045 (2018)
Laflorencie, Lemarié, Macé, PRR **2**, 042033(R), (2020)
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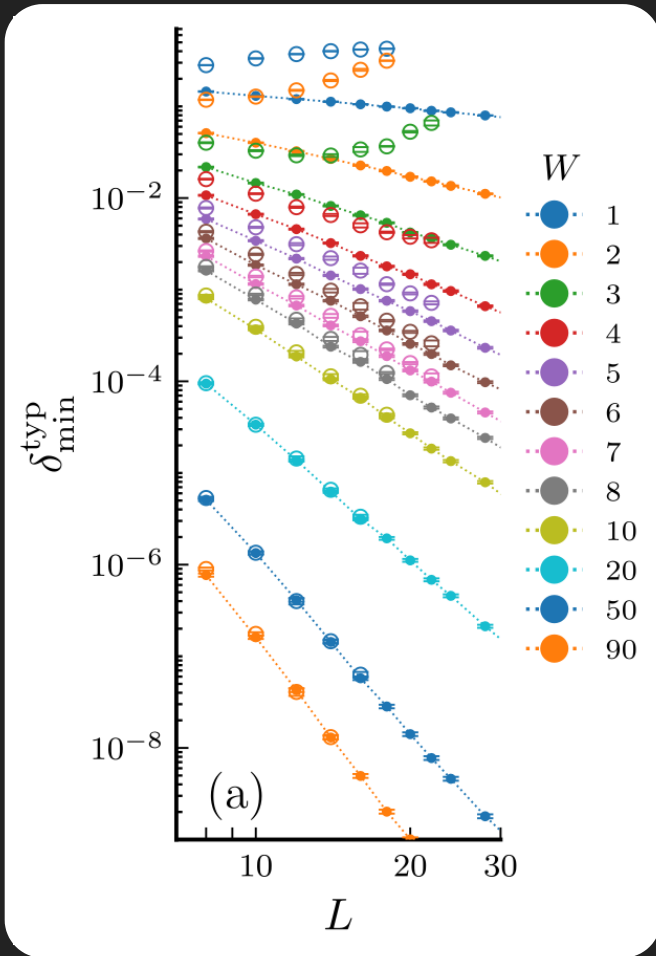
Challenge : high-energy eigenstate, lack of symmetries. Shift-invert ED $\rightarrow L = 22$



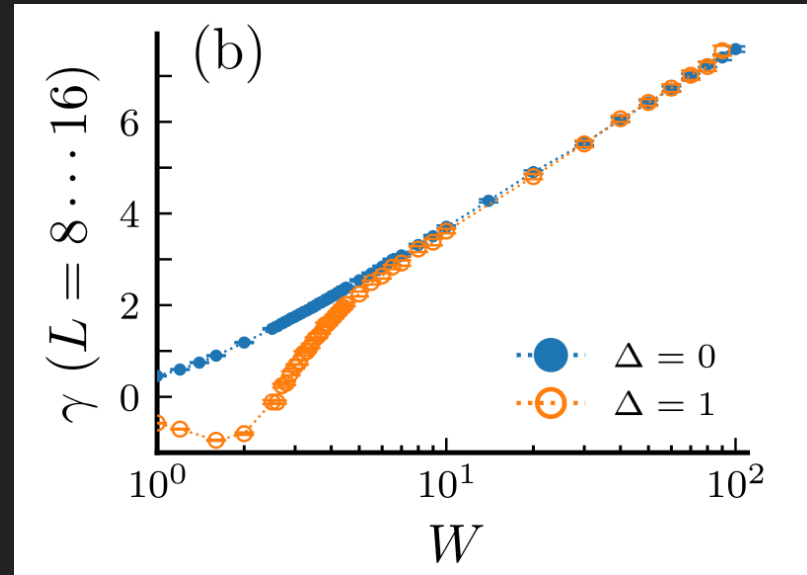
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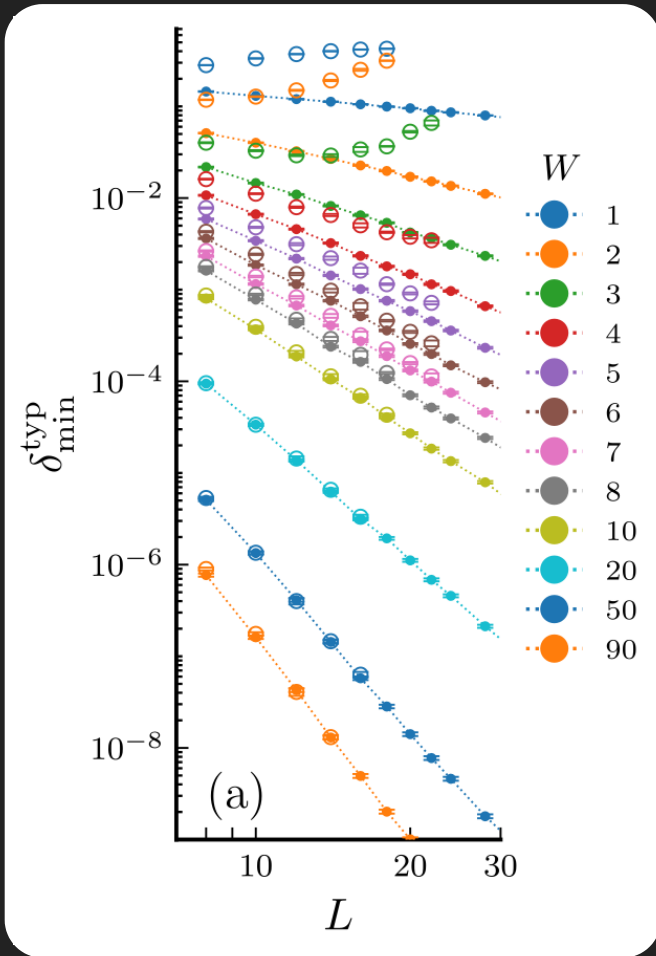
↓
disorder
increases



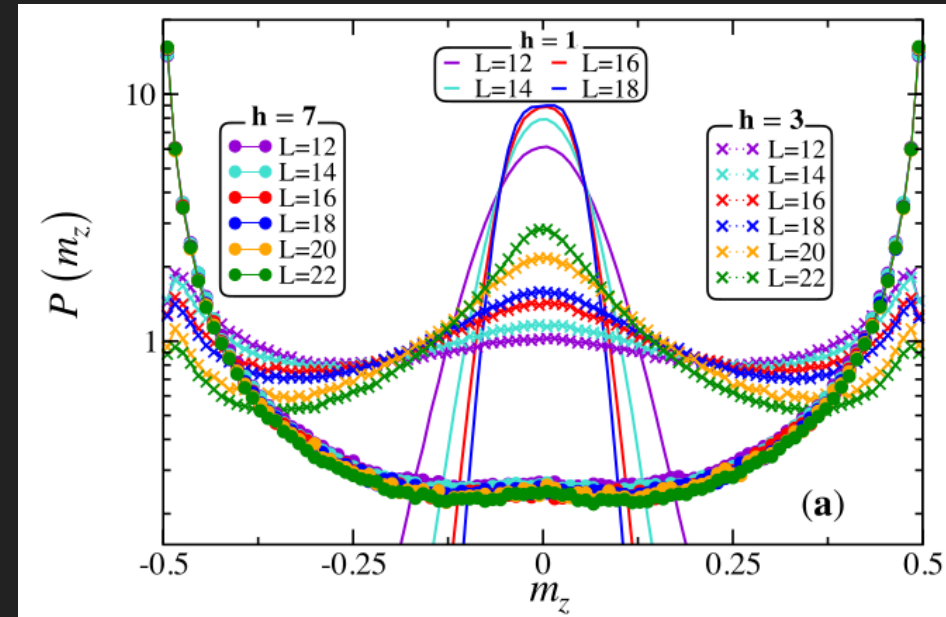
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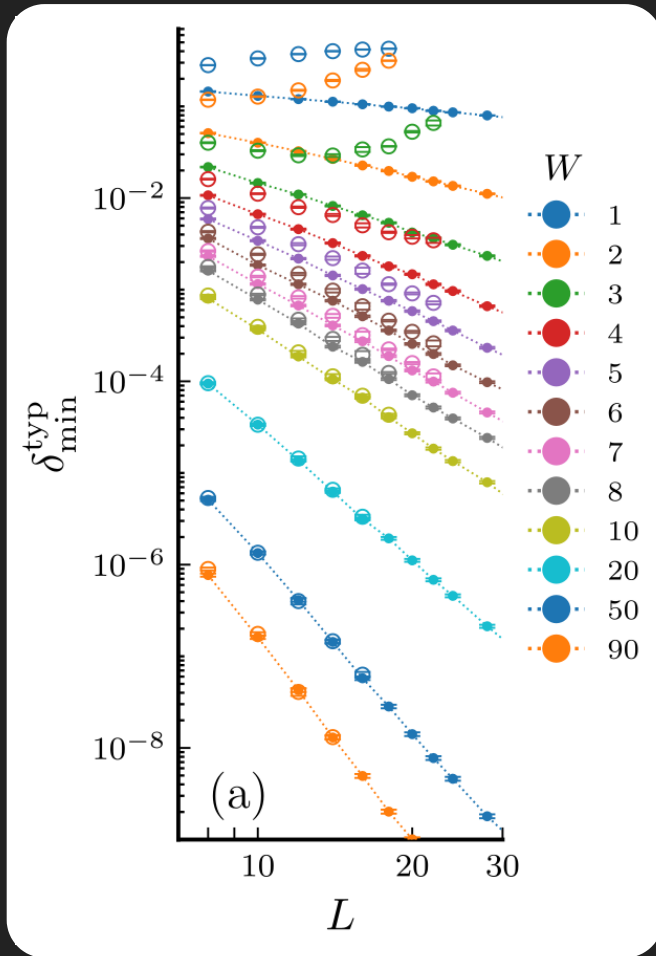
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Two limits captured

Can we say more?

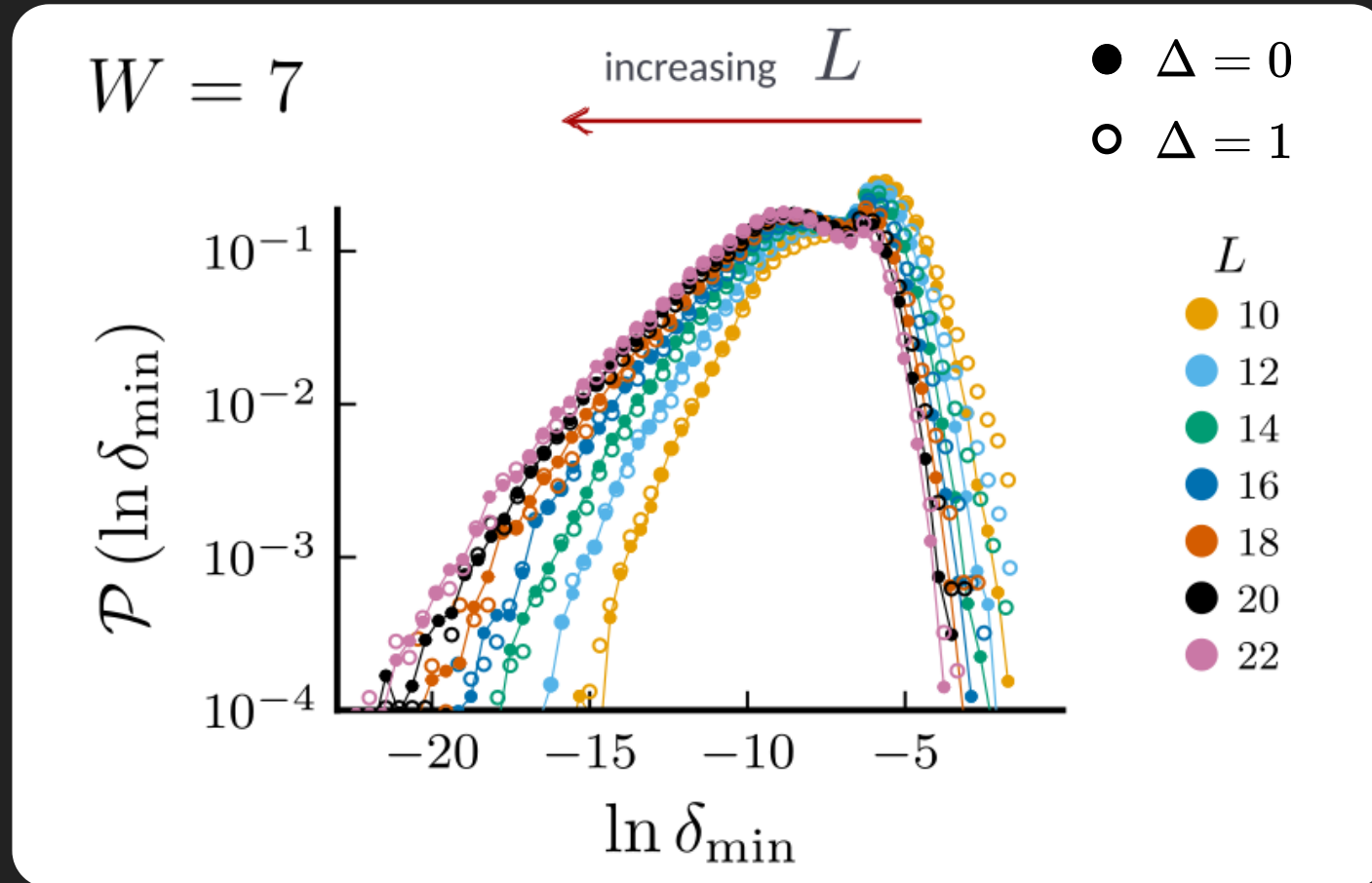
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CONSEQUENCES : HEISENBERG

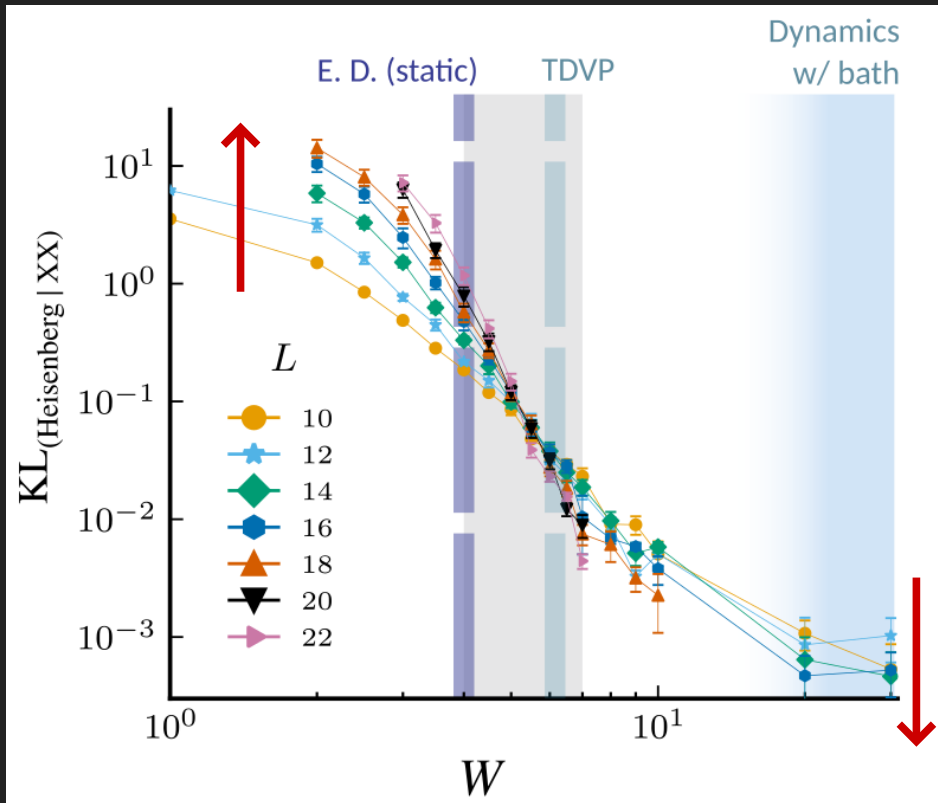
Conjecture : Gumbel (?) on the Ergodic side, Fréchet on the MBL side.

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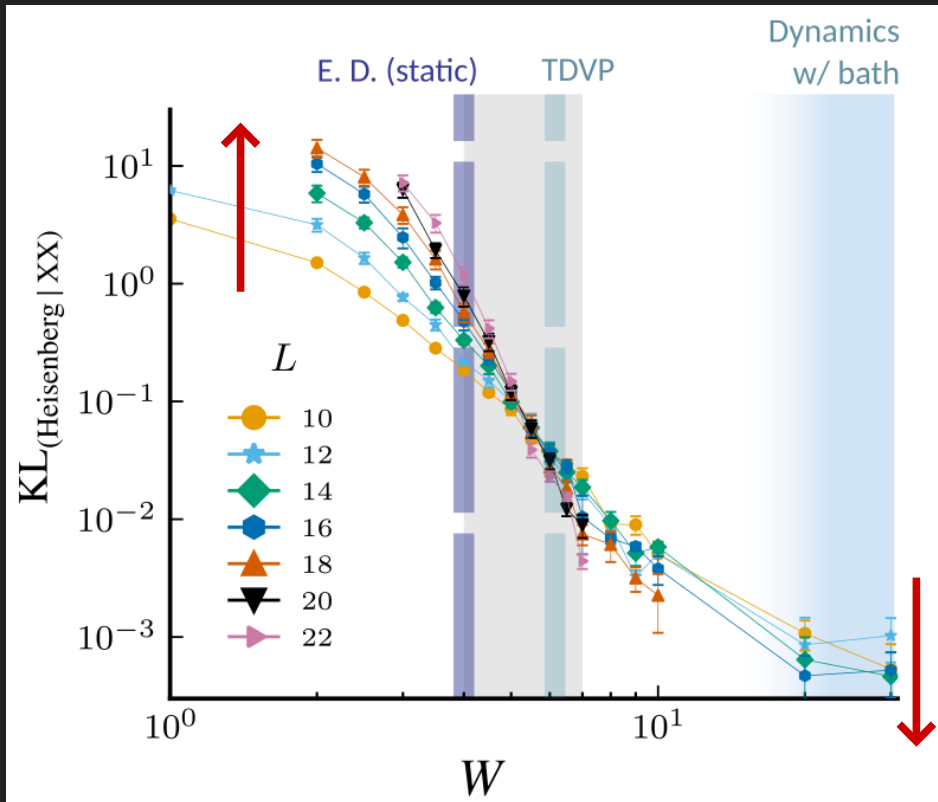
Kullback-Leibler divergence :

$$KL(p|q) = \sum_i q_i \ln \frac{q_i}{p_i}$$

S. Kullback and R. A. Leibler, The annals of mathematical statistics **22**, 79 (1951)
 JC, N. Laflorencie, arXiv:2305.10574

E. H. V. Doggen et al., PRB 98, 174202 (2018)
 See e.g. D. Sels, PRB 106, L020202 (2022)
 JC, N. Laflorencie, arXiv:2305.10574

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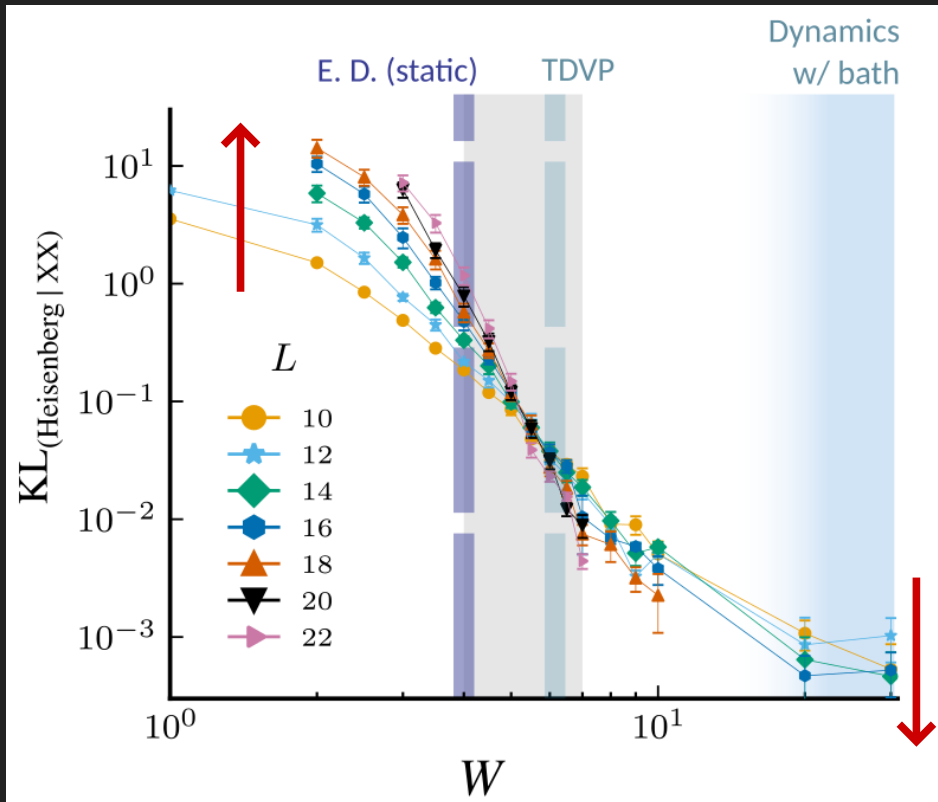


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- Transition in the extreme value distributions

CONSEQUENCES : HEISENBERG

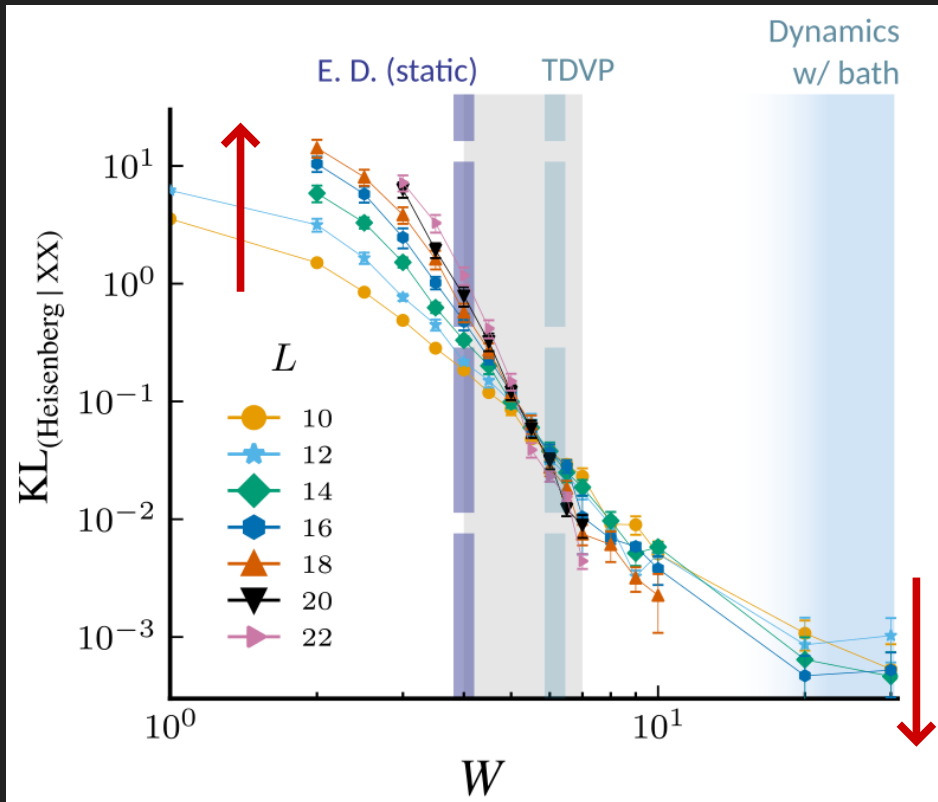


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- Coinciding with the MBL transition?

CONSEQUENCES : HEISENBERG



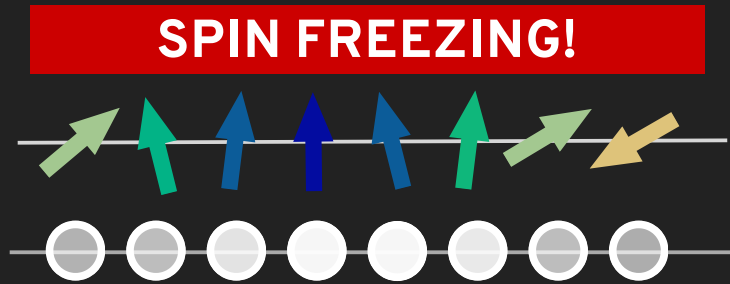
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- Transition in the extreme value distributions
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- Chain breaks as a tool?

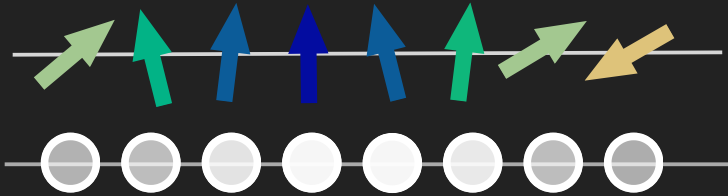
TAKE HOME MESSAGE

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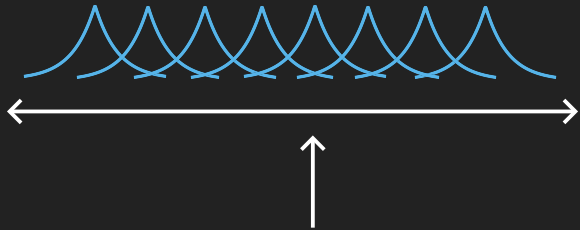
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SPIN FREEZING!

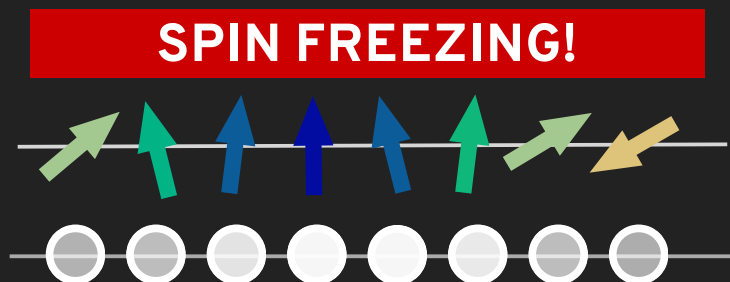


XX chain :

- controlled by largest cluster of occupied orbitals

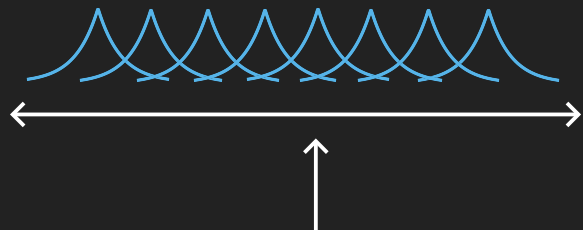


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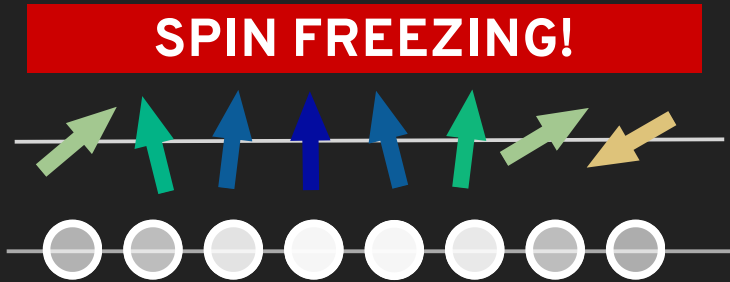
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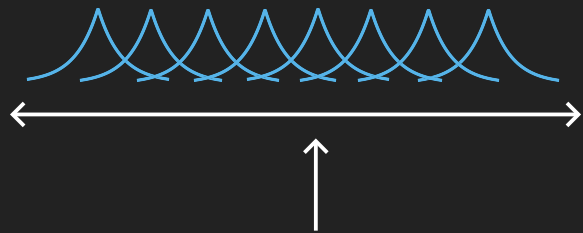
- Excellent fits & collapses with a Fréchet Law

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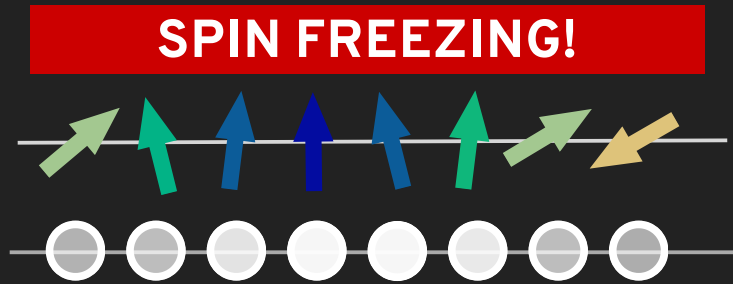


- Excellent fits & collapses with a Fréchet Law

Heisenberg chain at strong disorder:
Chain breaks!

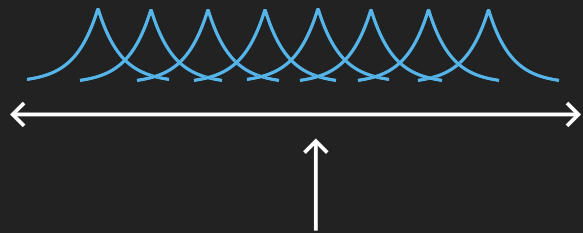


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Heisenberg chain at strong disorder:

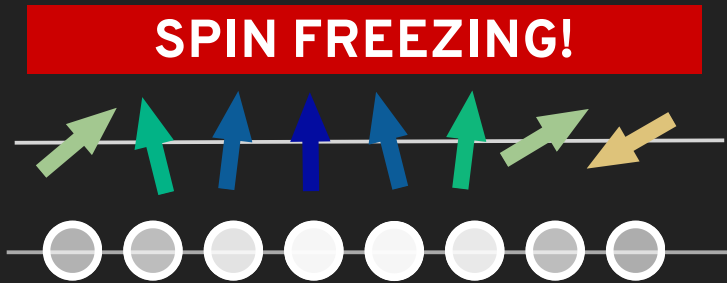
Chain breaks!



Comparing δ_{\min} deviation in Heisenberg vs Many-body Anderson:

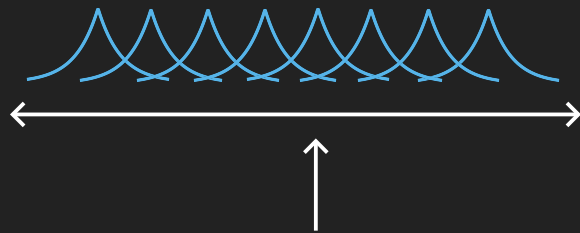
Extreme value transition characterized by the **KL divergence.**

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Comparing δ_{\min} deviation in Heisenberg vs Many-body Anderson:

Extreme value transition characterized by the **KL divergence.**

Thank you!

arXiv:2305.10574

