

# Parton cascades at DLA: the role of the evolution variable

André Cordeiro

In collaboration with:

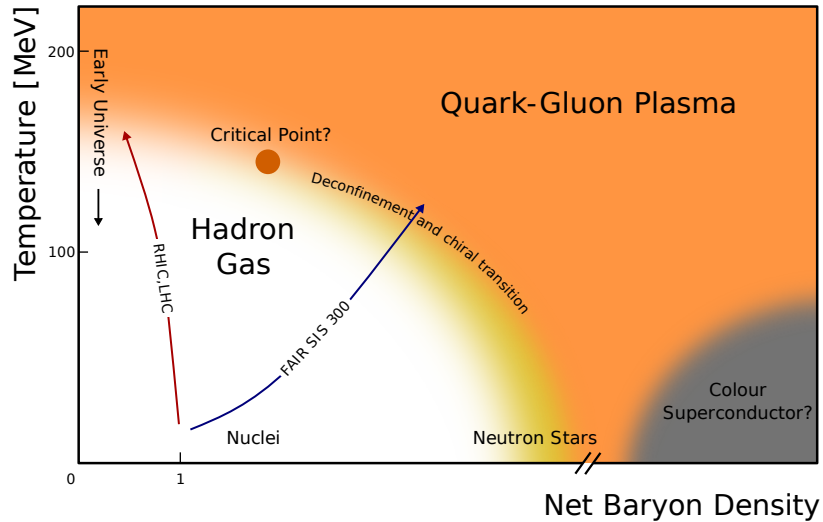
Carlota Andrés, Liliana Apolinário, Nestor Armesto,  
Fabio Dominguez, Guilherme Milhano



TÉCNICO  
LISBOA

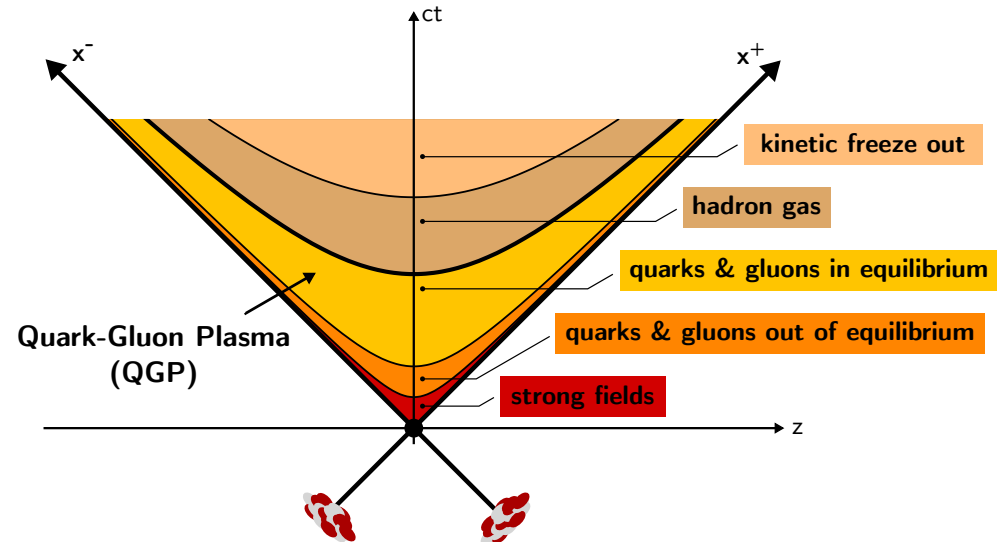
Excited QCD Workshop, 19<sup>th</sup> January 2024

# Heavy Ion Collisions

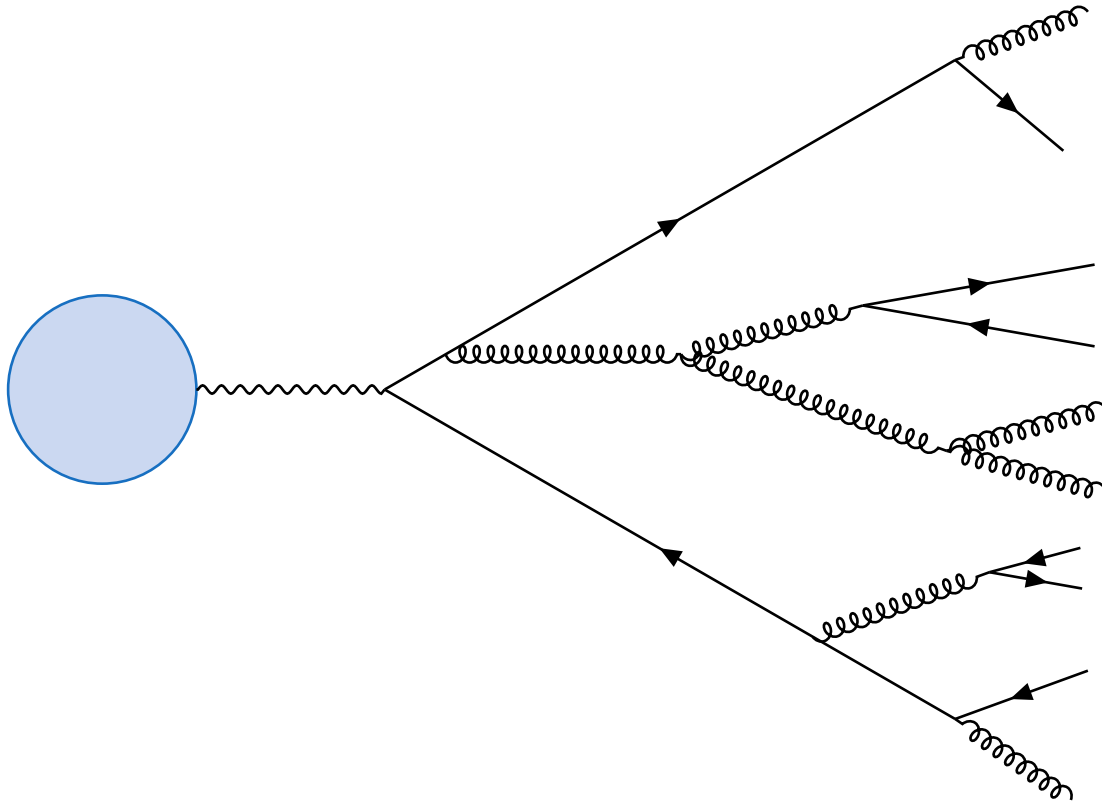


The spacetime evolution of QCD matter covers a wide range of time/energy scales

Heavy Ion collisions are valuable as a laboratory to study the QCD phase diagram

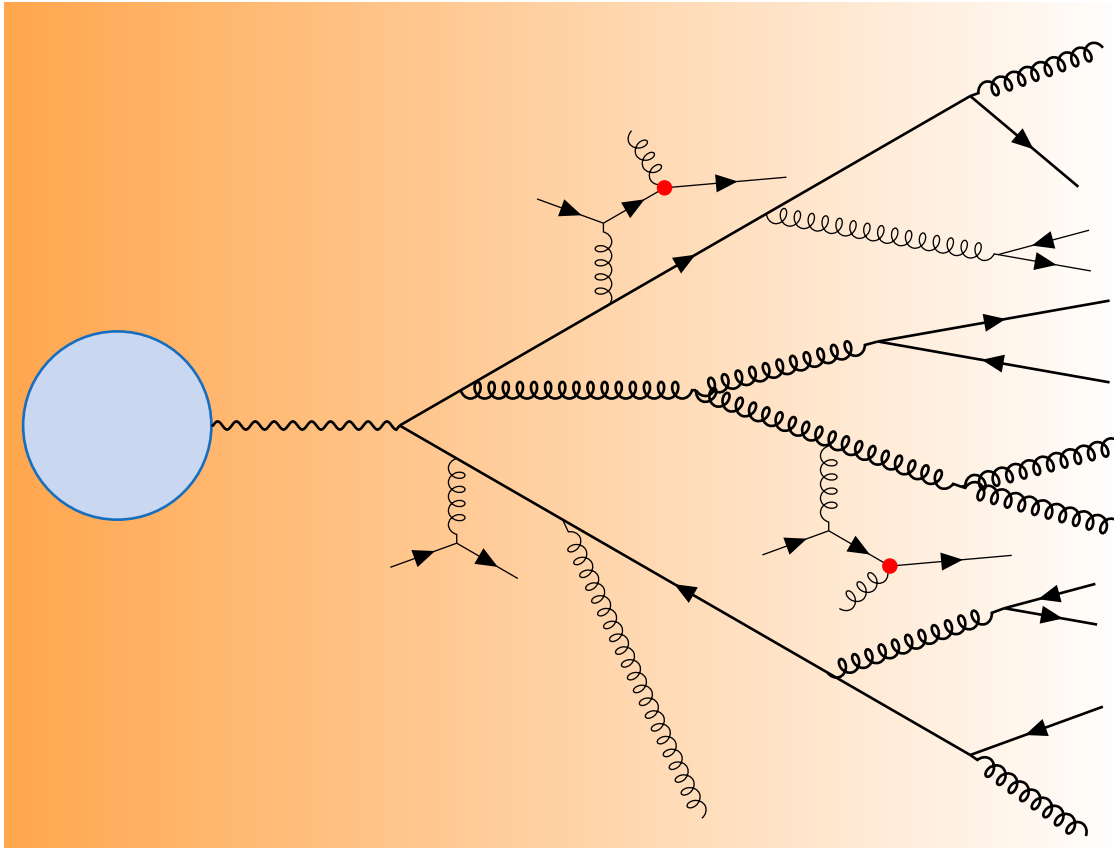


# Parton Showers in a Coloured Medium



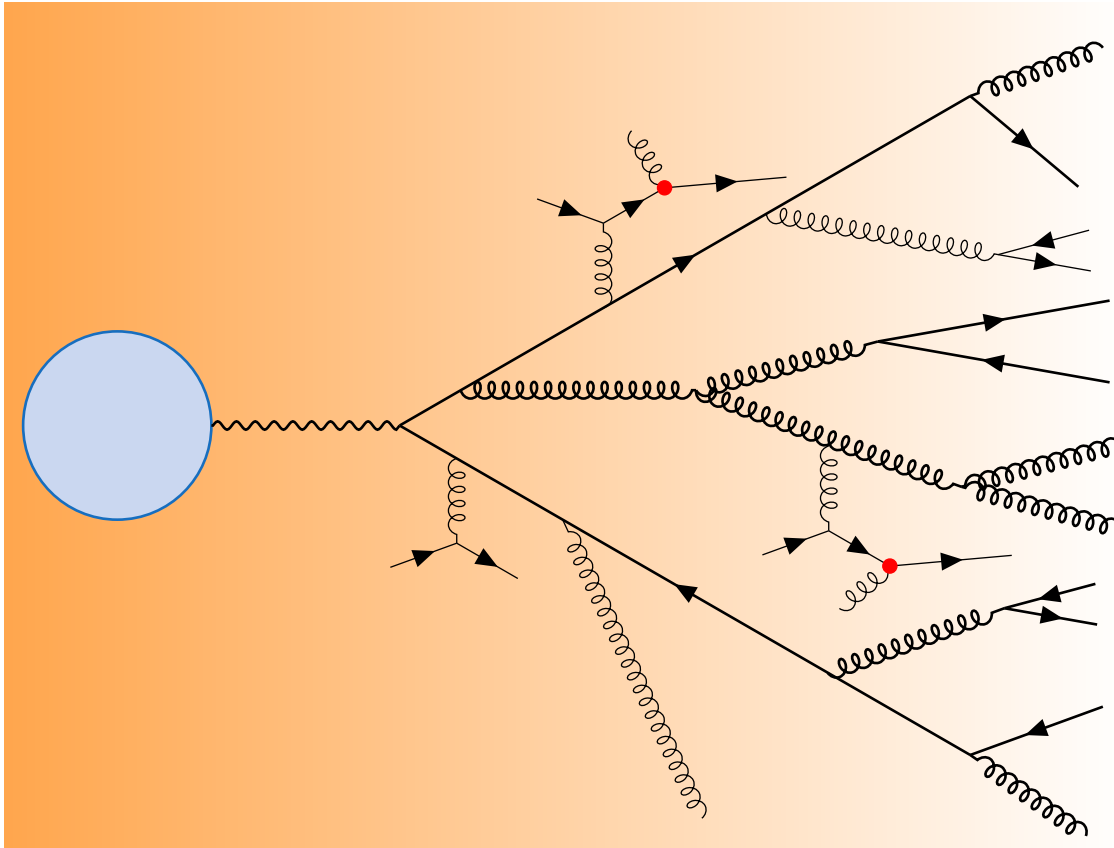
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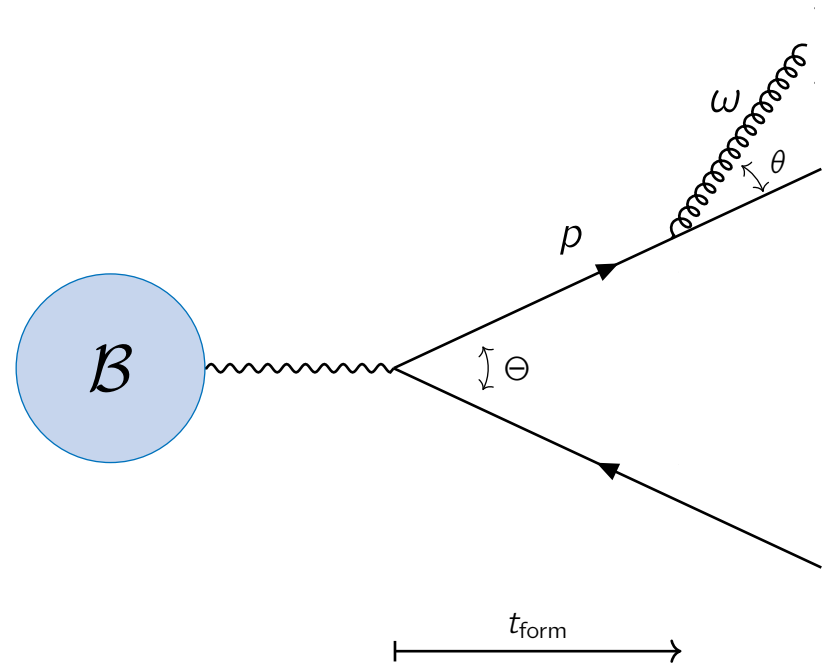
**Is jet quenching sensitive to the ordering of vacuum-like splittings?**

**First, a look at vacuum  
(proton-proton) showers**

# QCD Vacuum Splittings

Estimate some scales:

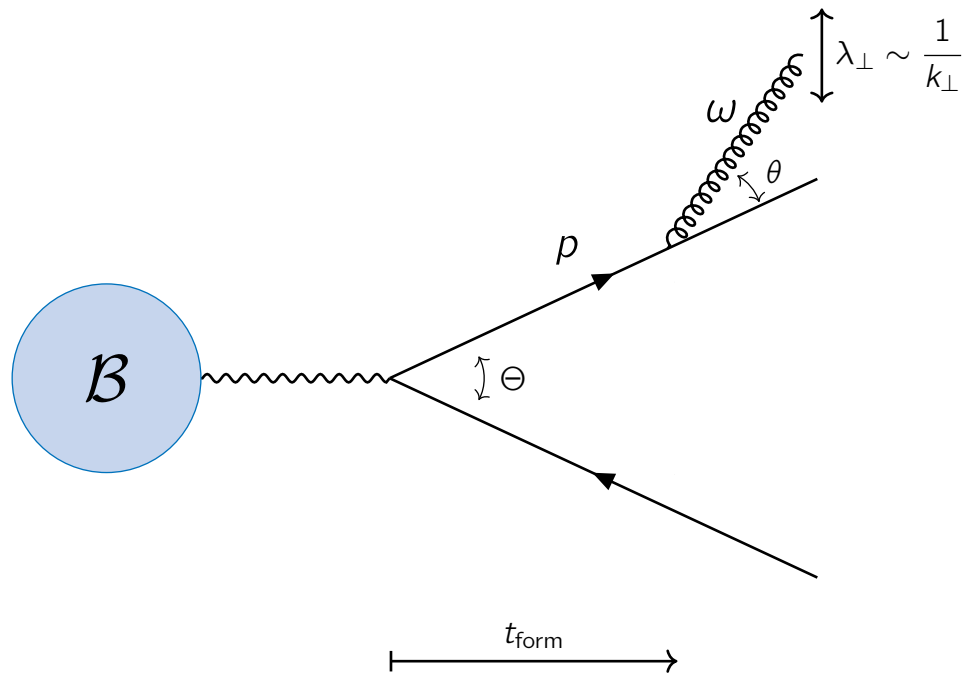
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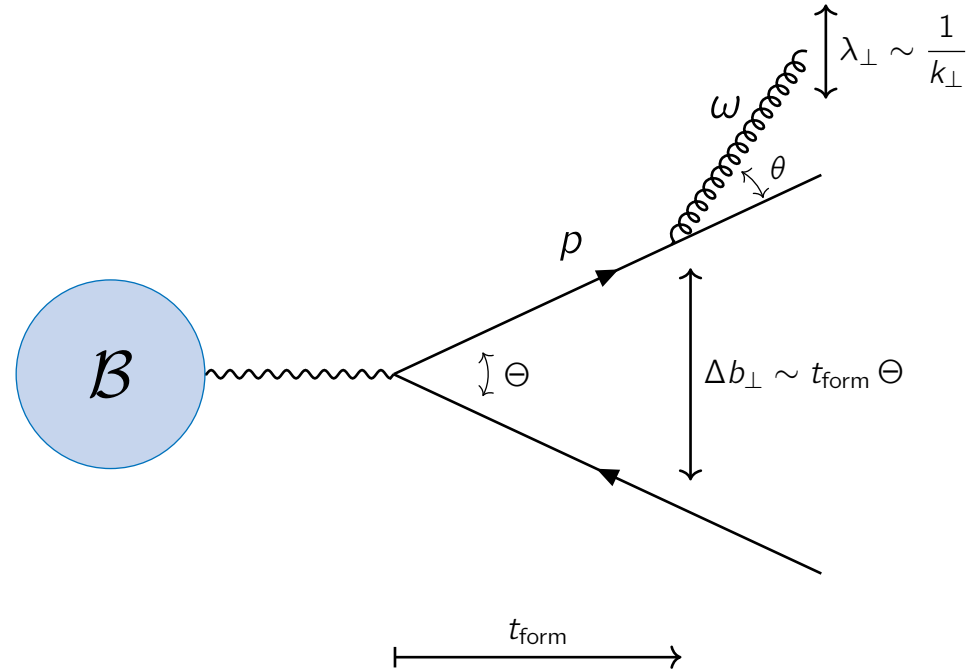




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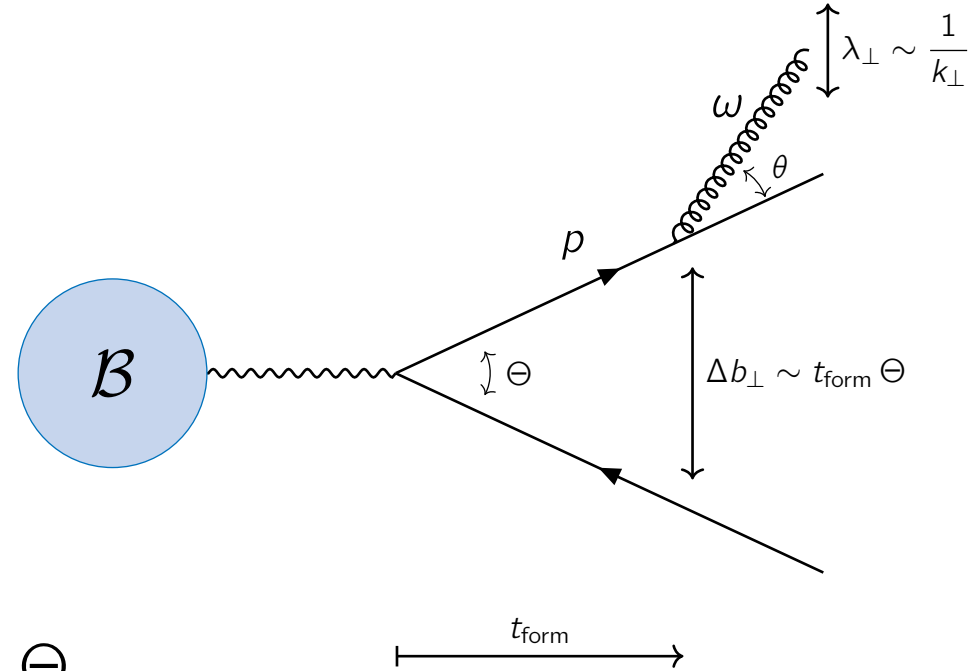
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**For an antenna:**  $\lambda_{\perp} < \Delta b_{\perp} \iff \theta < \Theta$

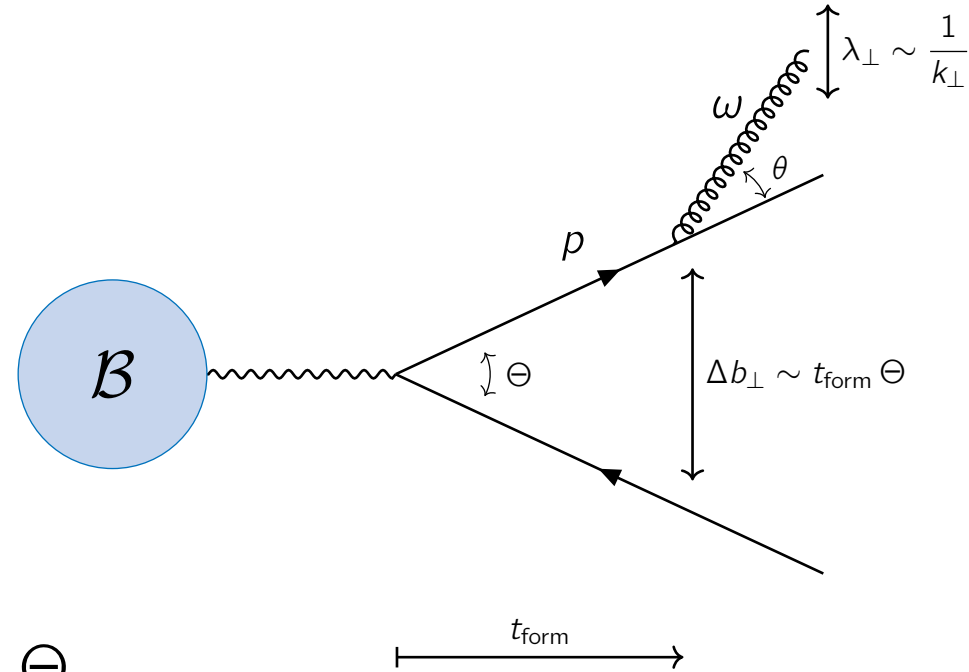
\* Larger  $\lambda_{\text{T}}$

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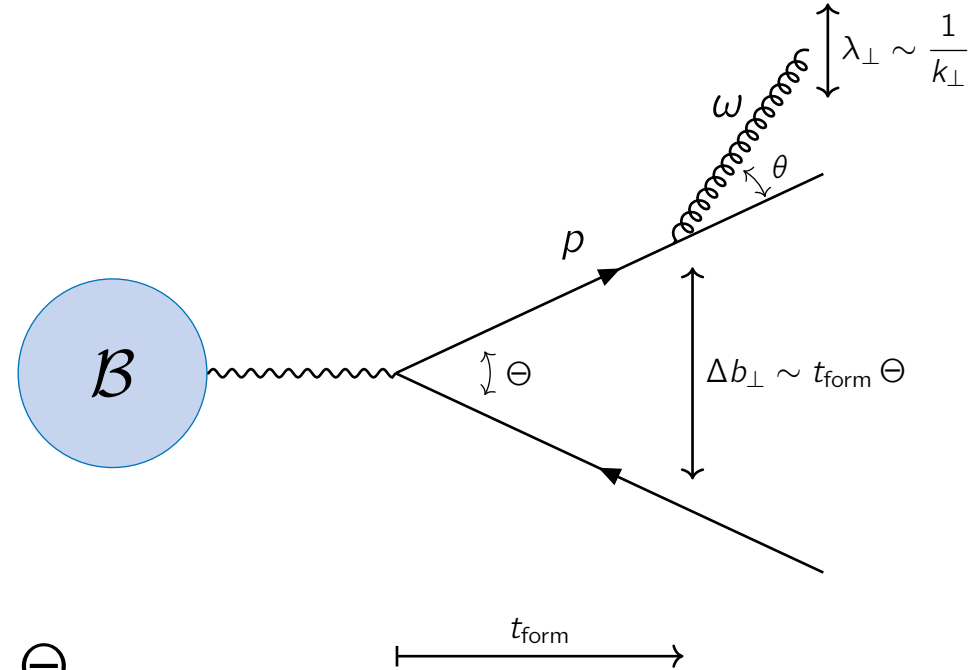
This is the angular ordering property of vacuum splittings → Showers are collimated

# QCD Vacuum Splittings

\*Can be generalised to non-singlets & multiple emissions

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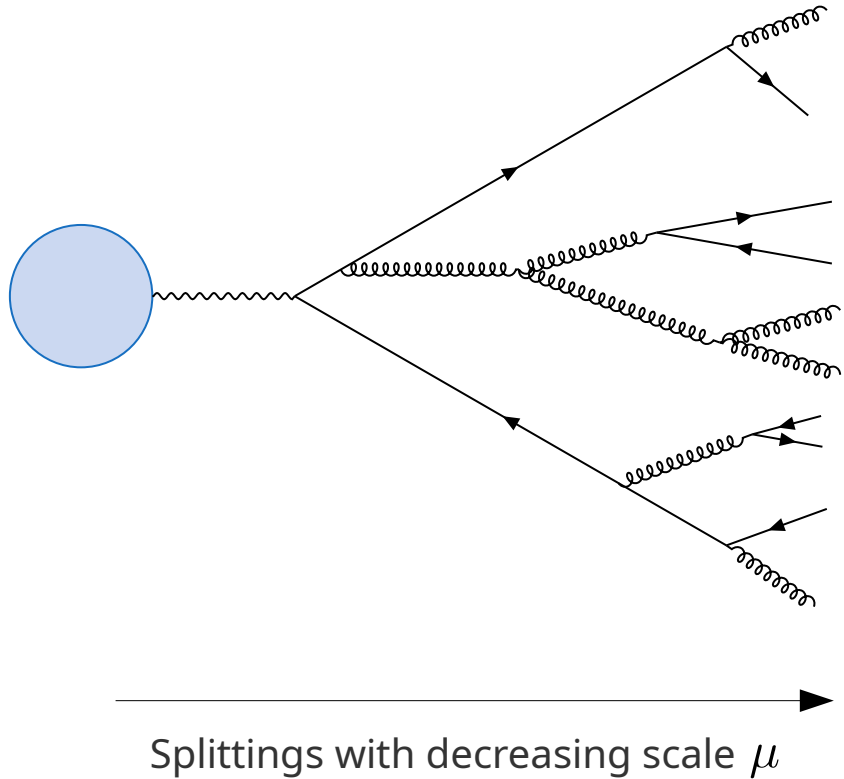
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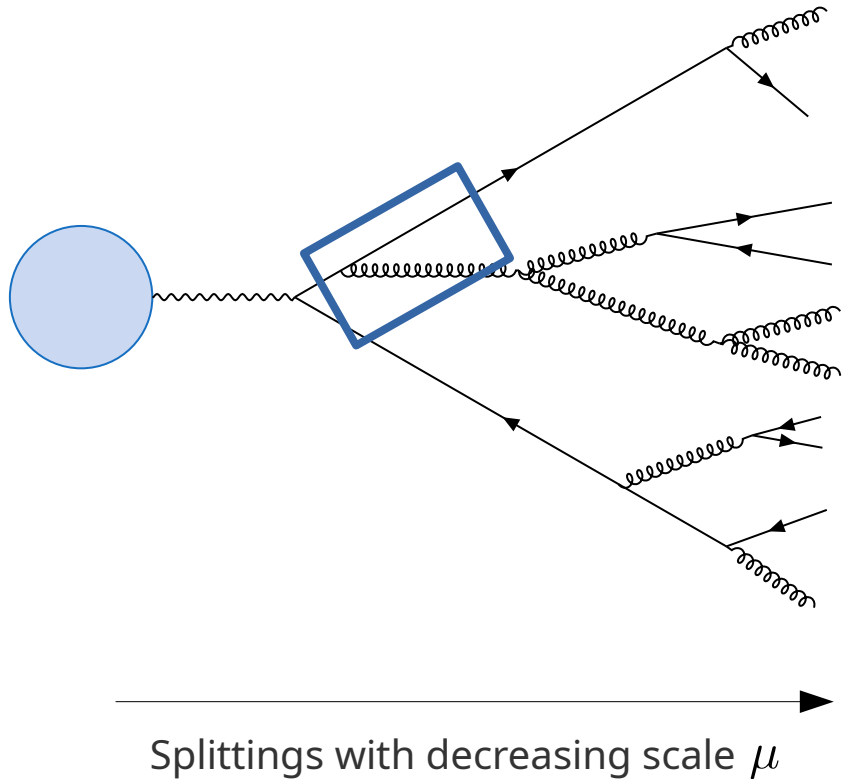
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# How to build a parton shower



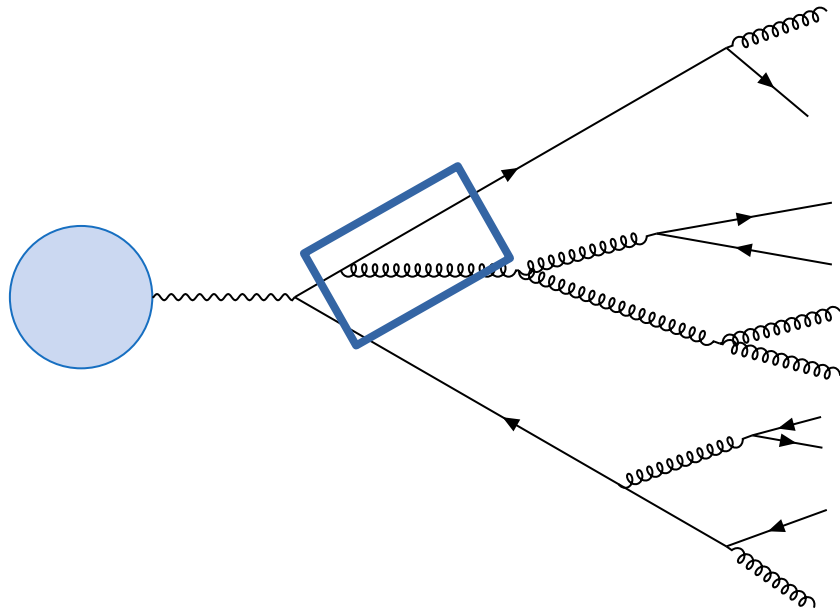
# How to build a parton shower

Building blocks: QCD splittings

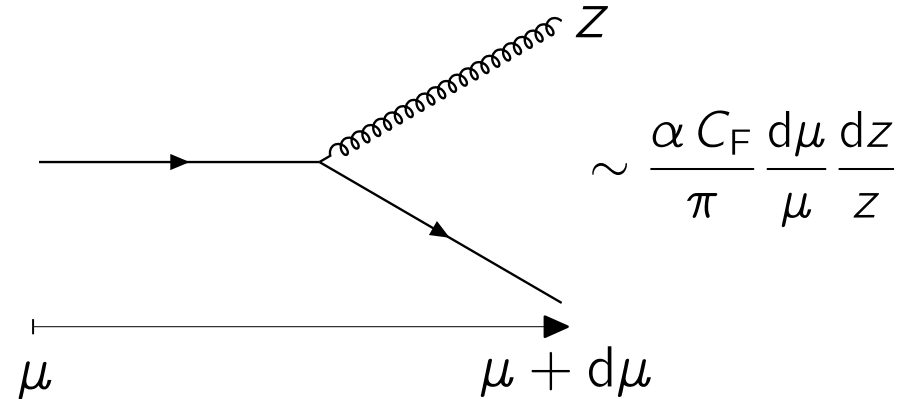


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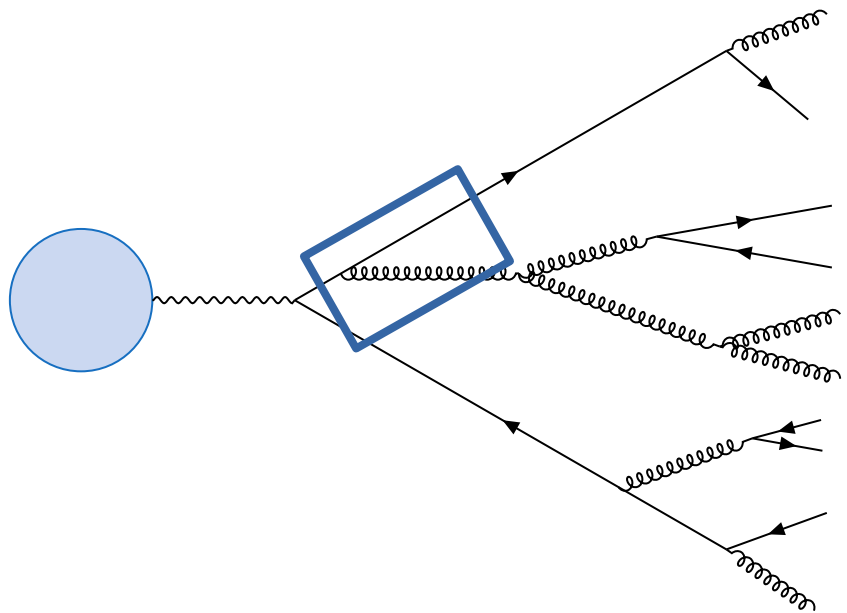


Splitting probability given by pQCD:



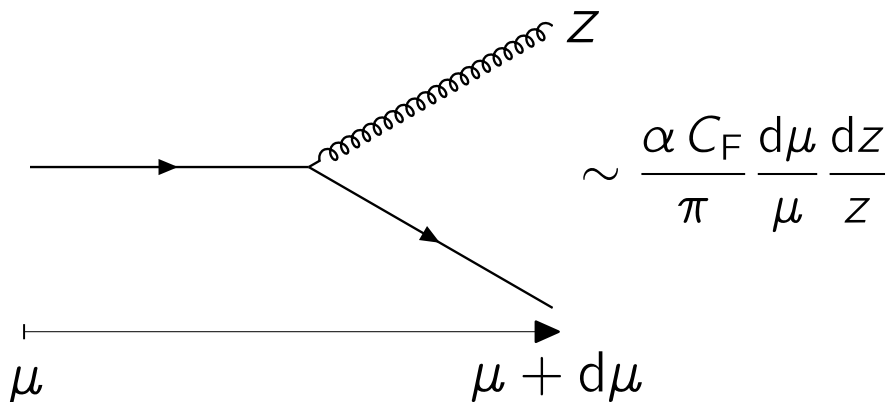
# How to build a parton shower

Building blocks: QCD splittings



→ Splittings with decreasing scale  $\mu$

Splitting probability given by pQCD:



Probability of not emitting until some scale  $s$  :

$$\Delta(s_{\text{prev}}, s) = \exp \left\{ -\frac{\alpha C_R}{\pi} \int_s^{s_{\text{prev}}} \frac{d\mu}{\mu} \int_{z_{\text{cut}}(\mu)}^1 \frac{dz}{z} \right\}$$

Yields the next emission scale  $s$  , given the previous scale  $s_{\text{prev}}$

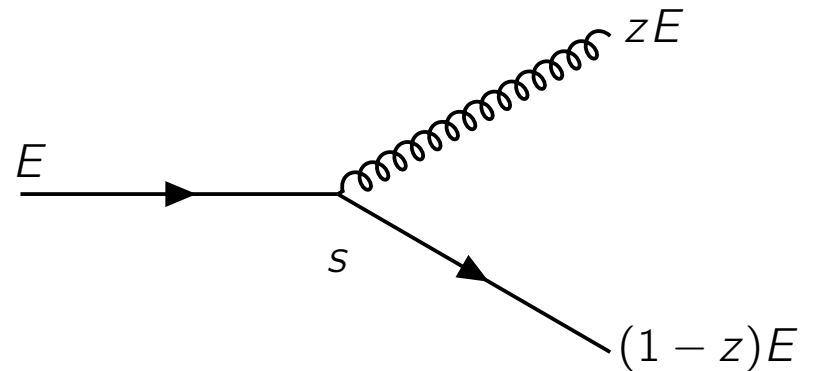


# Building differently ordered cascades

No-emission probability:

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**Splitting variables:**



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Interpretations for the scale:

$$s \rightarrow p^2 = \frac{|\mathbf{p}_{\text{rel}}|^2}{z(1-z)}$$

(Virtuality)

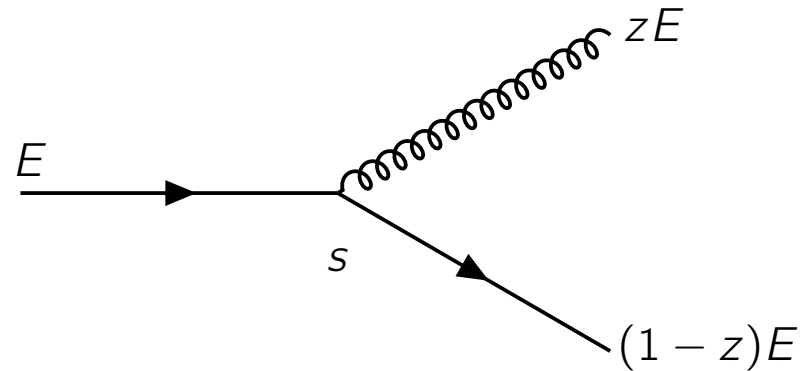
$$s \rightarrow t_{\text{form}}^{-1} = \frac{p^2}{E} = \frac{|\mathbf{p}_{\text{rel}}|^2}{Ez(1-z)}$$

(Formation time)

$$s \rightarrow \zeta = \frac{p^2}{E^2 z(1-z)} = \left( \frac{|\mathbf{p}_{\text{rel}}|}{E z(1-z)} \right)^2$$

(Angle)

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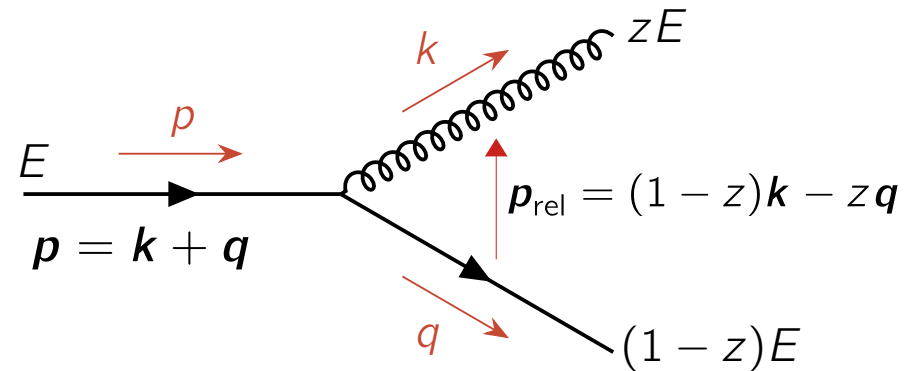
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To generate a splitting:



1. Sample a scale from  $\Delta(s_{\text{prev}}, s)$
  2. Sample a fraction from  $\hat{P}(z) \propto 1/z$
- Ensure that**  $|\mathbf{p}_{\text{rel}}|^2 > \Lambda^2$

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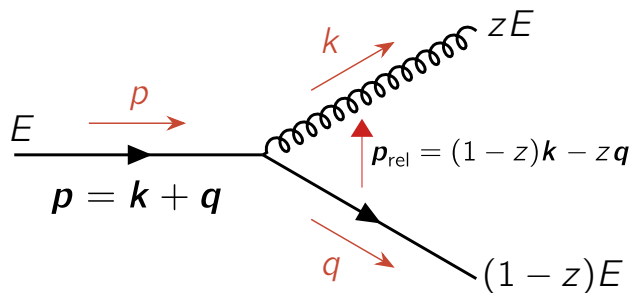
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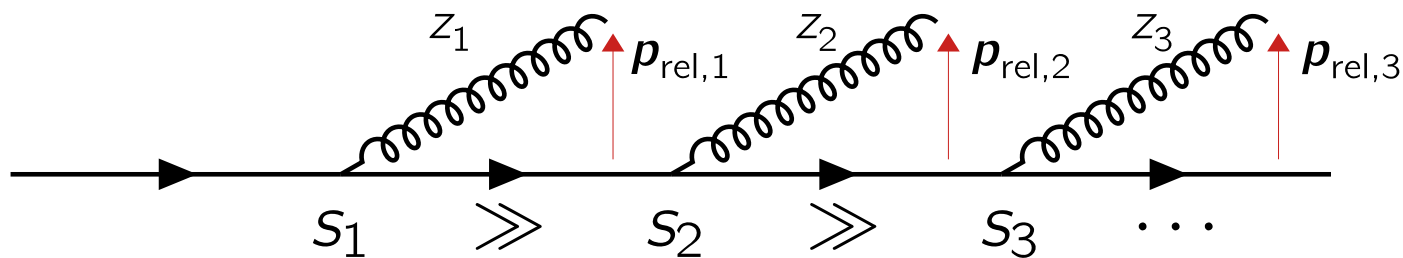
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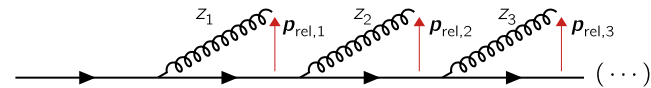


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This results in the strong ordering of scales

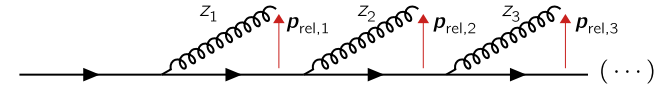
# Parton Shower Details



No-emission probability:

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# Parton Shower Details



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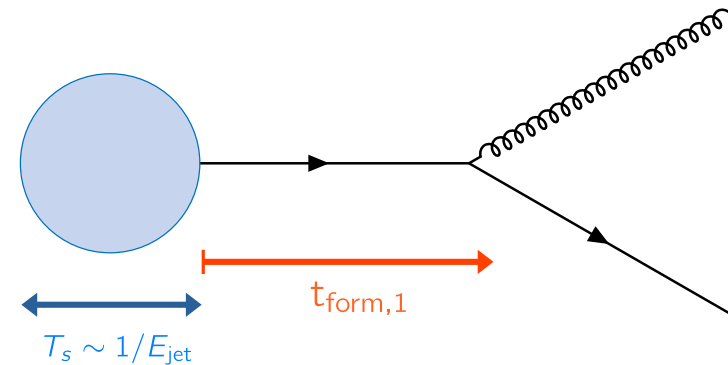
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- Splittings must happen above an hadronisation scale:  $|\mathbf{p}_{\text{rel}}|^2 > \Lambda^2$

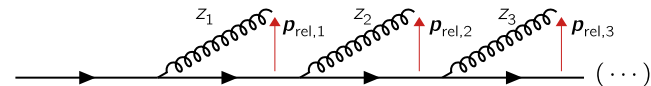
- This provides a **soft cutoff**:  $z > z_{\text{cut}}(s)$

e.g.: Formation time ordering  $|\mathbf{p}_{\text{rel}}|^2 > \Lambda^2 \iff z(1-z) > \frac{\Lambda^2}{t_{\text{form}}^{-1} E}$

- Initialisation condition for the shower:  $t_{\text{form}}^{-1} < E$



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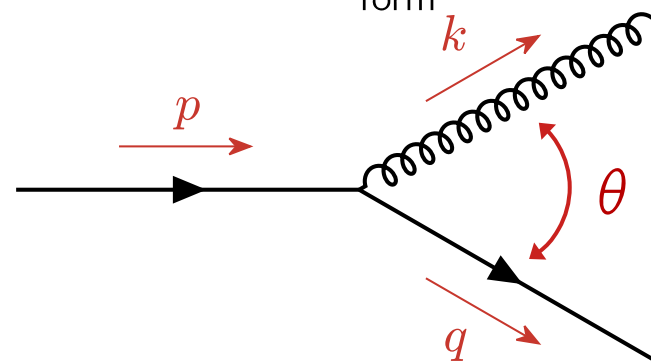
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- For consistency between orderings:  $\zeta < 4 \implies |\mathbf{p}_{\text{rel}}| < \frac{E}{2}$   
(Enforced via retries)



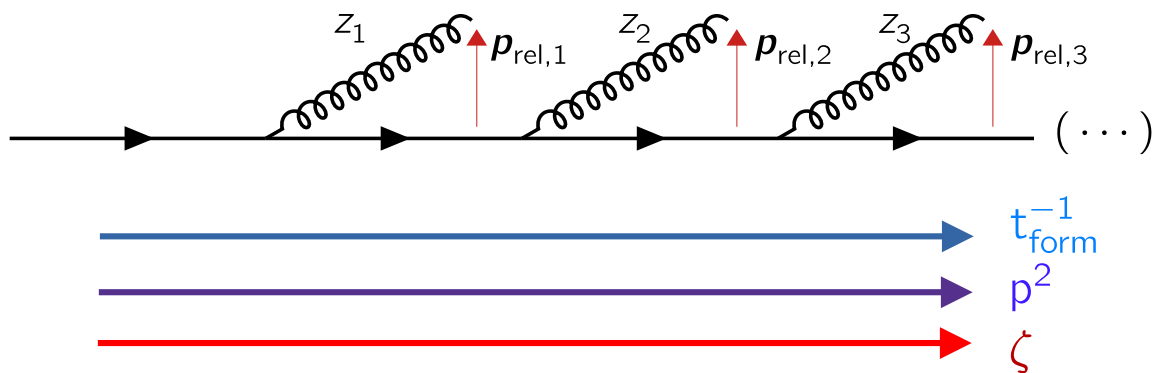
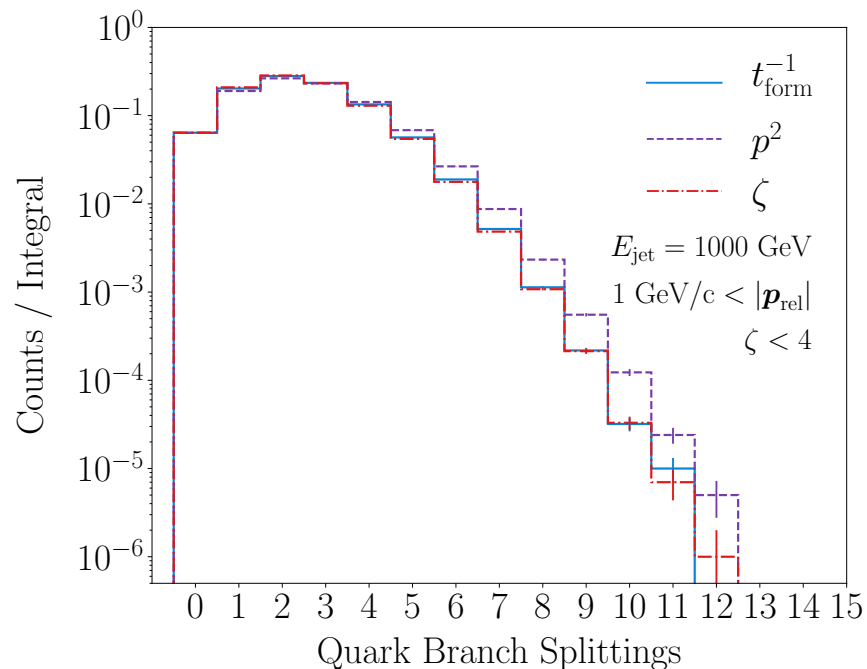
Massless Limit :  $\zeta \simeq 2(1 - \cos \theta)$  8

# Results (Work in Progress)



# Differences in Ordering Choices

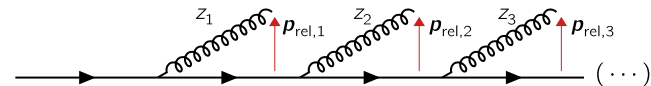
## Splittings along the quark branch



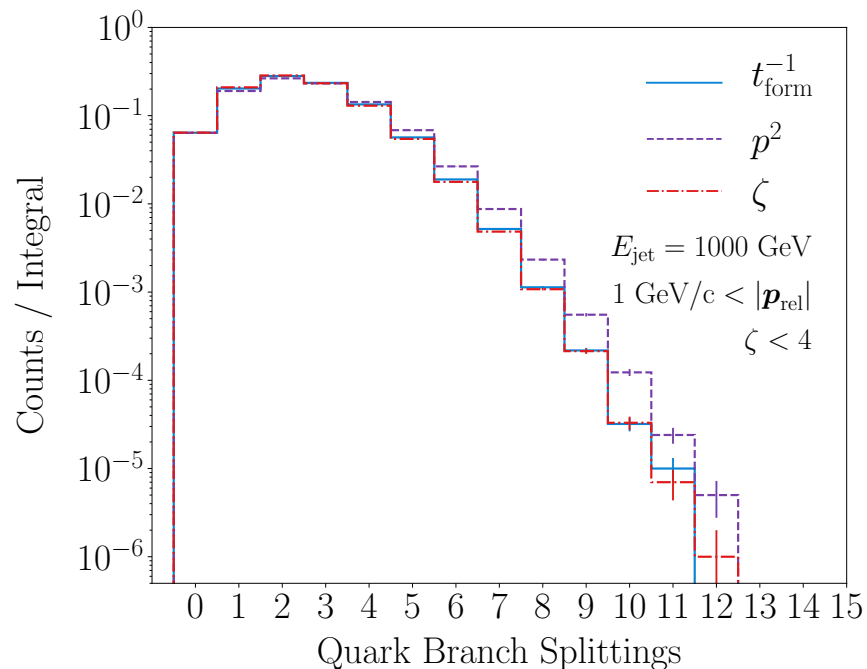
The strictly decreasing scale is different for the three algorithms

Different orderings  $\rightarrow$  Different phase-space for allowed splittings

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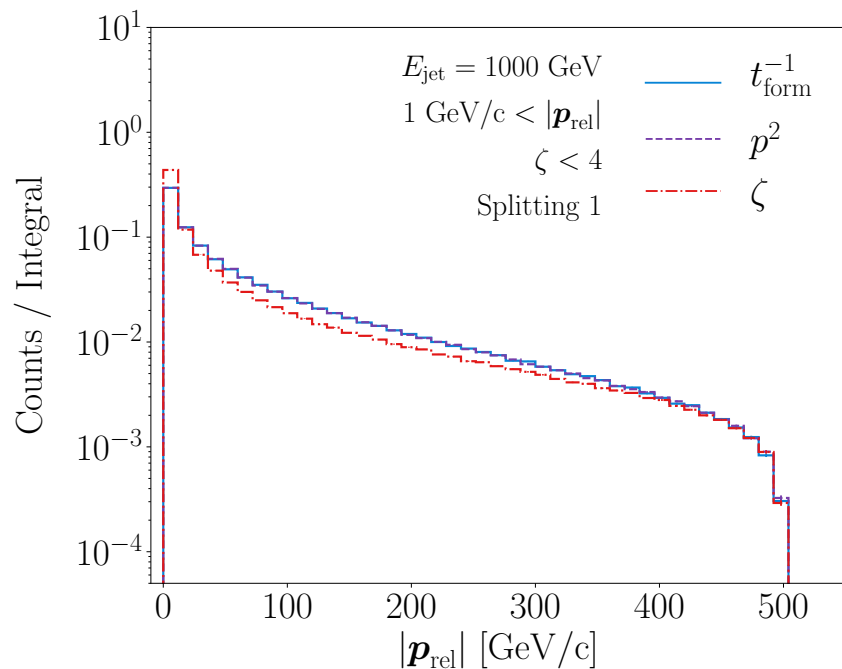


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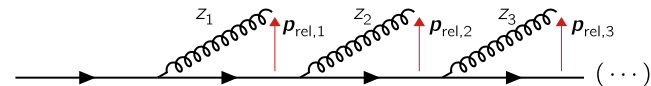
## Relative transverse momentum (1<sup>st</sup> splitting)



Transverse momentum distributions

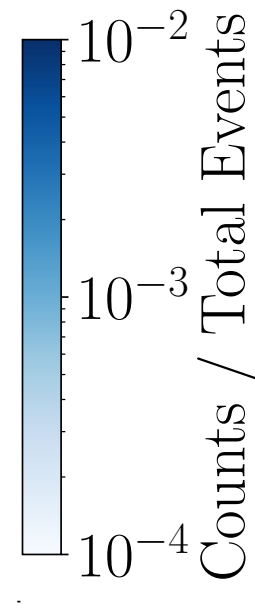
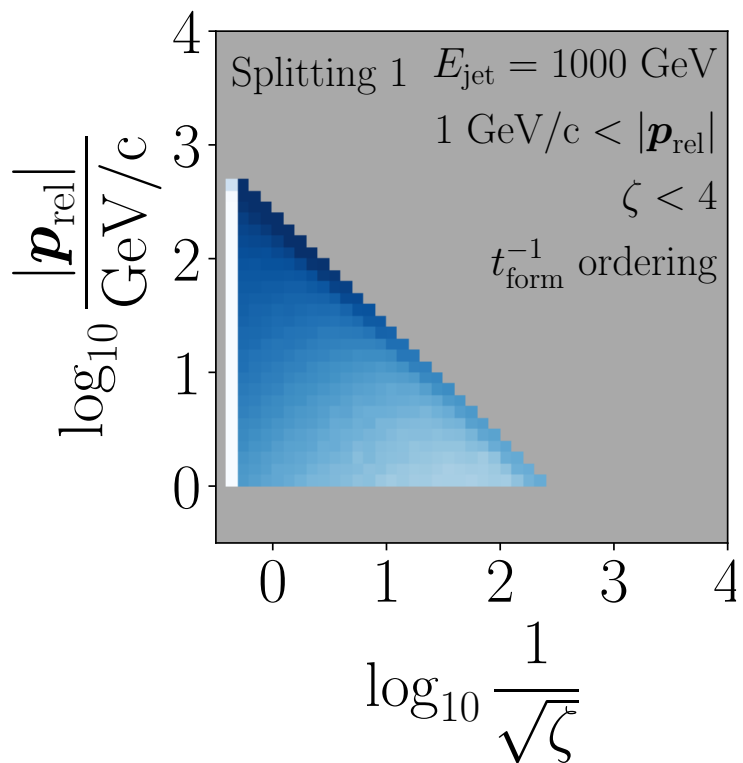
follow  $\frac{d p_{\text{rel}}^2}{p_{\text{rel}}^2}$

# Lund Plane Densities

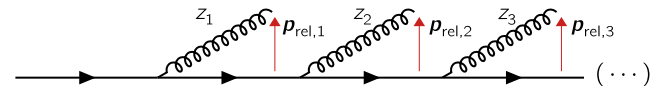


Consider the shower evolution along the quark branch:

\*Exaggerated scale

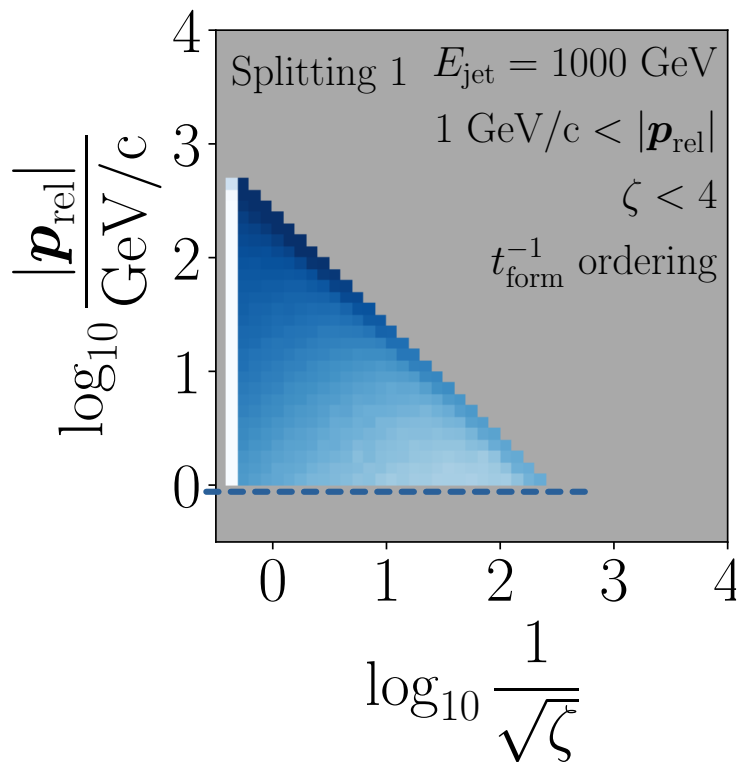


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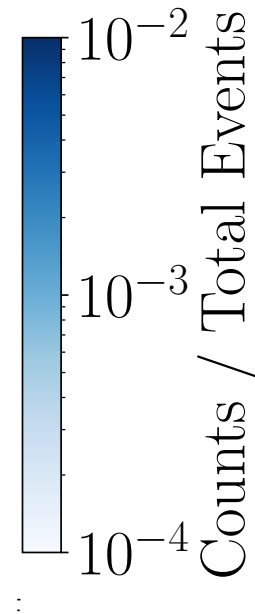
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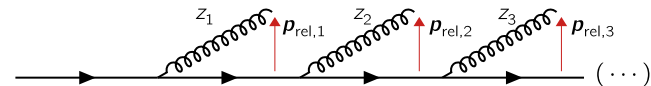


**Lund Plane Boundaries:**

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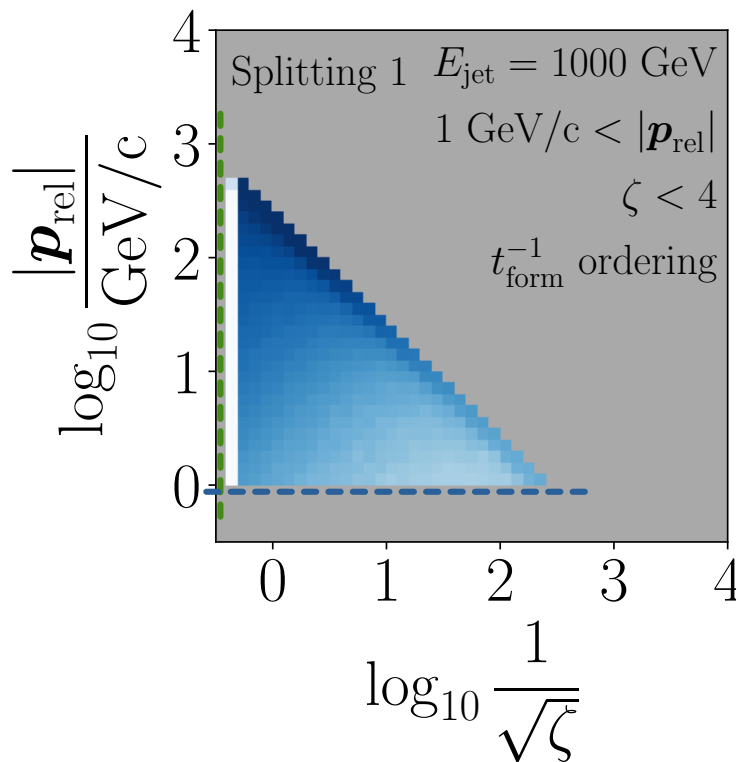


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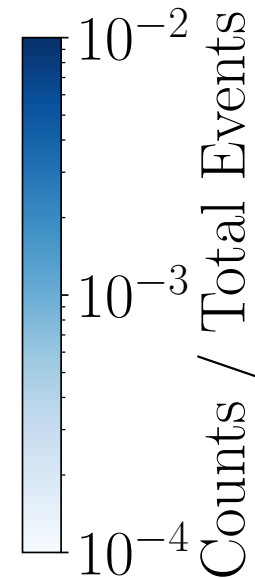
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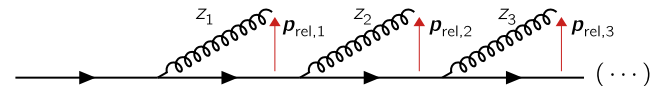


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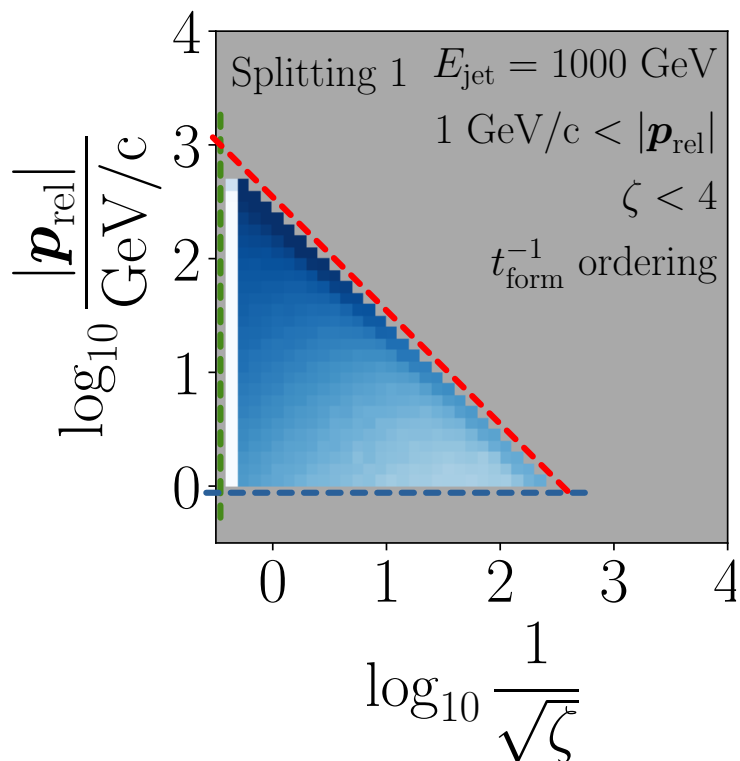


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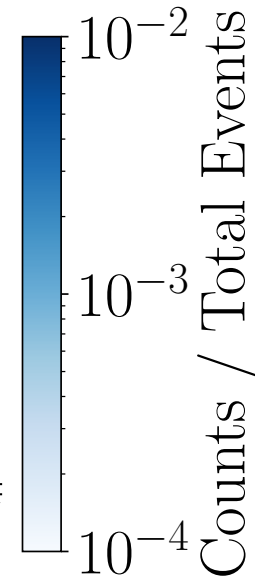


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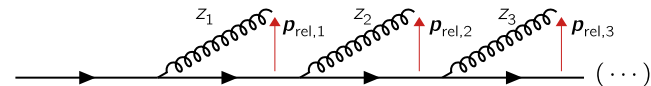
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- Energy constraint:  $\frac{E}{4} > E z(1-z) = |\mathbf{p}_{\text{rel}}| \frac{1}{\sqrt{\zeta}}$



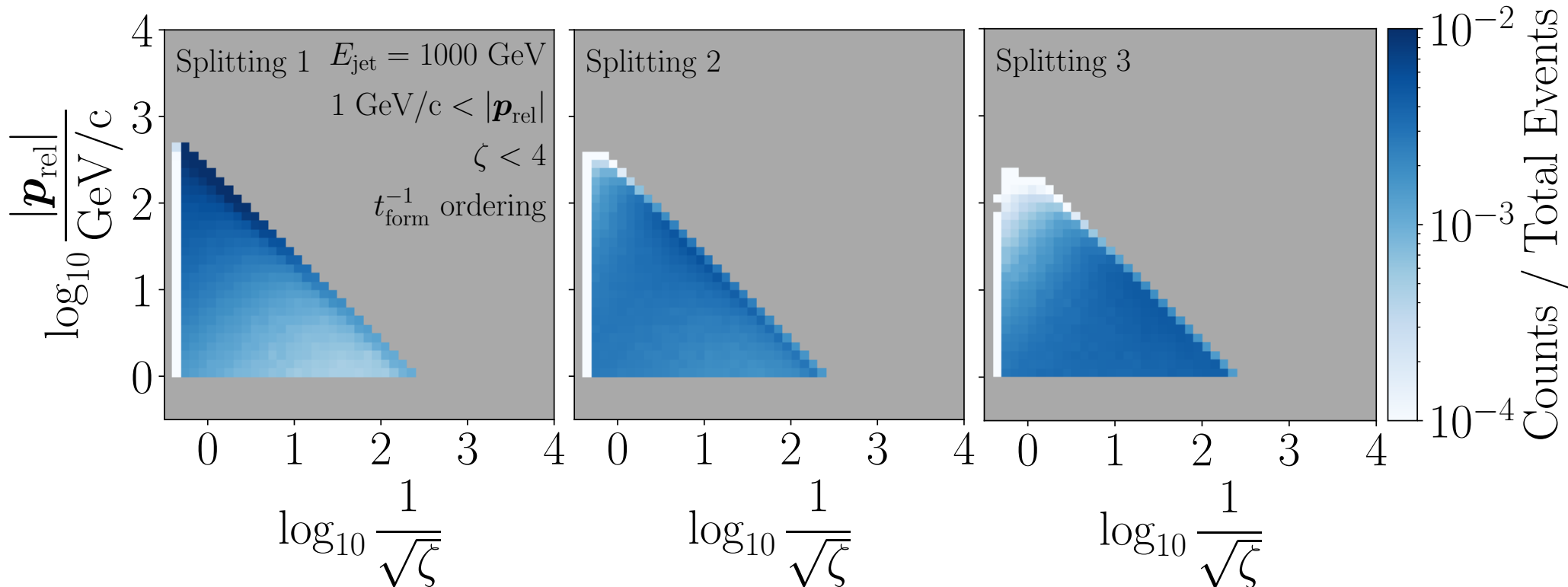
**Shower evolution:** Both transverse momentum and angle decrease.

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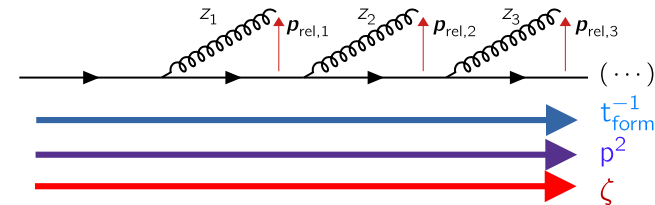
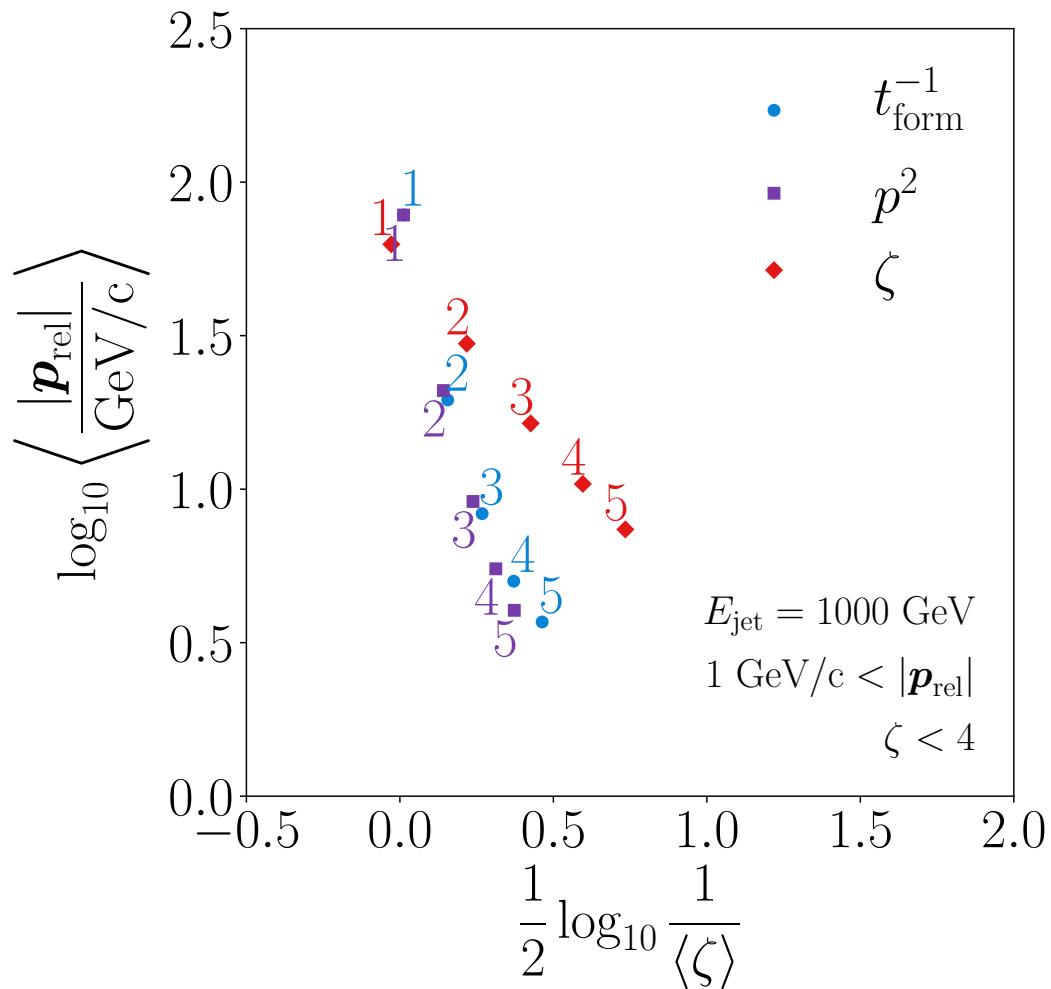
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# Lund Plane Trajectories

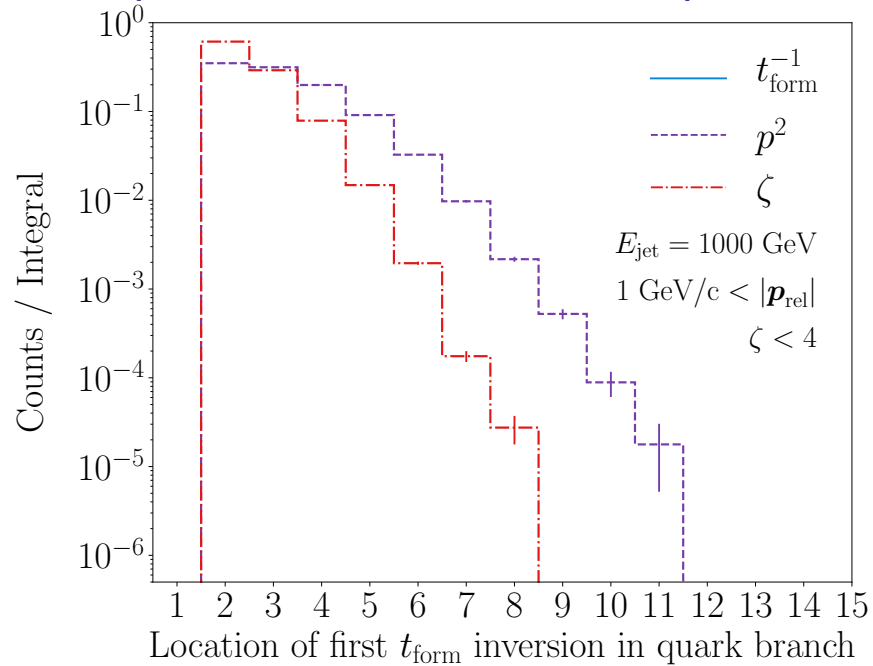


**Differences between  
 phase-space trajectories**  
**→ Uncertainty at DLA  
 Accuracy**



# Inversions in Kinematic Variables

(~ 30% Events with time inversions)

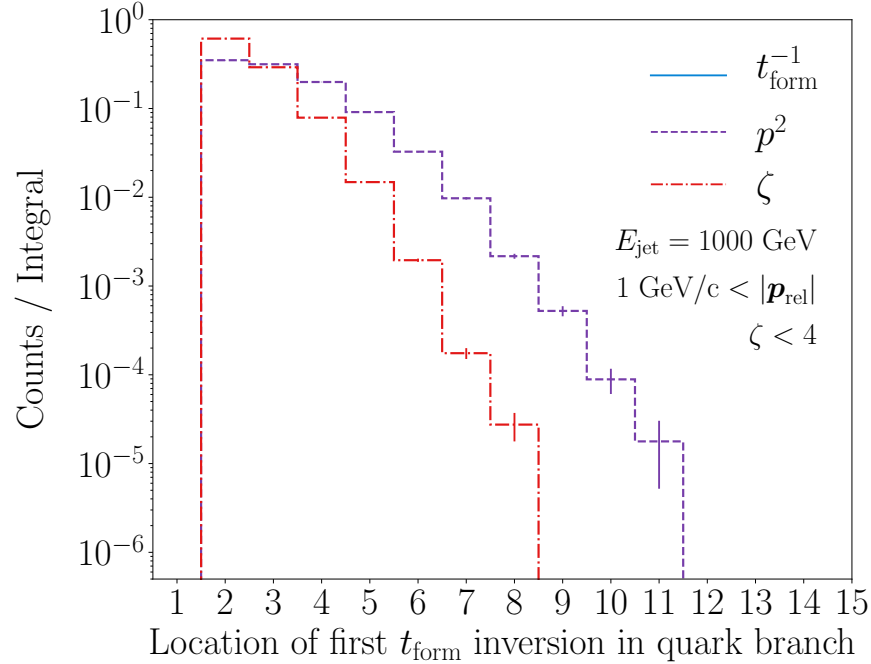


## Formation Time Inversions:

Splittings with a formation time shorter  
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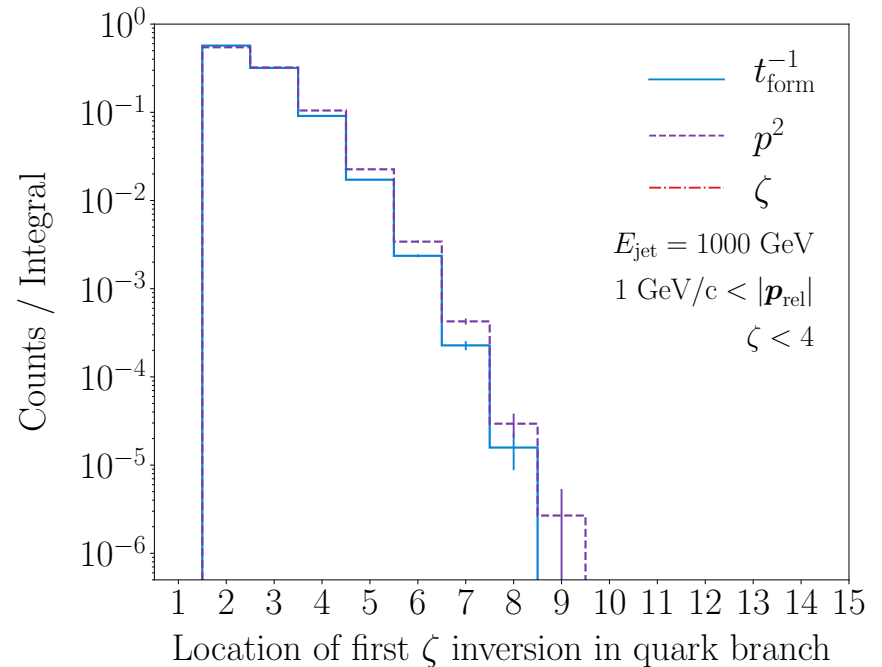
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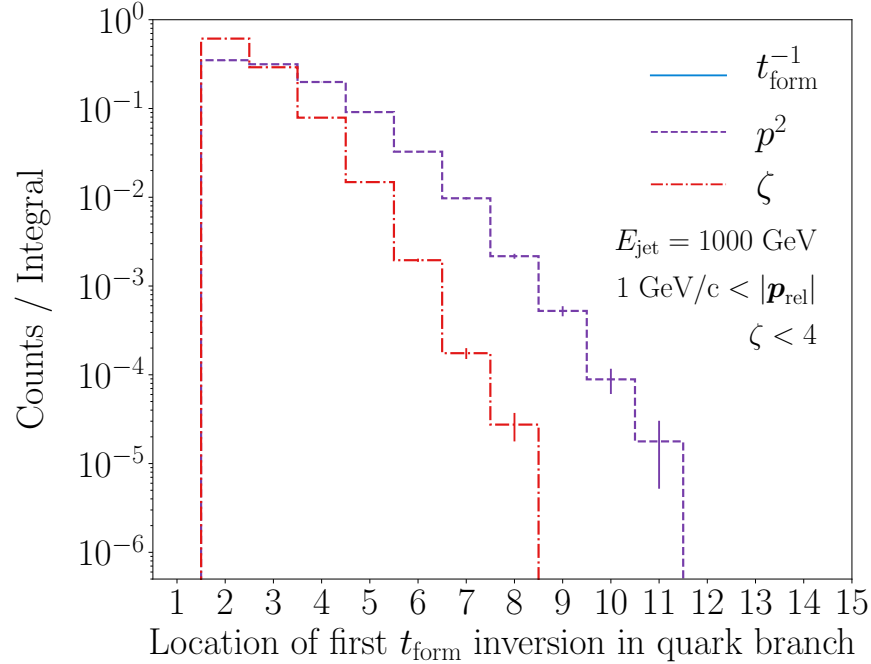
(~ 20% Events with  $\zeta$  inversions)



## Angular inversions

# Inversions in Kinematic Variables

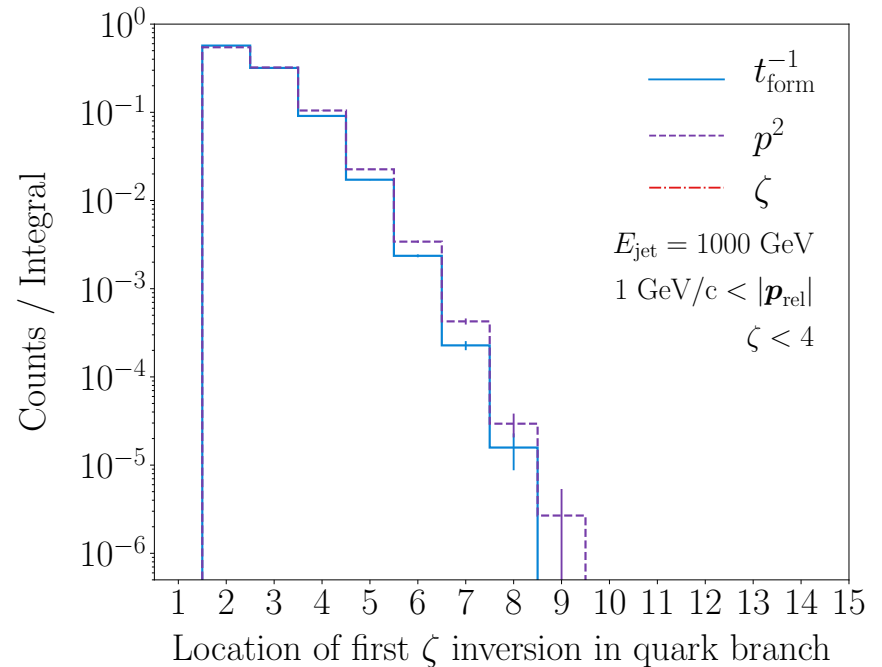
(~ 30% Events with time inversions)



## Formation Time Inversions:

Splittings with a formation time shorter than their immediate predecessor.

(~ 20% Events with  $\zeta$  inversions)

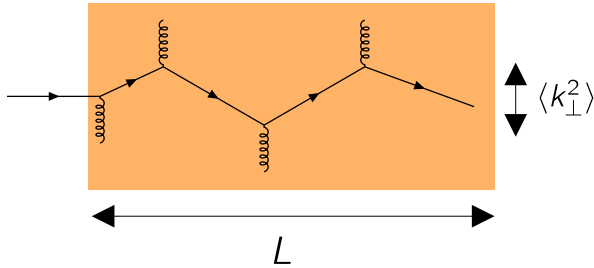


## Angular inversions

Can this discrepancy translate into differences in quenching magnitude?

**Now, a simple jet quenching model!**

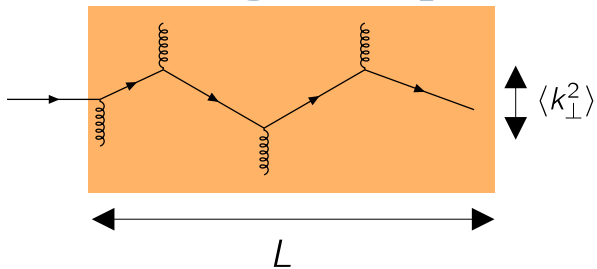
# Choosing a quenching condition



## Medium parameters (for a simple model):

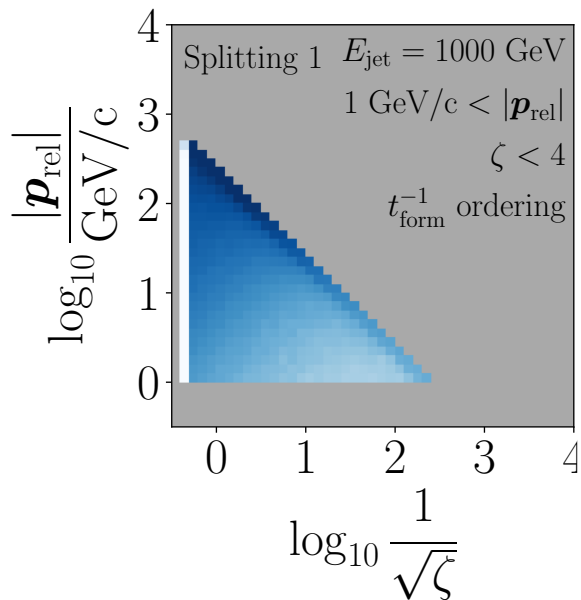
- Medium length:  $L$
- Transport coefficient:  $\hat{q} \sim \frac{\langle k_{\perp}^2 \rangle}{\lambda}$

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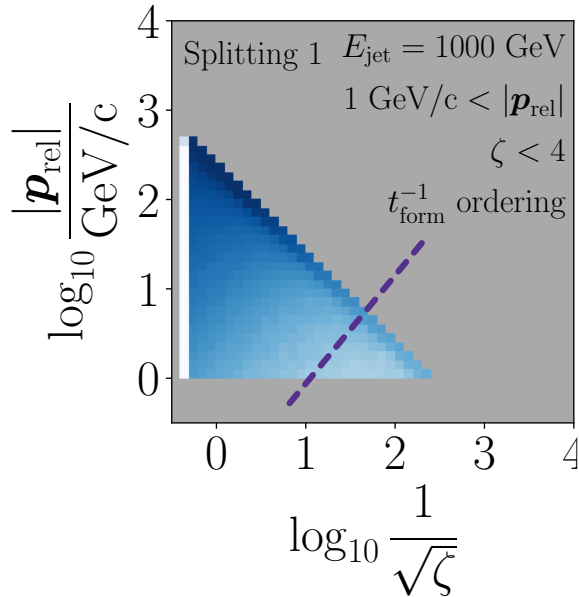
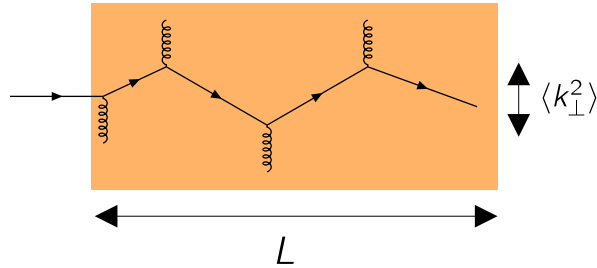


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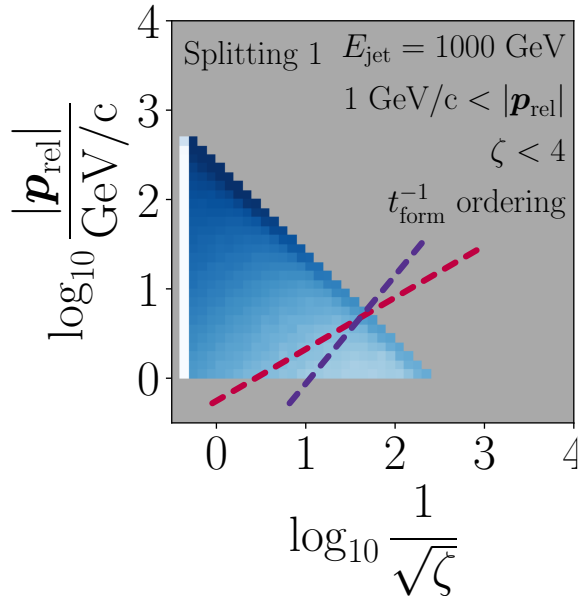
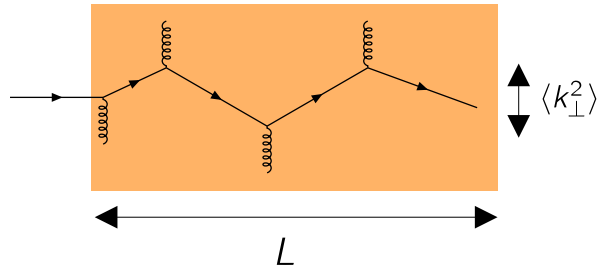
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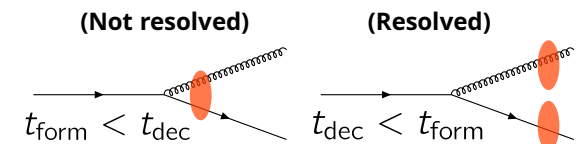
## Eliminate event if

- Splitting is inside the medium:  $t_{form} < L$
- Splitting transverse momentum is below medium scale:

$$|\mathbf{p}_{rel}|^2 < \hat{q} t_{form} \iff \underbrace{(\hat{q}\zeta)^{-1/3}}_{t_{dec}} < t_{form}$$

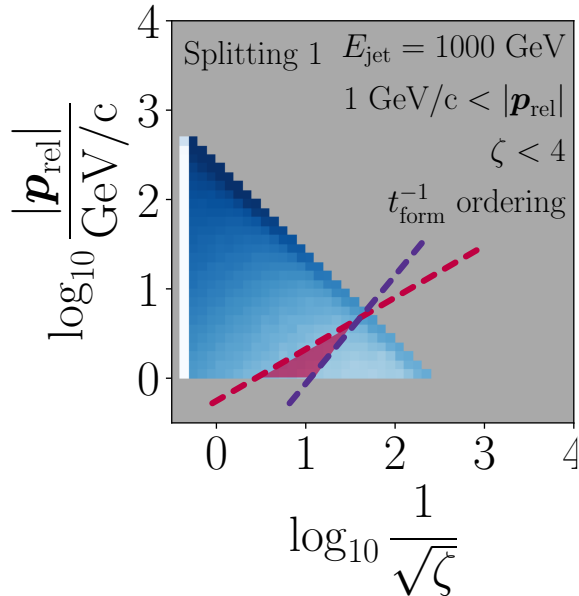
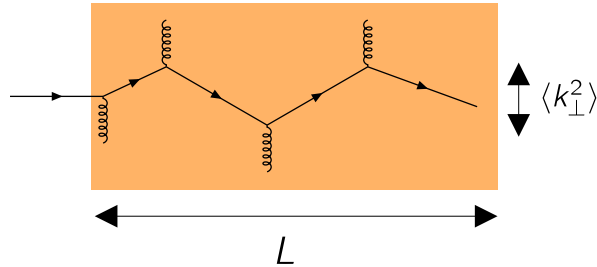
**Medium resolves splittings on the (de)coherence time scale**

→ Daughters lose energy individually (cf. antenna)





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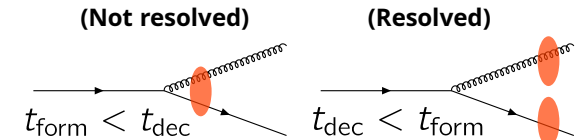
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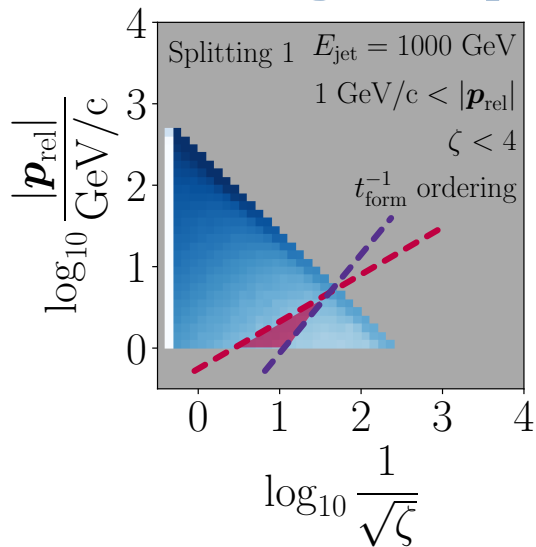
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**Eliminate events within this area:**

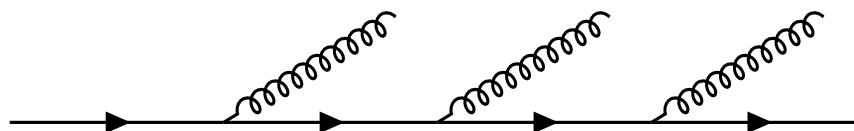
$$\mathcal{P}_{quench} = \Theta(L > t_{form} > t_{dec})$$

# Choosing a quenching condition

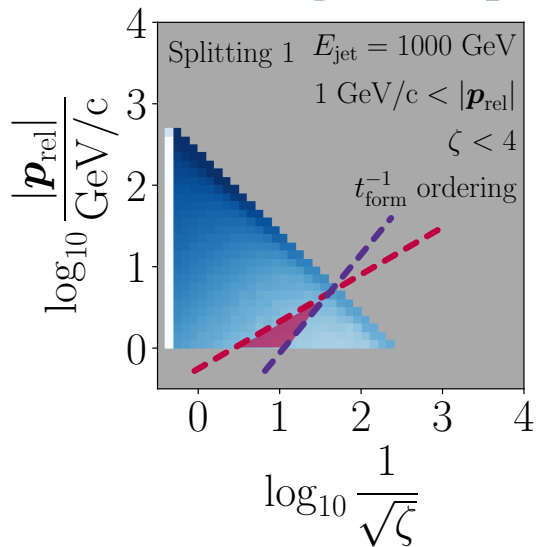


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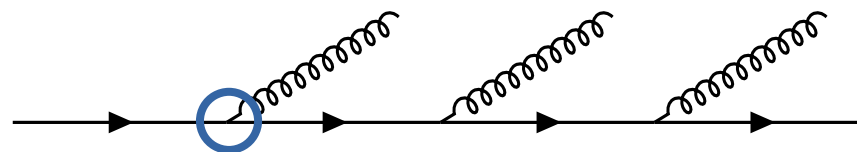
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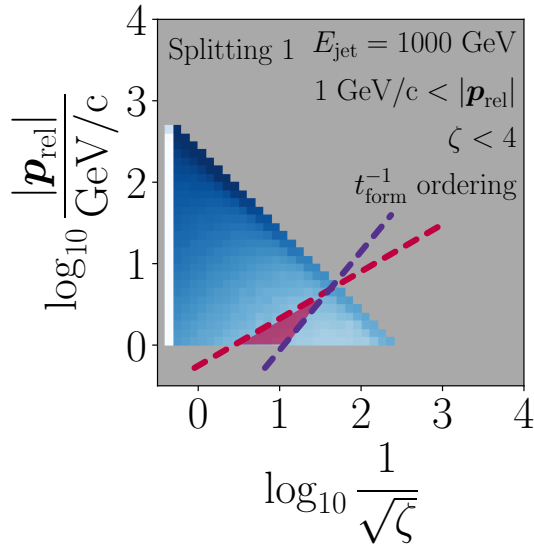
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**Two implementations:**



- Option 1: Apply only to first splitting

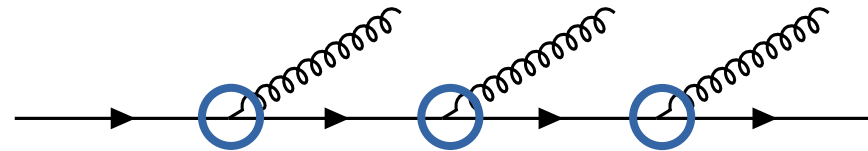
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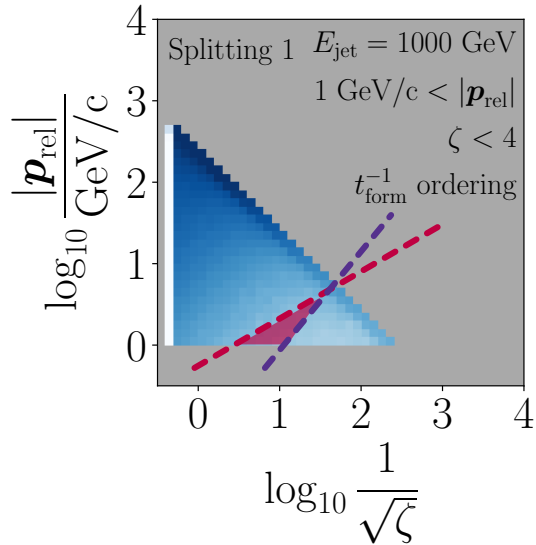
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- Option 1: Apply only to first splitting
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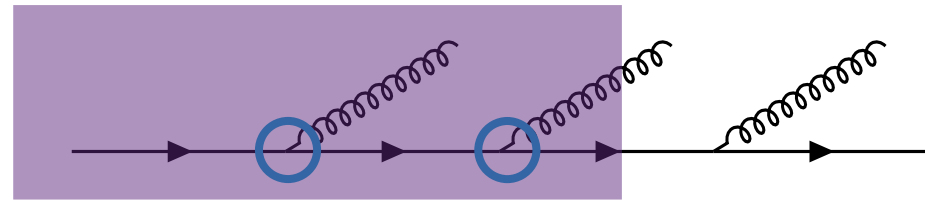
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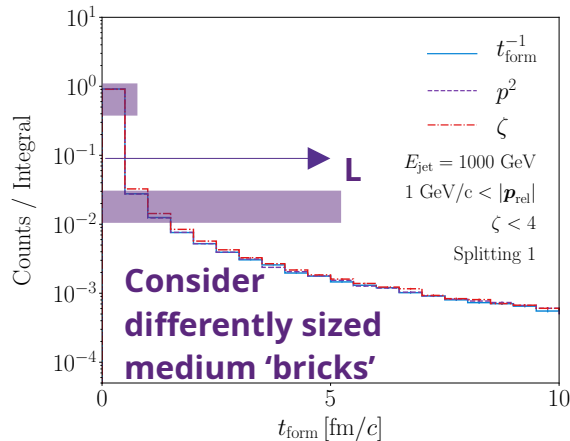
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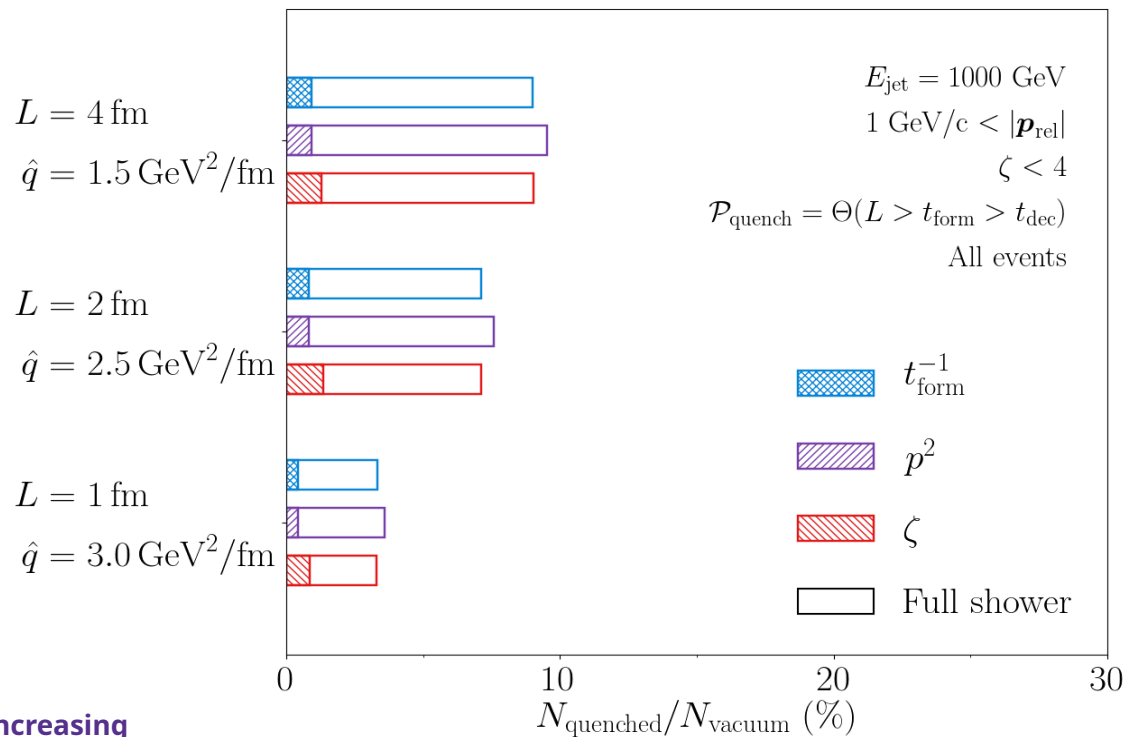
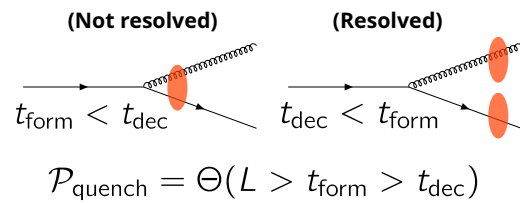


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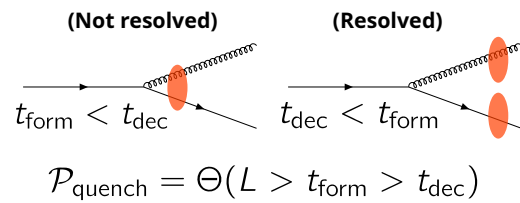
# Fraction of Quenched Events

## Percentage of events eliminated by the quenching condition

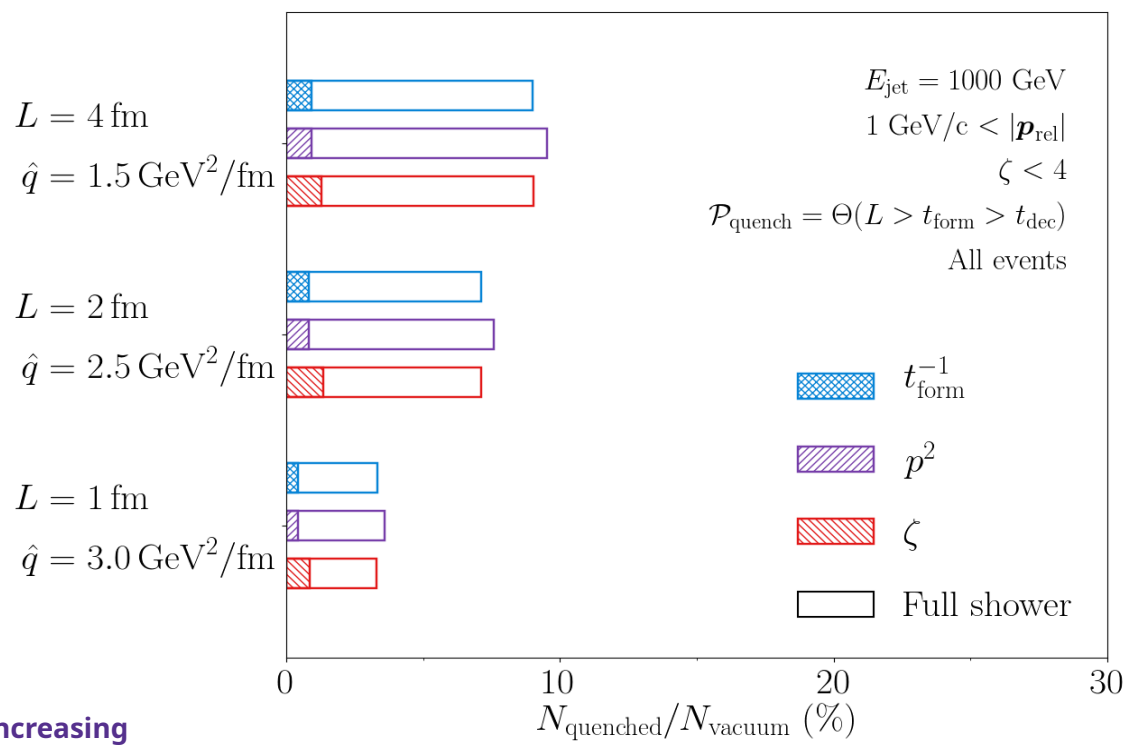


Increasing  
quenching effects

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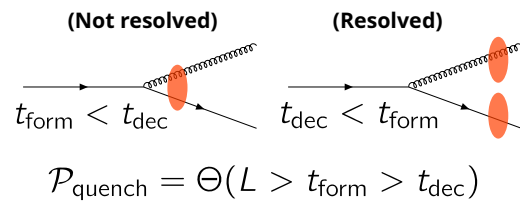


**Applying condition to the first splitting** → Significant differences in quenching between algorithms

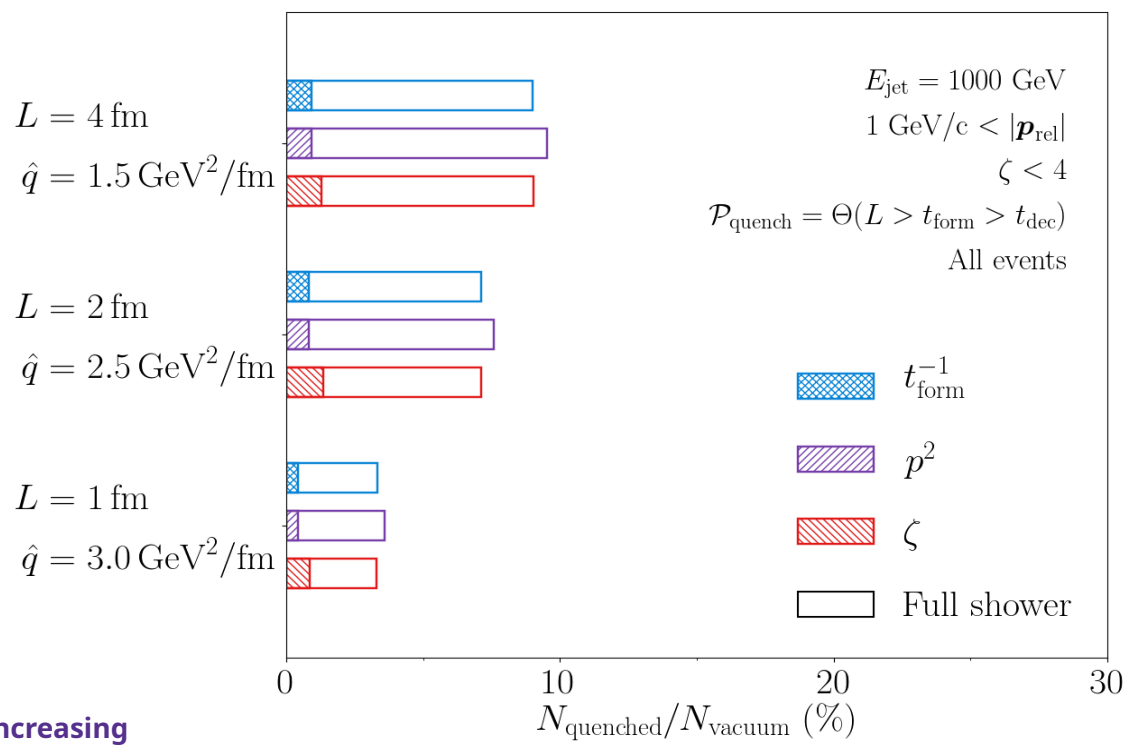
Differences are **seen to remain (for larger L)** when applying the condition to the full quark branch.

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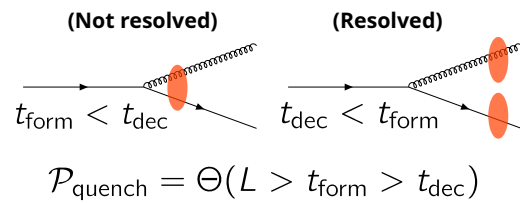
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**What role do time-inversions play in these quenching differences?**

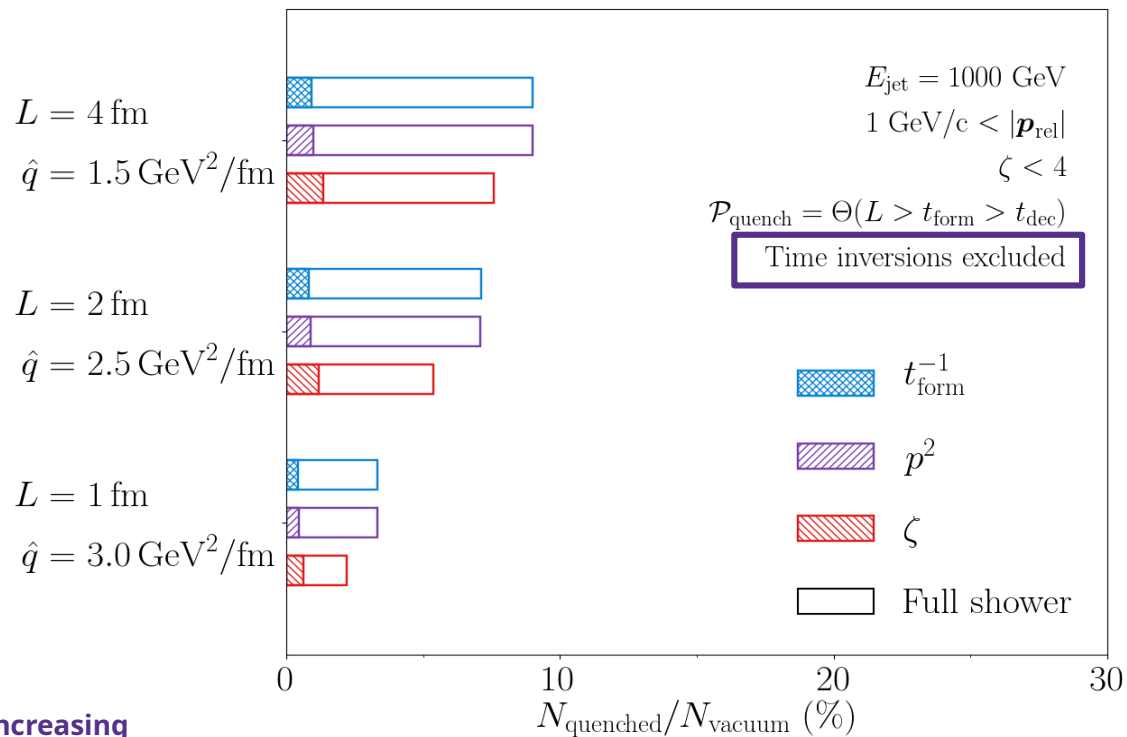


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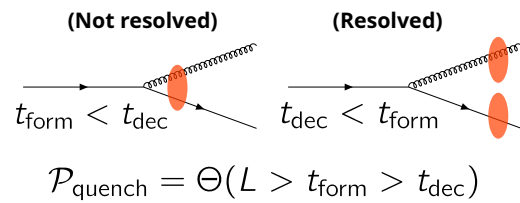
Discarding time-inverted events from the samples:

(Ad-hoc 'cut')



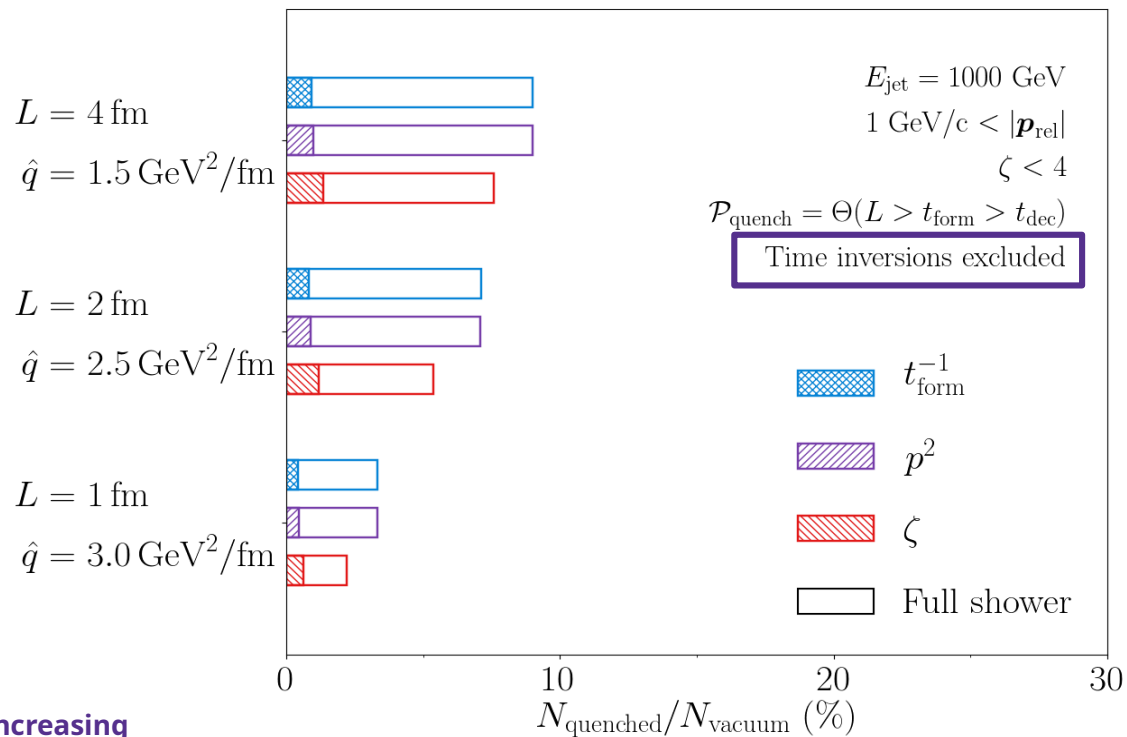
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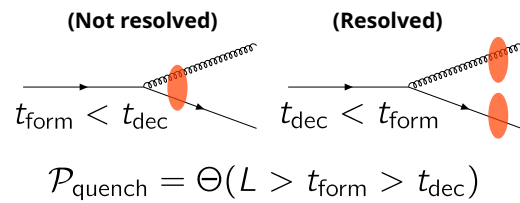
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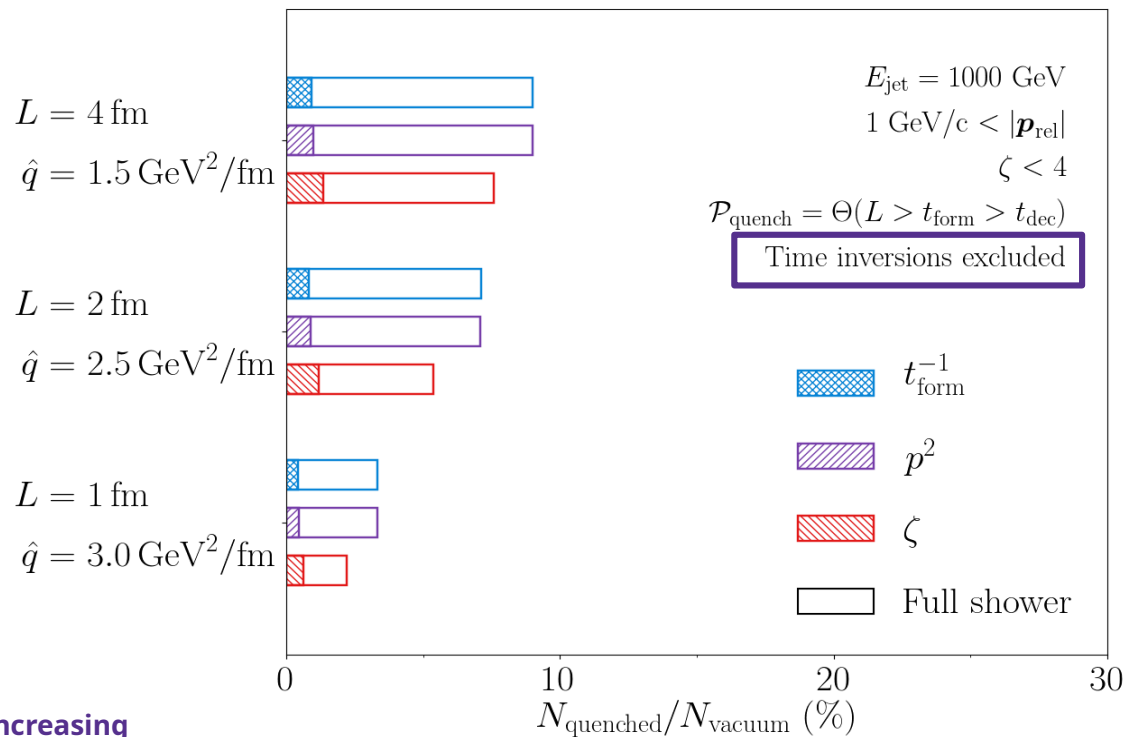
For angular ordered showers:  
 $\Rightarrow \zeta$  strictly decreasing  
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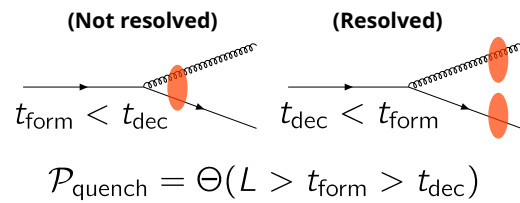


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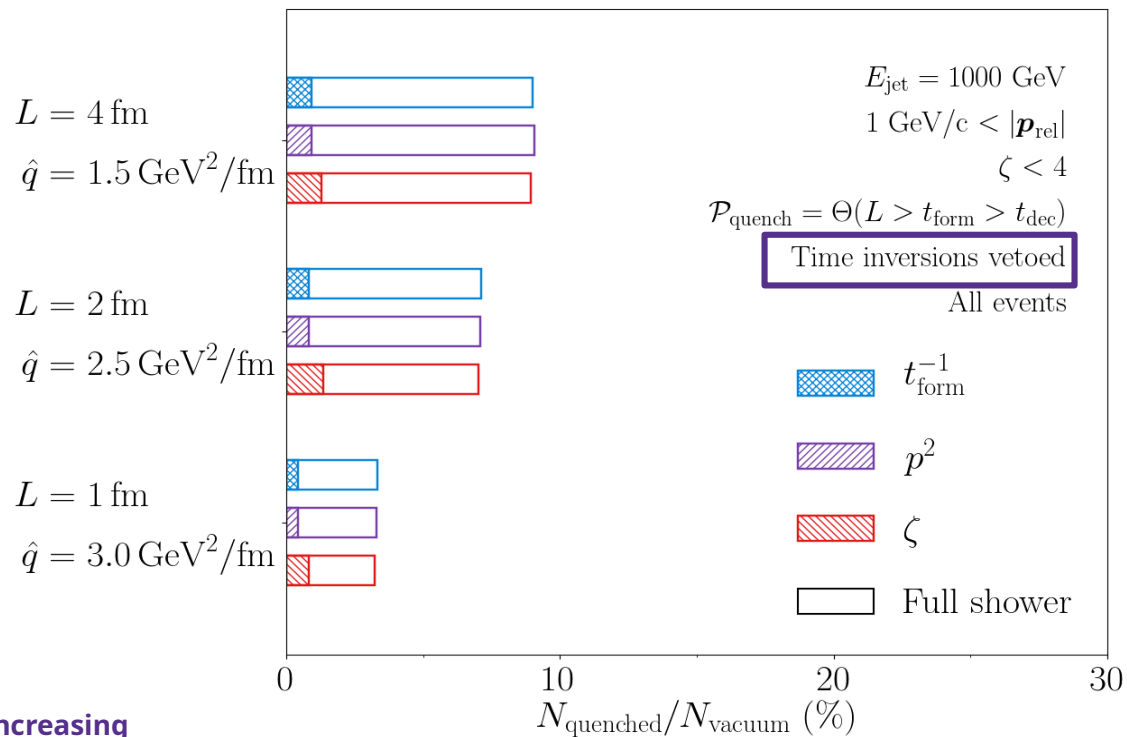
This is only one way of preventing inversions!

# Fraction of Quenched Events



Vetoing the time-inversions by retrial:

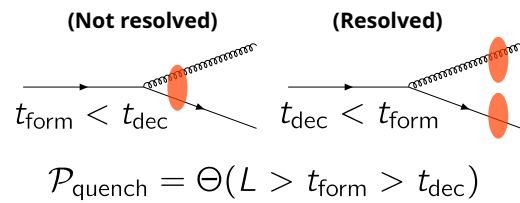
(Phase-space is adjusted splitting by splitting)



\*\*\* Time-inverted splittings are re-tried while generating the shower

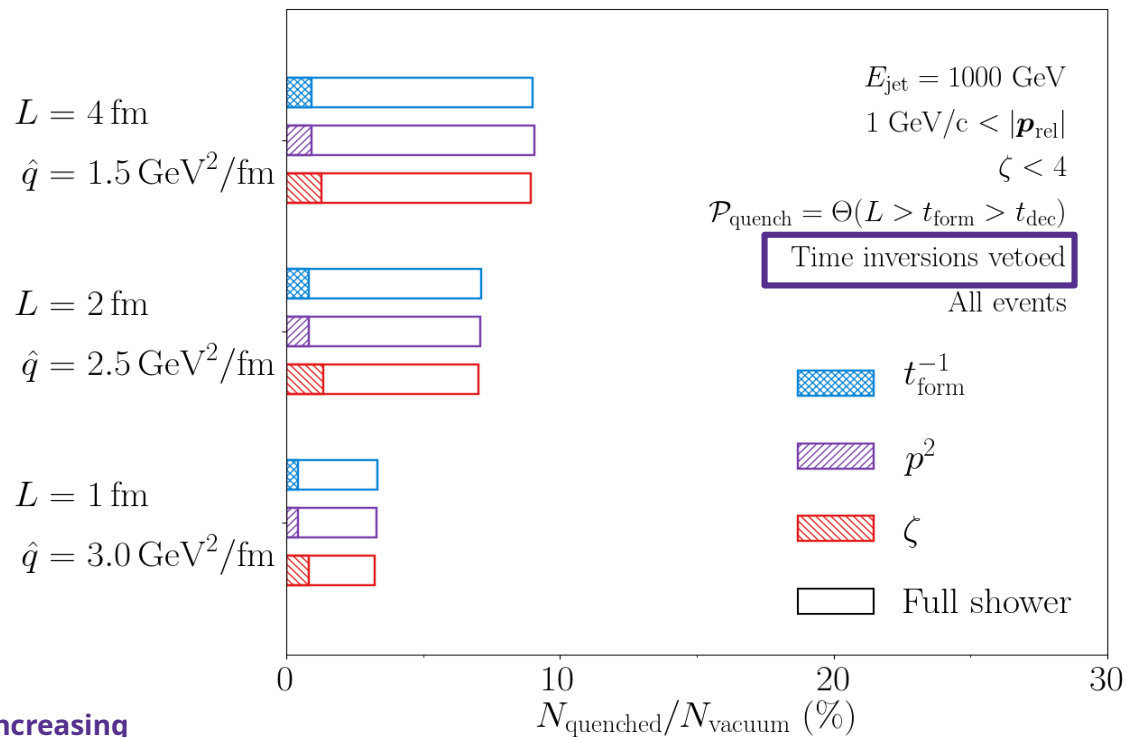
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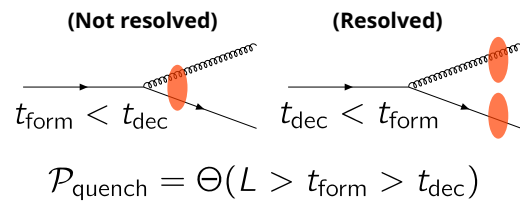
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Fraction of quenched events remains levelled across algorithms for the 'Full Branch' condition

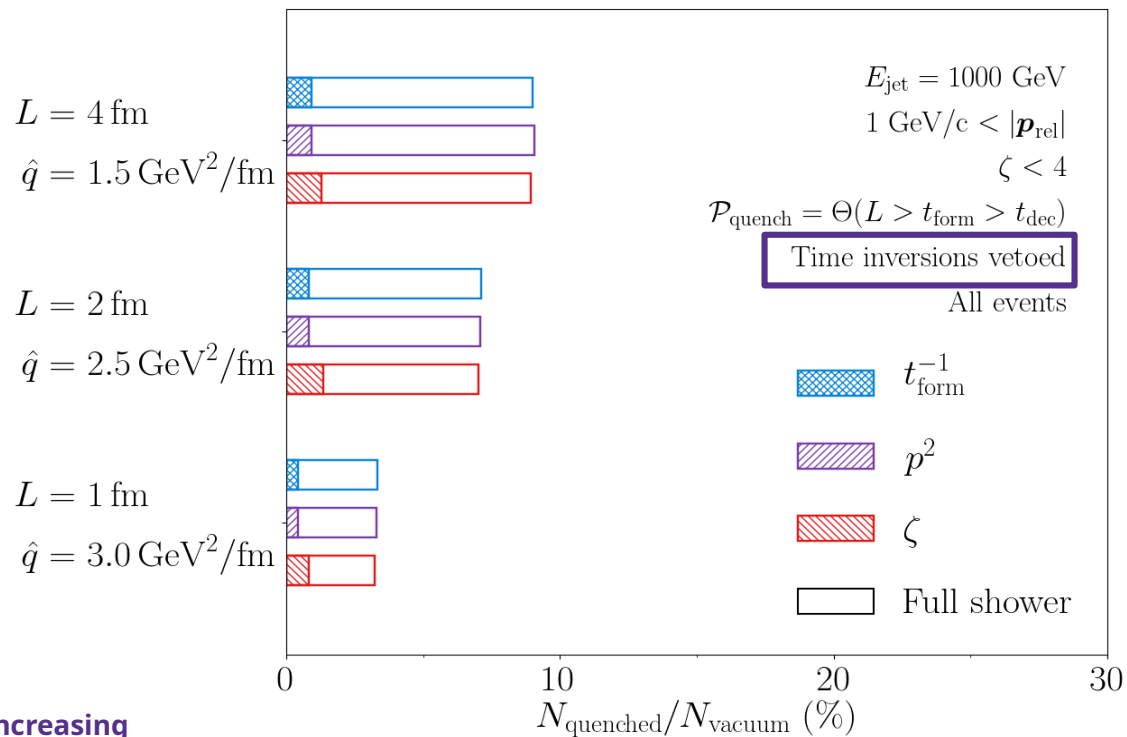
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The implementation details of the jet interface with a time-evolving medium are crucial!

# Summary

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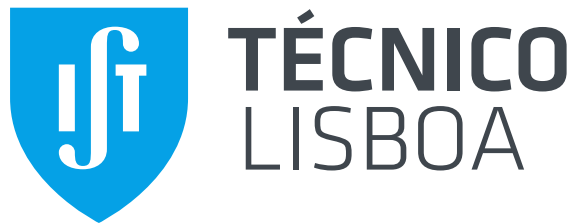
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# Acknowledgements



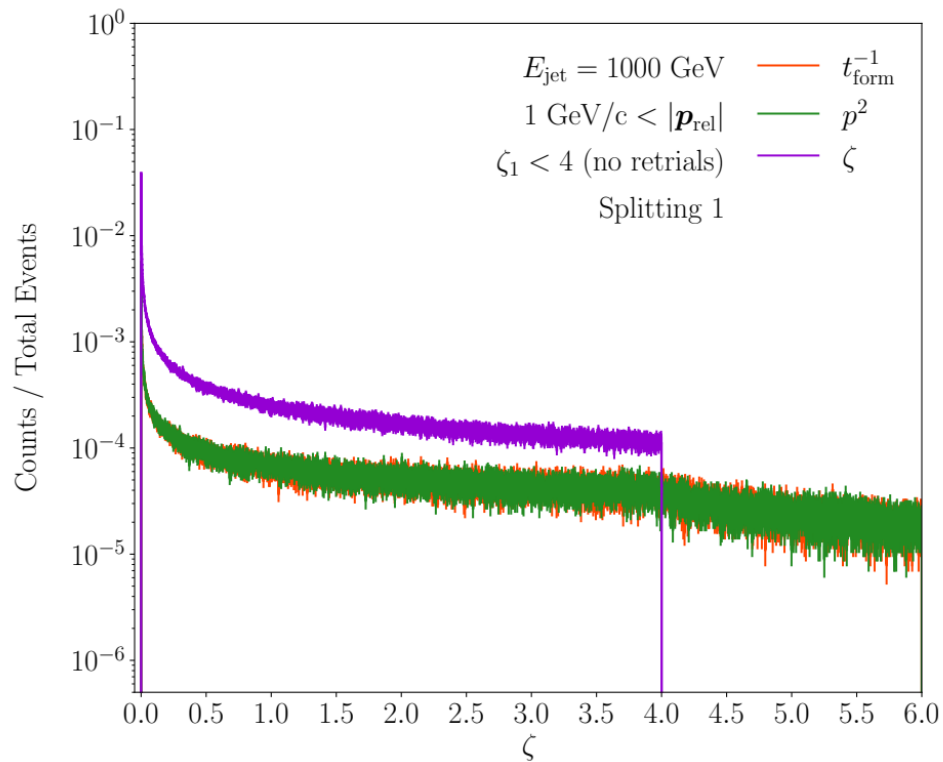
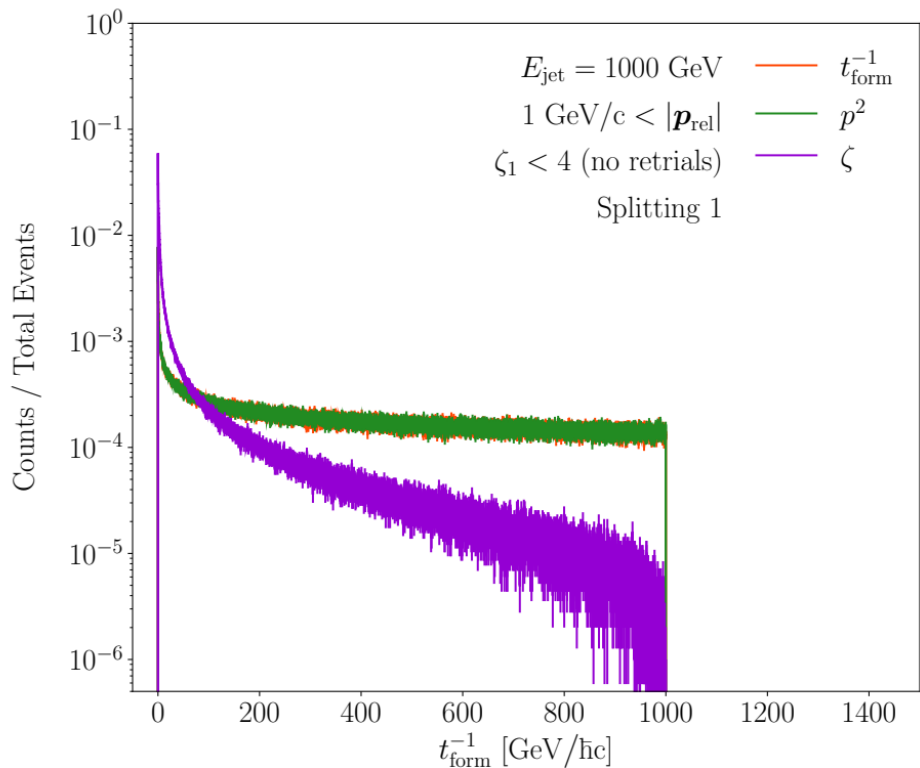
Fundação  
para a Ciência  
e a Tecnologia



REPÚBLICA  
PORTUGUESA

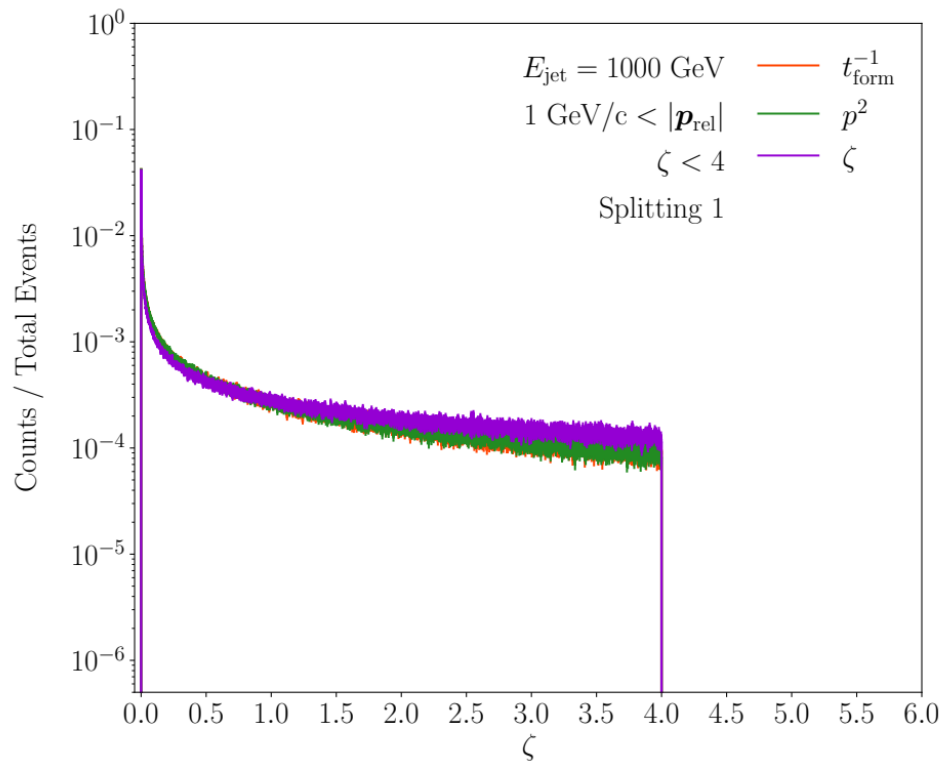
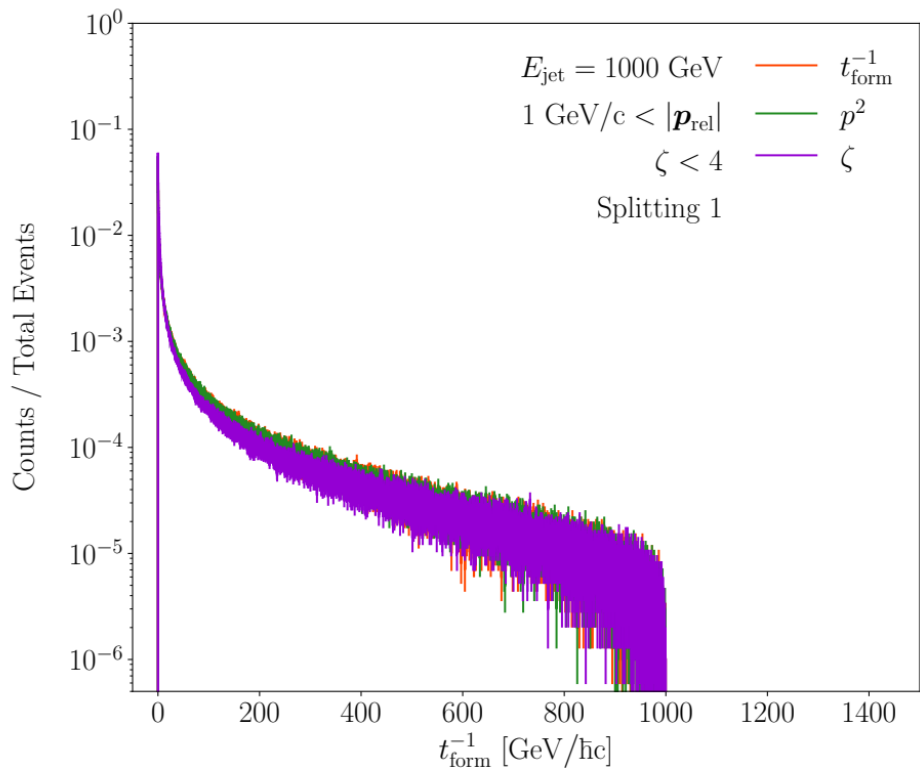
# Backup Slides

# Without the consistency condition



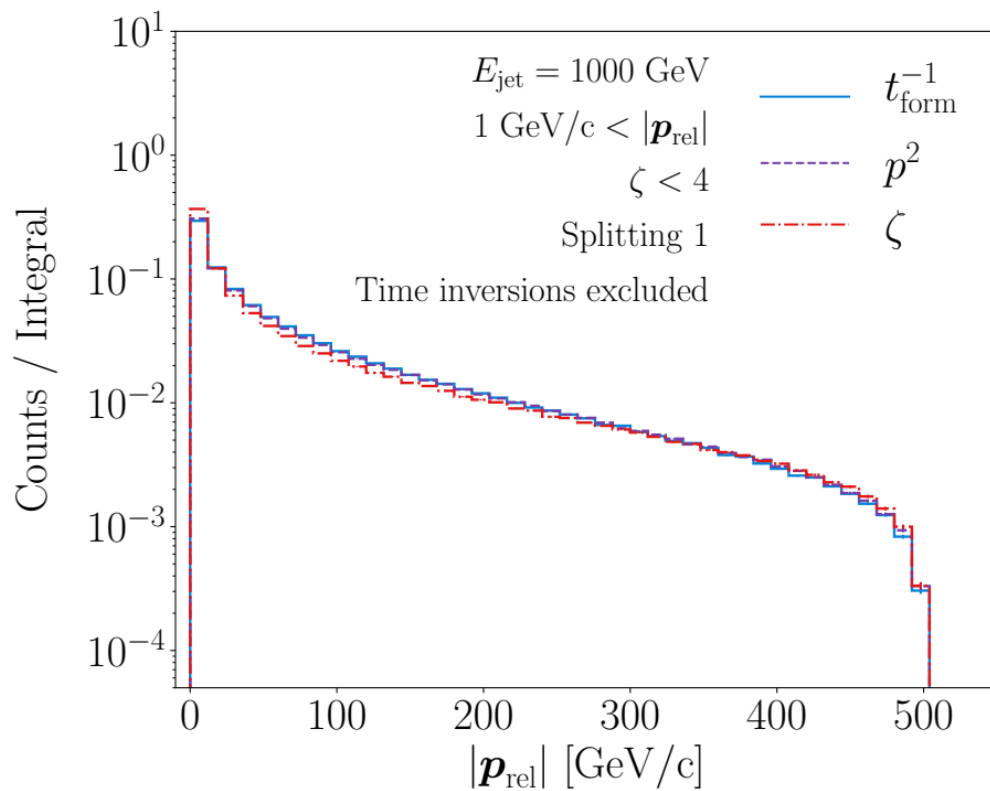
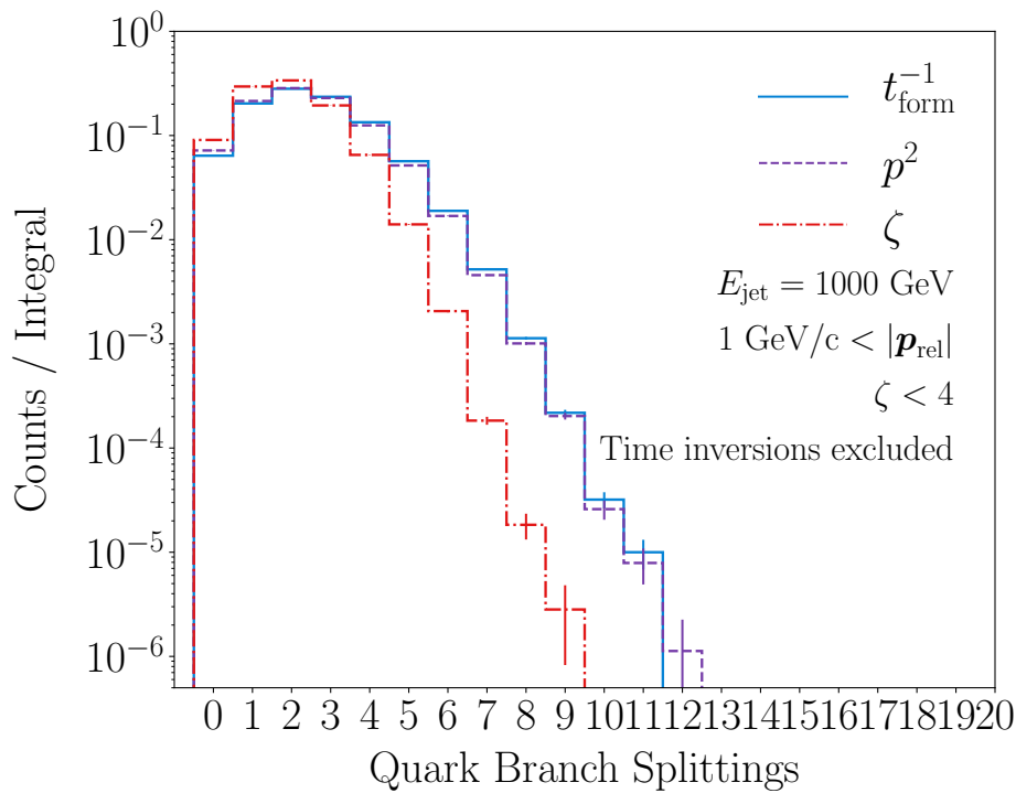
If the condition  $\zeta < 4$  is used simply to initialise the angular shower, the time and angle distributions do not behave consistently across algorithms

# With the consistency condition

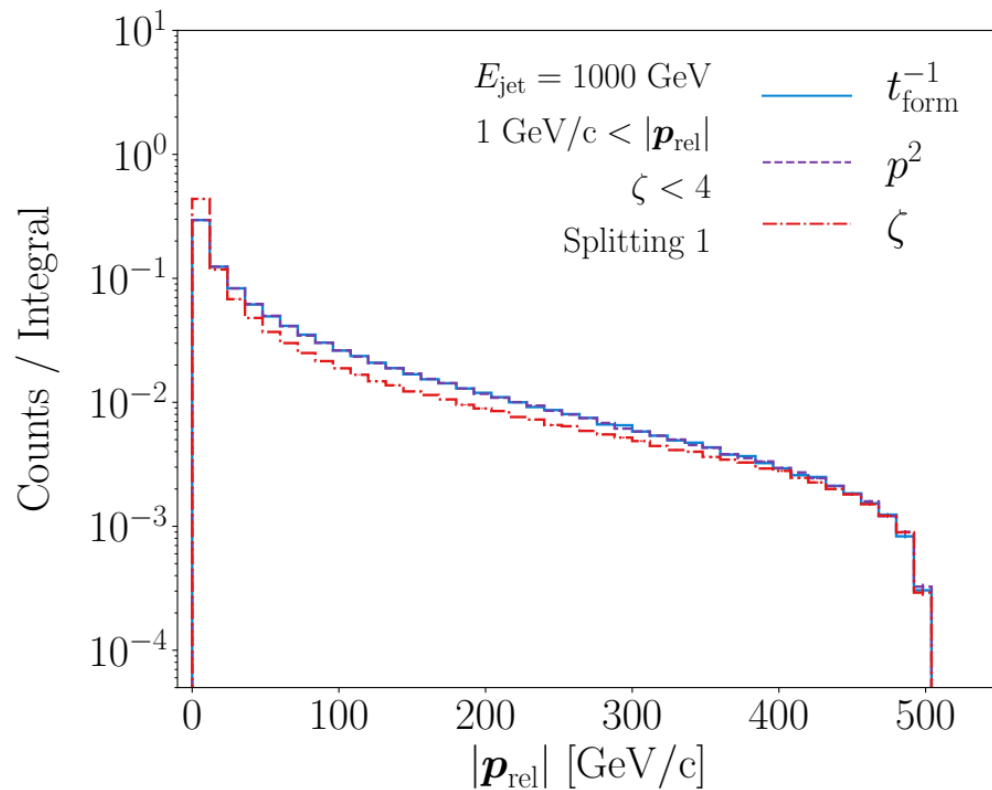
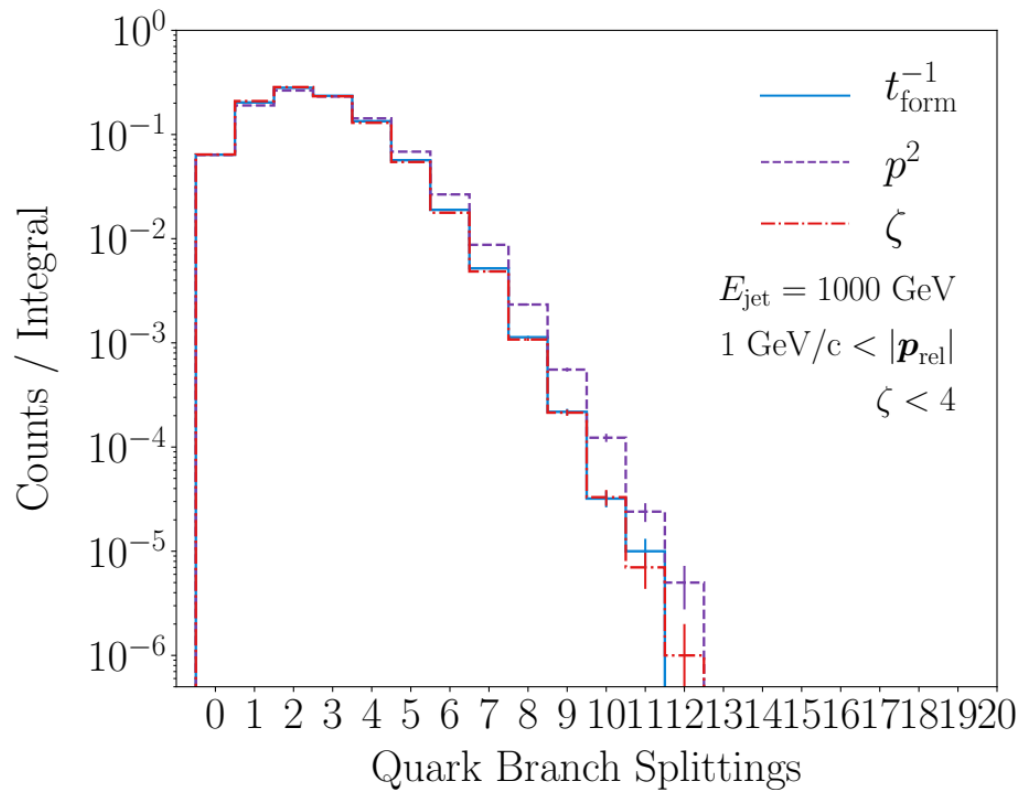


When the condition  $\zeta < 4$  is used as a veto for all emissions, the distributions become consistent.

# Excluding time inversions – 1D Distributions

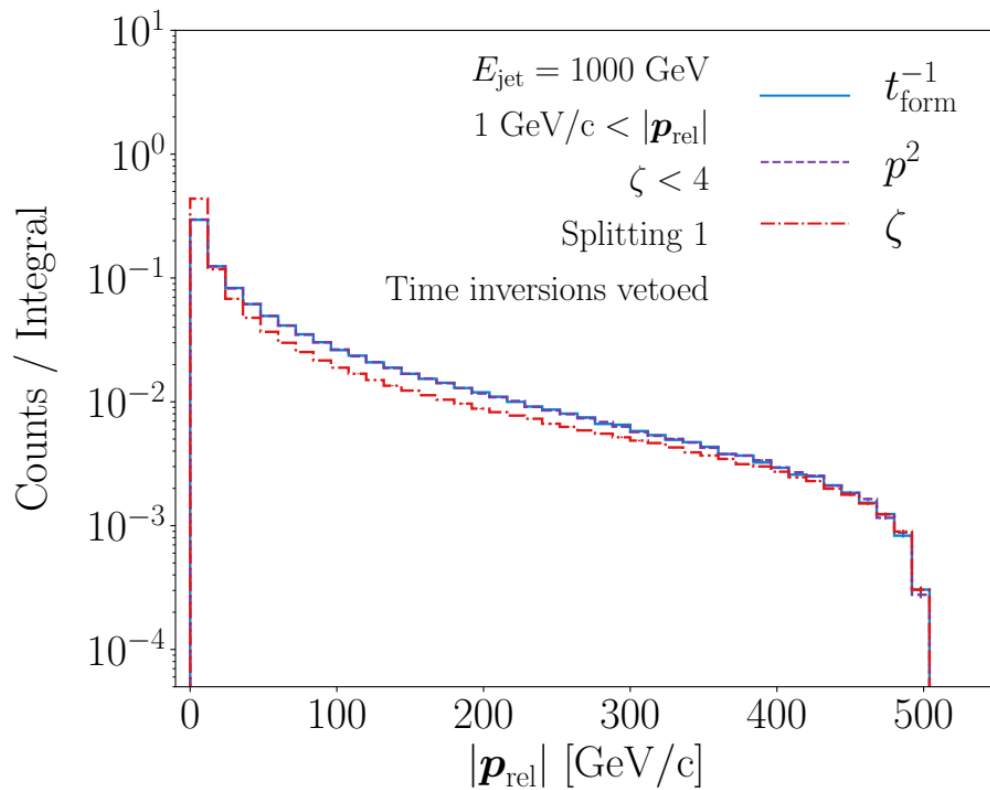
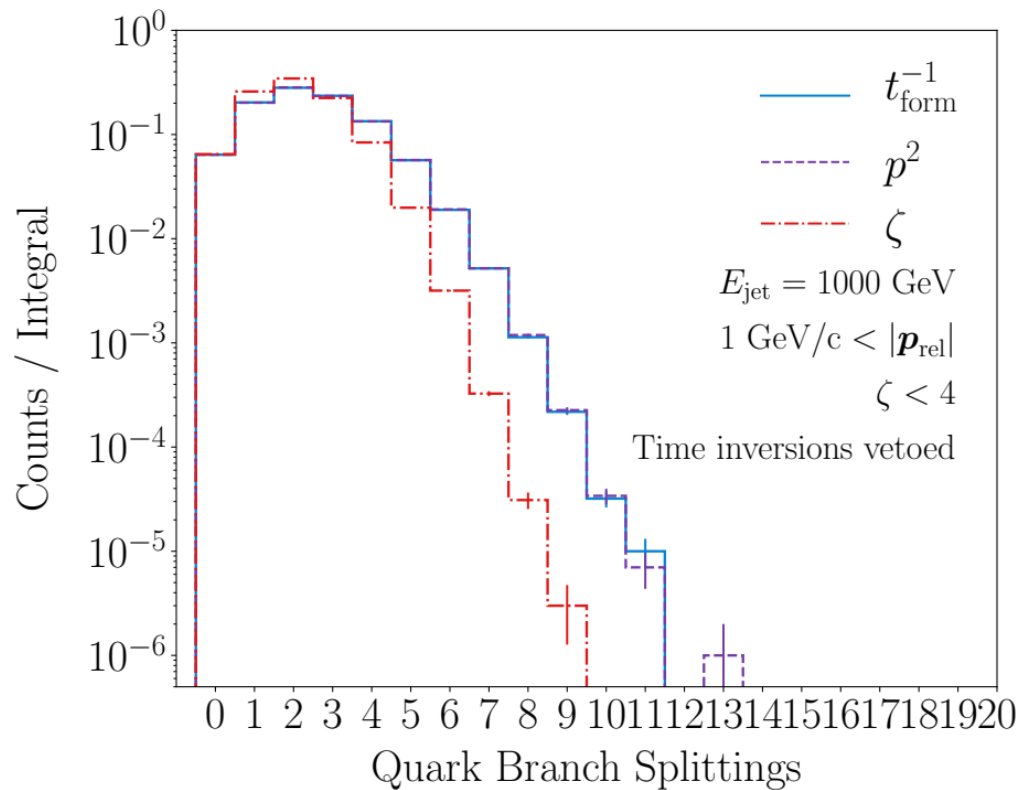


# Inclusive Sample - 1D Distributions



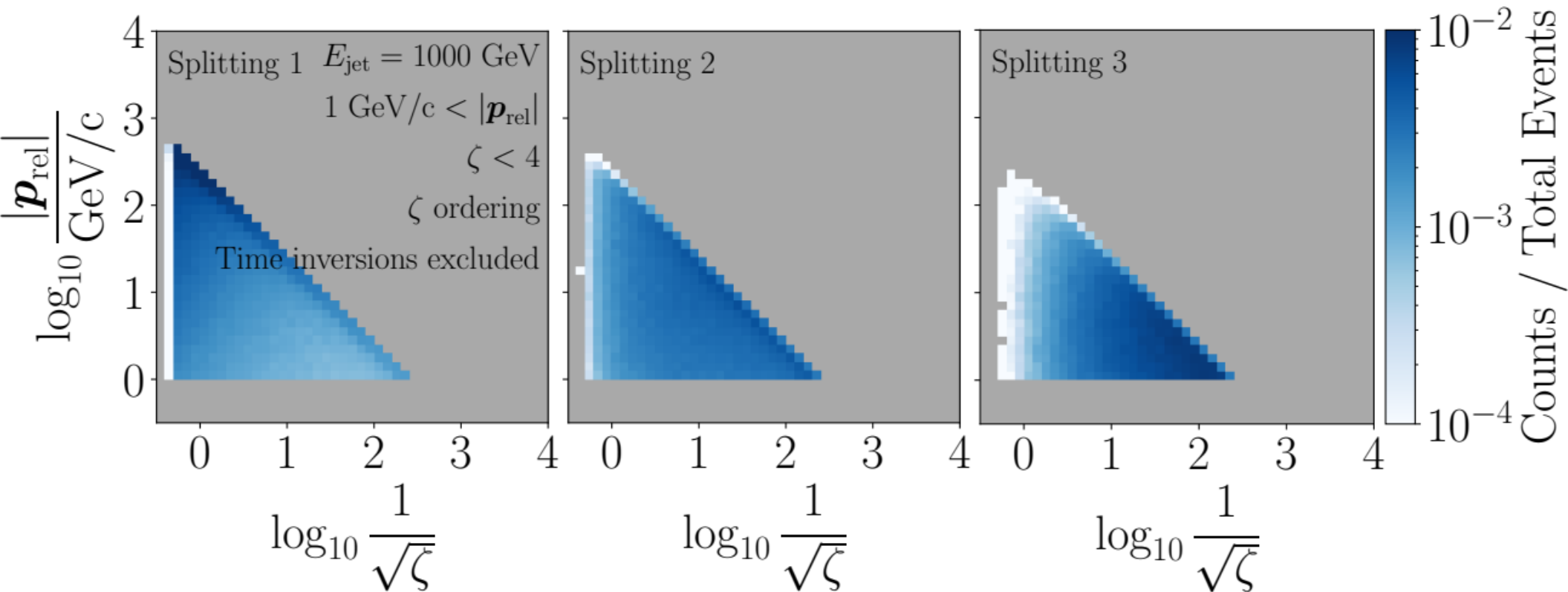


# Vetoing time inversions - 1D Distributions



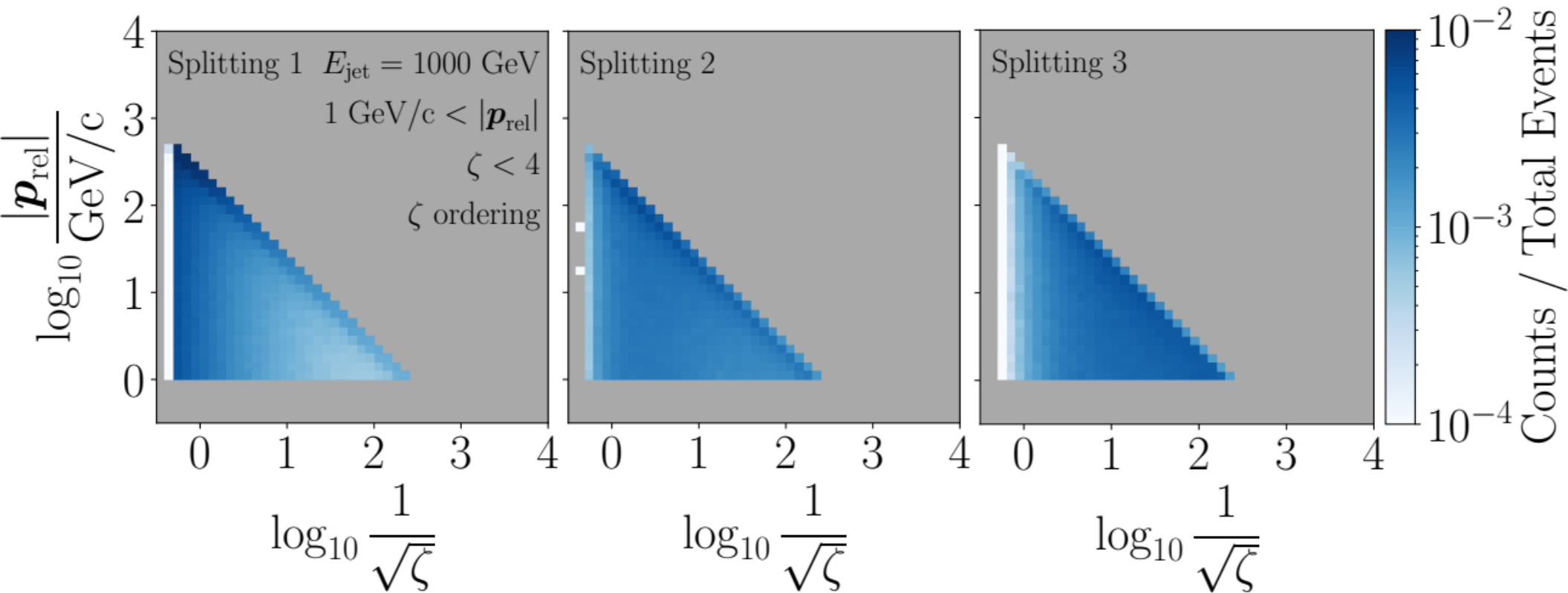
# Excluding time inversions - Lund Planes

\*Ordered in angle



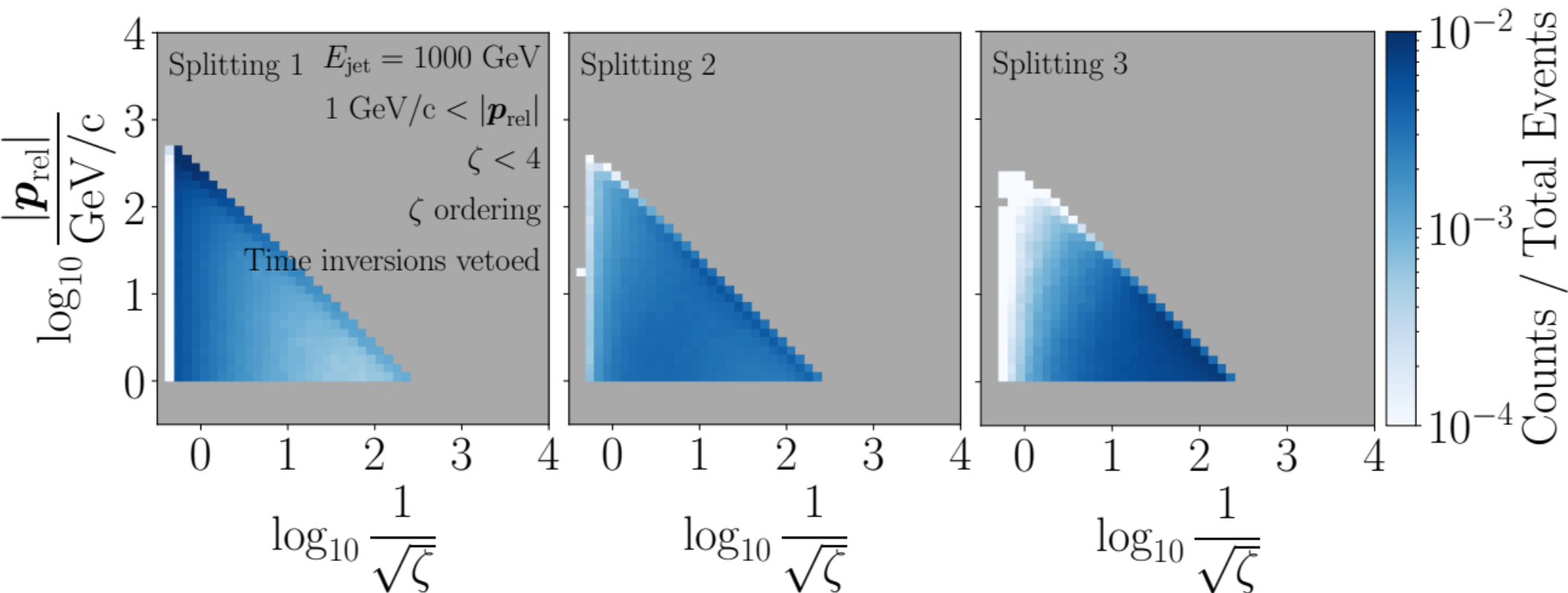
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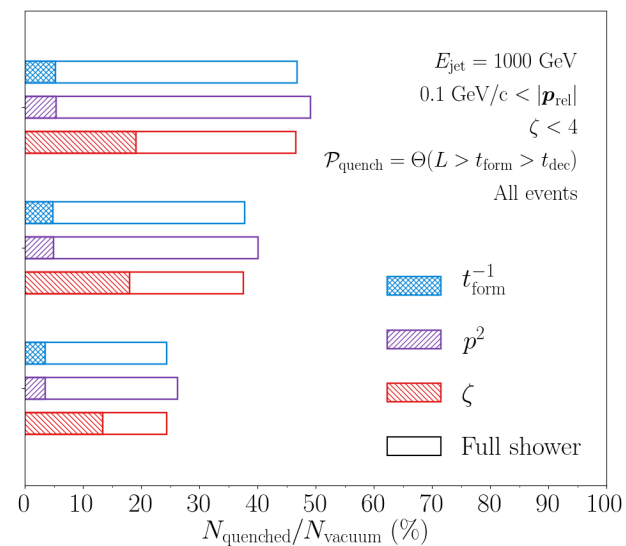
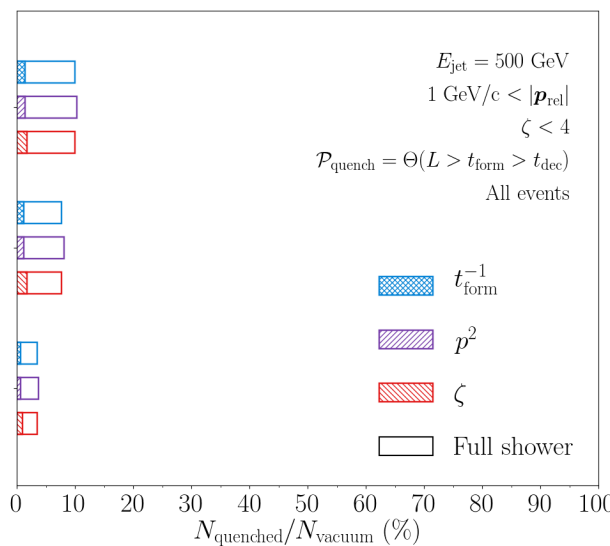
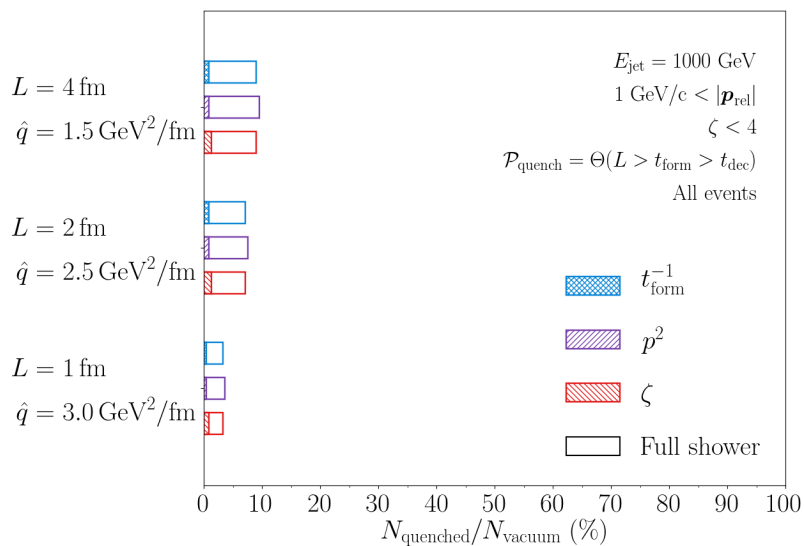


# Quenching Weights

$E_{\text{jet}} = 1000 \text{ GeV}, \Lambda = 1 \text{ GeV}$

$E_{\text{jet}} = 500 \text{ GeV}$

$\Lambda = 0.1 \text{ GeV}$



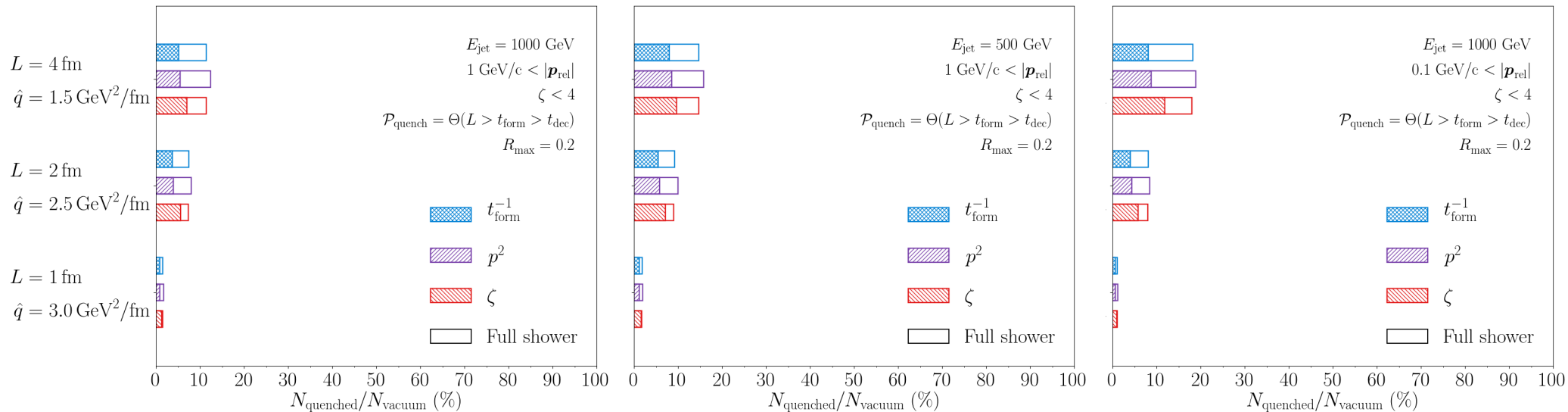
An apparent dependence on the hadronisation cutoff and initial jet energy

# Quenching Weights – Radius Cut

$E_{\text{jet}} = 1000 \text{ GeV}, \Lambda = 1 \text{ GeV}$

$E_{\text{jet}} = 500 \text{ GeV}$

$\Lambda = 0.1 \text{ GeV}$



**Cut all events whose quark branch has a splitting wider than  $R_{\text{max}} = 0.2$**   
**- This defines the new vacuum sample, and the quenching model is applied on top of this cut**

**An aggressive cut, but it returns independence of  $E_{\text{jet}}$  and  $\Lambda$ .**